

1 **An assessment of the role of surface sensible heat flux and the atmosphere**
2 **inversion on the breakup time in a highly complex terrain**

3 *

ABSTRACT

4 PENDING

5 **1. Introduction**

6 The formation of nocturnal Stable Boundary Layers (SBLs) and temperature inversions has im-
7 portant implications for urban areas, particularly those situated in complex terrain valleys. These
8 range from physical processes, including the modulation of momentum, heat, and moisture ex-
9 change, to practical ones associated with pollutant accumulation near the surface, as well as the
10 formation of Urban Heat Islands (UHI). Most pollutants emitted during the evening and through-
11 out the night, together with those that return to the surface due to the top-down contraction of
12 the residual layer, remain trapped in the valley atmosphere at least until the temperature inversion
13 ceases (e.g. Doran et al. 2003; Schnitzhofer et al. 2009; Fernando et al. 2010; Saide et al. 2011;
14 Liu et al. 2013; Hu et al. 2013; Rendón et al. 2014; Herrera-Mejía and Hoyos 2019). Depending
15 on valley geomorphological characteristics, and the meteorological conditions surrounding the
16 morning transition, pollutants may exit the valley atmosphere or recirculate. The latter could lead
17 to a gradual deterioration of the air quality, and in some cases, to the onset of critical air pollu-
18 tion episodes (Anquetin et al. 1998; Angevine et al. 2001; Henne et al. 2004; Rendón et al. 2015;
19 Czarnecka et al. 2019).

20 The inversion destruction, also referred to as the *inversion breakup* or the *inversion erosion*,
21 is linked to the occurrence of the necessary and sufficient conditions for the development of the
22 Convective Boundary Layer (CBL). The onset of the CBL leads to an efficient energy, moisture,
23 and pollutant exchange between the surface and the free atmosphere (Schnitzhofer et al. 2009;
24 Leukauf et al. 2015). Establishing under which conditions the nighttime inversion breaks up is not
25 a trivial issue, especially for urbanized valleys, where the complexity of the terrain and the urban
26 landscape directly influence the evolution of the SBL (Halios and Barlow 2018). Factors such as
27 valley geometry, which induces topographic shading, soil water holding capacity and moisture,

28 urban area fraction and roughness, and the overall valley circulation, are key in governing the
29 inversion breakup in complex-urban environments. Some authors have explored the influence of
30 these factors in the morning transition through both modeling (Whiteman and McKee 1982; Bader
31 and McKee 1985; Colette et al. 2003; Whiteman et al. 2004; Zoumakis and Efstathiou 2006;
32 Beare 2008; Rendón et al. 2014; Leukauf et al. 2015), and field experiments (Whiteman 1981;
33 Angevine et al. 2001; Halios and Barlow 2018; Nadeau et al. 2018), highlighting the need for
34 a better understanding and representation of the physical processes controlling the timing of the
35 transition for improving numerical weather prediction and air pollution models (Angevine et al.
36 2001; Colette et al. 2003; Beare 2008; Rendón et al. 2014).

37 Whiteman (1981) and Whiteman and McKee (1982) categorized the inversion-breaking pro-
38 cesses over mountainous valleys, predominantly rural, as a function of the interaction of two
39 mechanisms. The first mechanism is considered local, and corresponds to the CBL growth from
40 the bottom of the inversion resulting from radiative heating of the surface (Nadeau et al. 2018).
41 The second mechanism depends on the valley circulation, in which the replacement of air masses
42 at the valley bottom with air from the upper atmosphere through slope winds, causes the top of the
43 inversion layer to collapse. Depending on the relative role of each of the described mechanisms,
44 the breakup falls into three possible patterns, the first two resulting from the mechanisms previ-
45 ously mentioned, acting independently, and the third and last pattern to the simultaneous action of
46 both mechanisms (Whiteman 1981). The idealized modeling experiments by Bader and McKee
47 (1985) also suggest an essential role of both mechanisms in thermal inversion destruction.

48 Whiteman et al. (2004) reached different conclusions for high-latitude sinkholes, for which the
49 evidence suggests that inversion destruction occurs mainly through subsidence warming, indirectly
50 forced by the upslope flows. Valley geometry has a strong effect both on forcing the prevailing
51 breakup pattern and on the timing of the inversion breakup, with inversions persisting longer in

52 deeper valleys (Colette et al. 2003). The presence of extensive urban development over complex
53 topography leads to alterations in the surface-atmosphere exchanges, enhancing valley floor heat-
54 ing through the formation of UHI. The induced heat due to the UHI may intensify the thermal
55 turbulence production, accelerating the onset and growth of the CBL (Roth 2000; Rendón et al.
56 2015). Furthermore, slope winds tend to increase in magnitude as a consequence of strong tem-
57 perature gradients appearing across the urbanized hills. The latter favors the inversion-breaking
58 through the second and third breakup patterns. In addition to the above-mentioned processes, the
59 evidence hints that wind shear plays a vital role in driving the Turbulent Kinetic Energy (TKE),
60 and in the morning transition (Beare 2008).

61 From a broader perspective, regardless of terrain characteristics, and in terms of energy balance,
62 the inversion breakup occurs when the total energy provided to the valley atmosphere (Q_{prov} ,
63 following the notation used in Leukauf et al. (2015)) via surface sensible heat flux (H) is equal
64 to the energy required to erode the nocturnal SBL (Q_{req}) (Whiteman and McKee 1982; Angevine
65 et al. 2001; Leukauf et al. 2016). Following the inversion breakup, the additional energy injected
66 into the valley atmosphere is used, in part, to expand the CBL until the exchanges between the
67 surface and the free atmosphere peak, resulting in a more efficient pollutant vertical transport.
68 Leukauf et al. (2016) proposed a non-dimensional *breakup parameter* (B) defined as the ratio
69 between Q_{req} and Q_{prov} , combining the effect of the atmospheric stability and the surface heating
70 to assess the energy exchange processes. Leukauf et al. (2017) performed simulations, using the
71 Weather Research and Forecasting (WRF) model (Skamarock et al. 2008), in different energy
72 exchange scenarios characterized by different surface heating, initial atmospheric stability, and
73 terrain geometry, to test the dependence of the exported heating on B . Leukauf et al. (2017) found
74 that the amount of heat exported from the valley decreases exponentially as B increases and that
75 there is a critical condition over which the complete SBL neutralization never occurs.

76 Despite the importance of the relationship between Q_{req} and Q_{prov} to understand the timing of the
77 breakup and its potential modulating pollutant concentration, particularly for urbanized valleys,
78 there is insufficient observational evidence of this effect. Recently Halios and Barlow (2018)
79 studied the morning transition using ground-based remote sensing and in situ instrumentation in
80 central London, finding that buoyant production of TKE at the surface and shear production in
81 the upper half of the atmospheric boundary layer (ABL) erode the stable layer. They highlighted
82 the importance of regional flows, such as low-level jets, in determining the urban boundary layer
83 structure and growth.

84 The goal of this research is to gain insight into the above mentioned processes, based on the study
85 of the trade-off between observed proxies of the energy provided to the valley atmosphere as H ,
86 and of the energy required to erode the nighttime inversion, both key in the inversion-breaking
87 process in a low-latitude, highly urbanized valley. We also investigate the practical implications
88 of the relative variability of these two proxies regarding the air quality and explore the efficiency
89 of the energy injected into the atmosphere. Previous results suggest a substantial modulation of
90 the local air pollutant concentration associated with ABL variability (Herrera-Mejía and Hoyos
91 2019). We analyze nine months of ground based remotely sensed thermodynamic profiles and
92 in situ observations, including data from a microwave radiometer, a radar wind profiler, a 3D
93 sonic anemometer, automatic weather stations, and air quality monitoring sensors. Furthermore,
94 we evaluate the role of the valley circulation, primarily the vertical wind shear, and the regional
95 meteorological conditions at different levels in the atmosphere, on the SBL erosion efficiency.

96 Section 2 presents a detailed description of the study area, the observational dataset, and the
97 different thermodynamic proxies used for the assessment of Q_{req} and Q_{prov} . Section 3 includes the
98 results of the relationship between the selected Q_{req} and the Q_{prov} . The section explores the SBL
99 erosion efficiency, its dependence on local and regional meteorological conditions, the implica-

100 tions on air quality, and the intra-annual variability associated with the meteorological conditions
101 in the valley. Finally, section 4 presents the most important conclusions of the study.

102 **2. Methodology and observations**

103 The observational assessment of turbulent fluxes, vertical structure of virtual potential tempera-
104 ture (θ_v) and wind speed, regional meteorology, and air quality measurements, allows studying the
105 variability of the inversion breakup as a function of proxies representing Q_{req} and Q_{prov} , as well
106 as the SBL erosion, and the impacts on local particulate matter (PM) concentration in a narrow,
107 low-latitude valley.

108 The methodology includes i) the assessment of the magnitude and the intra-diurnal evolution of
109 the proxies for Q_{req} , ii) the evaluation of the surface H as a proxy for Q_{prov} , iii) the study of the
110 breakup time variability and the SBL erosion, iv) the assessment of the role of local and regional
111 meteorology on the SBL erosion efficiency, and v) the estimation of the breakup time impact on
112 the fine PM (PM_{2.5}) concentration near the surface. We also examine the Q_{req} - Q_{prov} relationship
113 from operational weather forecasts to explore whether the WRF model reproduces the observed
114 relationship under realistic simulation conditions.

115 We use in situ and ground-based remote sensing observations, as well as satellite and reanalysis
116 information from February 1 to November 21, 2018. Ground-based instrumentation is located
117 along and across the region of interest, a highly urbanized low-latitude narrow valley. Although
118 most of the data is available for a more extended time span, the analysis period is restricted by
119 the availability of turbulent fluxes. In the analysis, all days with available data are considered,
120 regardless of the meteorological conditions; in other words, we do not pre-select fair-weather days
121 for the analysis. In a low-latitude environment, such days correspond to less than 1% of the data,

122 which would restrict the study to very few cases, focusing on the exceptional circumstances and
123 not the expected scenarios. The datasets comprise wet, transition, and dry seasons.

124 To better understand when, and under which conditions, the inversion breakup occurs, we start
125 by assessing and comparing the characteristics of the stably stratified atmosphere of the Aburrá
126 Valley before sunrise, using thermodynamic profiles obtained from a Microwave Radiometer
127 (MWR), and the energy provided to the valley's atmosphere in the form of H . The thermodynamic
128 profiles allow establishing a proxy for the amount of energy required to erode the stably stratified
129 boundary layer, until the breakup occurs (Q_{req}). H is estimated using the eddy-covariance (EC)
130 technique based on the turbulent fluctuations of the wind speed, temperature, and humidity mea-
131 sured using a 3D sonic anemometer. The overall approach combines high frequency measurements
132 near the surface with macroscopic observations of the atmosphere in the vertical profile.

133 The inversion breakup assessment using a data analysis approach involves challenges associated
134 with the spatial representativeness of each of the variables considered in the study. The latter is
135 arguably the main reason why similar studies in the literature follow a modeling-based approach
136 rather than an observational one. Turbulent fluxes estimated from in-situ observations generally
137 represent the local-scale variability conditioned to the intrinsic heterogeneities of the terrain. In
138 contrast, the thermodynamic profiles represent the macroscopic features of the atmosphere. There-
139 fore, it is necessary to evaluate whether the observational datasets used in this study are suitable
140 for the primary purpose and whether they reproduce the overall dynamics around the morning
141 transition. Notwithstanding the potential limitations, the analysis using independent and different
142 nature datasets constitutes a robust assessment of the inversion breakup process. The covariabil-
143 ity among the datasets used in the study would imply a coherent response or connection among
144 different atmospheric scales considered.

145 *a. Geographical Context*

146 The inversion breakup is studied using information from the Aburrá Valley. The valley is located
147 in Colombia, in the Andes Cordillera between 6°N and 6.5°N and 75.3°W and 75.6°W (see Figure
148 1) and it is aligned predominantly south-to-north, south-to-northeast. The widest cross-section of
149 the valley, from ridgeline to ridgeline, is 18.2 km, with a relatively flat section of approximately 8
150 km at the bottom. The narrowest section is around 3 km wide. The highest peak, approximately at
151 3110 m.a.s.l, is in the western hill. The basin outlet is at 1290 m.a.s.l.

152 From the point of view of turbulent exchanges, the complexity of the region is due to the rugged
153 topography and the highly urbanized area. More than four million people are settled in an area
154 of 1152 km². Additionally, the urban area reaches, in some cases, three-quarters of the hill-slope
155 extension. The high population density in this geographic setting leads to several environmental
156 challenges. These include the recurrent onset of critical air quality episodes due to the high demand
157 for fossil fuels linked to motor vehicles and industries and the limited ventilation of the valley's
158 atmosphere for pollutant dispersion.

159 *b. Proxies for Q_{req}*

160 An accurate determination of Q_{req} depends, first, on a precise theoretical definition of the time
161 at which the atmosphere becomes neutrally stratified, establishing the initial conditions for the
162 boundary layer growth phase. Following the work by Angevine et al. (2001), for flat terrain,
163 numerous authors (e.g. Beare 2008; Nadeau et al. 2018) have defined the inversion breakup as the
164 onset of the CBL, the time at which the nocturnal inversion in the surface layer has been eroded
165 and turbulent eddies from the surface reach a certain depth (e.g., Angevine et al. (2001) used 200m
166 above ground level). To some extent, the size of the eddies may be assessed observationally.

167 The appropriate selection of a proxy for the Q_{req} is not straightforward, mainly because such
 168 an index should adequately represent the entire volume of the valley's atmosphere. We consider
 169 thermodynamic indices such as the change of θ_v in the vertical at the lower-troposphere ($\Delta_z\theta_v$)
 170 and the Convective Inhibition Energy (*CINE*). $\Delta_z\theta_v$ has been extensively used as a proxy for the
 171 inversion strength (e.g. Whiteman 1981), considering that $\partial\theta_v/\partial z > 0$ corresponds to stable strat-
 172 ification, $\partial\theta_v/\partial z = 0$ to neutral conditions, and $\partial\theta_v/\partial z < 0$ to an unstable atmosphere (Peppler
 173 1988; Curry and Webster 1999). Whiteman (1981) calculate $\Delta_z\theta_v$ as the difference between θ_v
 174 near the surface and at the top of the inversion. We consider $\Delta_z\theta_v$ for different atmospheric layers,
 175 where z corresponds to the height in meters above the surface. $\Delta_z\theta_v$ is computed as the difference
 176 between θ_v at height z and θ_v at height $z - \Delta z$, $\theta_v(z) - \theta_v(z - \Delta z)$. We consider $\Delta z = 200\text{m}$ thick
 177 layers, with the only exception for $\Delta_{200}\theta_v$, computed as $\theta_v(200) - \theta_v(50)$ to avoid the potential
 178 effects of the roughness sublayer.

179 Furthermore, *CINE* indicates the amount of energy inhibiting the updraft of air parcels, and is
 180 also an indirect measurement of the lower troposphere stability: as the stability of the atmosphere
 181 increases, *CINE* becomes more negative. Conversely, unstable atmospheres correspond to *CINE*
 182 nearing zero. *CINE* is estimated as

$$CINE = \int_{SFC}^{LFC} g \frac{T'_v - T_v}{T_v} dz,$$

183 where *LFC* is the level of free convection, *SFC* is the surface level, T_v is the virtual temperature
 184 of the environment, and T'_v is the virtual temperature of the parcel (Peppler 1988). On occasions,
 185 the *LFC* largely exceeds the depth of the valley, where the trade winds advect eastward the *H*
 186 and the pollutants emitted at the surface. As a result, the amount of energy required for *CINE* to
 187 become zero is larger than the amount of energy required to erode the stability within the valley.
 188 Also, the mechanisms that lead to changes in the thermodynamic profile above the valley may not

189 be fully linked to the turbulent exchanges near the surface, but the forcing could be associated
190 with the synoptic scale. To address this potential issue, and to have more control over the *CINE*
191 integration height, the *LFC* is forced to a maximum of 1200 m, which is the average depth of the
192 valley (see Figure 1). The modified index is referred to as *CINE*₁₂₀₀.

193 The proxy for Q_{req} , on a daily time scale, corresponds to the maximum $\Delta_z\theta_v$, and the mini-
194 mum *CINE* and *CINE*₁₂₀₀ after sunrise, considering the topographic shading. $\Delta_z\theta_v$ and *CINE* are
195 computed using thermodynamic profiles obtained using a MWR.

196 MWR DATA

197 An MP-3000A Microwave Radiometer (MWR), manufactured by Radiometrics, measured the
198 vertical profiles of temperature and relative humidity in the Aburrá Valley up to 10 km from the
199 top of SIATA's main operation center, approximately 60 m above the surface (see Figure 1a). The
200 site is inside a sports complex, surrounded by different types of urban land uses. The MWR is a
201 passive remote sensor that measures the radiation emitted by atmospheric gases using 31 different
202 submillimeter-to-centimeter wavelengths. The MWR is useful for the retrieval of the thermody-
203 namic state of the atmosphere at different levels, thus allowing the assessment of atmospheric
204 stability. The MWR provides vertical profiles with a 2-minute temporal resolution and variable
205 spatial resolution: 50 m from the surface to 500 m, 100 m up to 2 km, and 250 m up to 10 km. The
206 lower-troposphere retrievals (below 4 km above the surface) are used to calculate the proxies for
207 Q_{req} . The representativeness of the temperature and moisture profiles obtained using the MWR
208 has been previously assessed using radiosonde measurements, showing high correlations for all
209 the variables, including θ_v and $\Delta_z\theta_v$, in the lower troposphere (Roldán-Henao et al. 2020).

210 *c. Inversion breakup time*

211 Based on the different proxies for Q_{req} , the breakup time is assumed to have occurred when $\Delta\theta_v$
212 becomes zero after having been positive, or when $CINE = 0$ or relatively close to zero (we use
213 $CINE < 20 \text{ J kg}^{-1}$ as a threshold). Figures 2a and b depict the steps used to assess the strength
214 of the stability (Q_{req}) and the breakup time for a particular day, using $\Delta_{200}\theta_v$. In the example, the
215 proxy for the strength of the stability is recorded as the maximum positive $\Delta_{200}\theta_v$ after sunrise,
216 which corresponds to the gradient that must be neutralized to reach instability. According to Figure
217 2a, the value representing the strength of the instability for that particular date is 1.26 K. Figure
218 2b marks the breakup time, after 11:00 LT, corresponding to the moment when $\Delta_{200}\theta_v = 0$.

219 *d. Proxy for Q_{prov}*

220 The proxy used for the daily Q_{prov} (see Figures 2c) is the time-integrated surface H from the
221 moment used to record the strength of the stability (maximum Q_{req} after sunrise) until the inver-
222 sion breakup (see Figures 2a and b, respectively). The proxy is based on a modified version of the
223 methodology presented in Angevine et al. (2001), which is based on the boundary layer growth
224 equation in Garratt (1992). The methodology assumes that the temperature profile is mainly modi-
225 fied from the bottom as a function of sensible heating. This encroachment approach may introduce
226 biases, given that it does not account for the θ_v profile modification in the lower troposphere, hence
227 the stability modulation due to horizontal thermal advection. However, the evidence suggests that
228 θ_v in the lower-troposphere, 500 m above ground, mainly varies as a result of vertical processes,
229 with temperature changes lagging those in layers closer to the ground (not shown).

230 The H is estimated using the EC technique with a block-averaging period of 30 minutes. An EC
231 tower equipped with a CSAT3 ultrasonic anemometer (Campbell Scientific) is used to obtain the
232 three wind components (u' , v' , and w') and sonic temperature (T'_s) with a sampling rate of 20Hz.

233 The instrument is installed 10 m above the surface, in a mast located next to the local airport (see
234 Figure 1a). The absence of tall buildings in the surrounding area prevents the data from being
235 strongly affected by the local circulation. Raw data are stored at full resolution in 24-hour files,
236 and the statistical first- and second-order moments are calculated using 30-min block periods. Af-
237 ter applying standard data quality controls (periods flagged by the instrument, checks for large
238 data gaps and consistency limits, and data despiking), a coordinate double-rotation was applied
239 for each 30-min interval to obtain the along-wind u and cross-wind v components. The latter en-
240 sures that the magnitude of the mean lateral and vertical components of the velocity vanishes for
241 each of the averaging interval (e.g. McMillen 1988; Finnigan et al. 2003; Stiperski and Rotach
242 2016). Additional post-processing procedures include flux corrections (Webb et al. 1980; Moore
243 1986; Schotanus et al. 1983), and assessing of non-stationarity (following Foken and Wichura
244 (1996)). A detailed description of the post-processing steps will be included in a separate publica-
245 tion (Herrera-Mejía et al. in preparation).

246 *e. Role of breakup time in air quality*

247 The breakup time is closely related to thermodynamically driven vertical dispersion processes.
248 Consequently, the ABL neutralization may play a vital role in pollutant concentration. The assess-
249 ment of the impact of the inversion breakup time on the air pollution near the surface of the valley
250 is accomplished through the study of conditional probability density functions (CPDFs) of PM_{2.5}
251 concentration as a function of the breakup time. In the Aburrá Valley, PM_{2.5} is the most critical air
252 pollutant. All of the criteria air pollutants defined by the United States Environmental Protection
253 Agency (U.S. EPA), except for lead, are routinely monitored in the region, in a 37-station compre-
254 hensive and accredited monitoring network. For this study, data from four in-situ PM_{2.5} stations
255 equipped with a U.S. EPA Federal Equivalent Method (FEM) Met One Instruments BAM-1020

256 monitor is used. These stations are located along the base of the valley, some of them intentionally
257 selected far from the location of the MWR and the CSAT3 (see Figure 1a), so as to indirectly test
258 the representativeness of the proxies for Q_{req} and Q_{prov} for the entire valley. Retrievals from three
259 Vaisala CL51 ceilometers (910 nm wavelength) are used to illustrate the structure of the vertical
260 profile of aerosols in different Q_{req} scenarios. The ceilometer used is installed at the same site
261 as the MWR (see Figure 1a). Ceilometers provide information regarding the laser-pulse energy
262 backscattered by clouds and other atmospheric components, including aerosols, expressed as the
263 backscattering attenuated coefficient (Emeis et al. 2009; Kambezidis et al. 2012; Wiegner et al.
264 2014).

265 *f. Role of local and regional meteorology*

266 One of the challenges in better understanding the ABL in complex terrain, and in particular, the
267 morning transition, relates to the multiscale nature of the processes that modulate the phenomena.
268 The overall behavior of the atmosphere at different levels exhibits signs of multiscale interaction,
269 both in time and space. This is particularly true for the ABL over complex terrains, where the
270 diurnal cycle, including the transition from the SBL to the CBL, is modulated by processes evol-
271 ving on different temporal and spatial scales (e.g. Serafin et al. 2018; De Wekker and Kossmann
272 2015). The large-scale modulation of the ABL could occur directly through kinetic energy transfer
273 between different scales of motion or indirectly via regional and large-scale changes in the con-
274 ditions that favor or inhibit cloud formation, leading to surface radiative forcing. For this reason,
275 it is essential to explore not only the turbulent exchanges, but also the potential role of the valley
276 flow and the synoptic atmospheric circulation on the breakup time.

278 The methodology includes the evaluation of the contribution of the synoptic scale in modulating
 279 the morning transition. In the tropical environment, pressure patterns persist for very long periods,
 280 and the pressure and geopotential height gradients are typically weak, even during the passage of
 281 storm systems or perturbed weather in general. For this reason, the wind field is more useful than
 282 the pressure or geopotential height fields in describing the synoptic conditions in the tropics. In
 283 this regard, the velocity potential anomalies summarize the effect of synoptic-scale convection,
 284 allowing the tracking of upper-level divergence or convergence.

285 ERA5 reanalysis data (Hersbach et al. 2020) is used to calculate the velocity potential and the
 286 stream function corresponding to the reanalyzed wind fields. For the estimation of both variables,
 287 the wind field is separated into two components, the rotational (\vec{V}_{rot}) and the divergent component
 288 (\vec{V}_{div}) of the flow. The divergent winds are used to calculate a velocity potential (χ), satisfying
 289 that the winds flow out low potentials and their speed is proportional to their gradient ($\vec{V}_{div} = \nabla\chi$).
 290 In other words, χ is obtain as the solution to the equation $\nabla^2\chi = \nabla \cdot \vec{V}$ (Krishnamurti et al. 2013;
 291 Laing and Evans 2015). We also use the Outgoing Longwave Radiation (OLR) from the NOAA
 292 daily interpolated dataset (Liebmann and Smith 1996), to assess the role of synoptic forcing on
 293 the evolution of the intra-valley ABL.

294 In the assessment, three different atmospheric scenarios are considered, corresponding to cases
 295 when the nighttime inversion is strong (high Q_{req}) and i) the magnitude of the energy forcing
 296 provided to reach the morning transition via sensible heating is lower than the 33th percentile
 297 among all days considered in the study, ii) times when the energy provided to reach the transition
 298 is between the 33th and the 66th percentile, iii) and days for which the energy provided to the
 299 system, before reaching the transition, is larger than the 66th percentile.

300 RADAR WIND PROFILER (RWP) DATA

301 The RWP uses refractive index variations caused by changes in humidity, temperature, and
302 pressure, to retrieve vertical profiles of winds (Lau et al. 2013). The Aburrá Valley wind profiler,
303 a RAPTOR VAD-BL by DeTect Inc., works at a nominal frequency of 1290 MHz, reaching up to
304 approximately 8 km above the surface under high humidity conditions. The RWP is designed to
305 measure the wind profile in various operation modes that differ in their vertical resolution, as well
306 as in the atmospheric domain sensed. The operation of the RWP includes two overlapping modes:
307 in the higher resolution mode (60 m), the RWP retrieves the wind profile from 77 to 3500 m, and
308 in the lower resolution mode (72 m), from 2500 to 8000 m. In the present study, only data from
309 the higher resolution mode is used. The temporal resolution is five minutes.

310 WRF

311 We used the output of the operational daily 00Z WRF (version 3.7.1) 24-hour forecasts for three
312 years. The model configuration includes three nested domains with 18 (191 x 191), 6 (82 x 118)
313 , and 2 (136 x 136) *km* grid spacing, and 40 vertical levels up to 50 *hPa*. The description of the
314 domains and the model setup are described in detail in Herrera-Mejía and Hoyos (2019).

315 **3. Results**

316 One of the main challenges in assessing the inversion breakup from an observational perspective
317 is to ensure that measurements from in-situ sensors and ground-based remote sensing equipment
318 represent the overall ABL variability in the valley. The latter is not only a challenge but a source
319 of uncertainty in all ABL observational studies. While this challenge is difficult to overcome, it is
320 possible to evaluate the holistic coherence and consistency among all variables in the dataset, and
321 their capability to represent the ABL dynamics. Considering that most variables used in this study

322 are obtained using different measurement techniques, high co-variability and interdependence in
323 the dataset would indicate a skillful representation of the zeroth- and first-order valley dynamics,
324 including the major spatial and temporal scales of variability.

325 Figure 3a depicts the interrelationship between the hourly H , $\Delta_{200}\theta_v$, and $CINE_{1200}$. Evidently,
326 negative values of $\Delta_{200}\theta_v$ correspond to near-zero values of $CINE_{1200}$, and conversely, positive
327 values of $\Delta_{200}\theta_v$ are associated with negative values of $CINE_{1200}$. The Pearson correlation be-
328 tween $\Delta_{200}\theta_v$ and $CINE_{1200}$ is -0.80. Figure 3b shows the correlations among all the Q_{req} proxies
329 considered, including $\Delta_{200}\theta_v$, $\Delta_{800}\theta_v$, $\Delta_{Total}\theta_v$ ($\theta_v(1200) - \theta_v(50)$), $\Delta_{Sup}\theta_v$ ($\theta_v(800) - \theta_v(50)$),
330 $CINE$, $CINE_{1200}$, and $CINE_{1500}$. The magnitude of the linear correlations among all variables,
331 except between $CINE$ and $\Delta_{Sup}\theta_v$, is over 0.6, emphasizing the high covariance of the virtual
332 temperature in the lower levels of the troposphere, below LFC . This result serves as additional
333 evidence of the strong dependence of the temperature profile on the surface heating, rather than on
334 thermal advection.

335 Figure 3a shows that, for strong surface forcing to the ABL (high values of H), the likelihood of
336 positive values of $\Delta_{200}\theta$ is very low. In other words, it is unlikely to have stable stratification when
337 the heating is strong. Correspondingly, when H is larger than 150 Wm^{-2} , the average of $CINE_{1200}$
338 is -7.5 Jkg^{-1} . The large spread of the data at the bottom portion of the H - $\Delta_{200}\theta$ scatterplot is a
339 consequence of the fact that the transition from a stable to an unstable atmosphere is a cumulative
340 process, which does not depend exclusively on the magnitude of the surface forcing at a given
341 time, but also on multiple additional factors. Furthermore, the observed relationship between H
342 and the incoming radiation is very high, with a correlation of 0.89. Despite the different nature of
343 the physical principles used to measure the different variables, the high covariability corresponds
344 well with the expected behavior of the ABL, suggesting a clear link between the radiative forcing
345 and H with the evolution of the nocturnal inversion within the valley. Based on these results,

346 it is possible to follow the described observational approach. Additionally, considering the high
347 correlations in 3b, the subsequent results are obtained using two proxies of Q_{req} : a near-surface
348 stability proxy ($\Delta_{200}\theta_v$), and a lower-troposphere column integrated proxy ($CINE_{1200}$). Results
349 using other proxies are similar and do not alter the main conclusions of this study.

350 *a. Q_{req} Vs. Q_{prov}*

351 Figures 4a and b show the relationship between Q_{req} , using $CINE_{1200}$ and $\Delta_{200}\theta_v$ as proxies,
352 respectively, and Q_{prov} prior to the temperature profile neutralization, as defined. Each point on
353 the scatterplot corresponds to a specific day between February and November 2018. Both diagrams
354 show a remarkable correspondence between the two selected proxies for Q_{req} , providing evidence
355 that Q_{prov} indeed is required to be higher when the magnitude of Q_{req} is large, regardless of the
356 proxy used. The Q_{req} - Q_{prov} relationship is not linear. There appears to be a threshold in the
357 strength of the inversion (Q_{req}), over which there is a considerable spread in the Q_{prov} before
358 neutralization, implying that, in some cases, for the same Q_{req} the magnitude of Q_{prov} could be
359 four to six times larger than usual before achieving SBL neutralization. The latter suggests the
360 existence of a heating efficiency similar to the findings of Leukauf et al. (2017). This is explored
361 further in subsection b.

362 Figure 5, similar to 4b but calculated using information from the WRF forecast runs. The dia-
363 grams for the WRF runs show, in general, a similar behavior to the observations, but with a larger
364 spread Q_{prov} for large Q_{req} , hinting to a larger variability in the heating efficiency in the models.

365 The relationship observed in both diagrams in Figure 4 suggests a different state of the atmo-
366 sphere for cases corresponding to the lower and upper parts of the scatter plots. To further explore
367 this behavior, Figure 6 shows the comparison of the state of the atmosphere on two contrasting
368 days, corresponding to the larger circles in Figure 4. The first case, with a high Q_{req} , corresponds

369 to February 22, 2018 (see Figures 6a, b, c and, d), and the second case, with a low Q_{req} , corre-
370 sponds to October 12, 2018 (lower panels) (see Figures 6e, f, g, and, h). The Figure includes the
371 evolution of the θ_v profile, from 05:00 LT to 14:00 LT, the vertical profile of wind speed and wind
372 direction, the ceilometer backscattering intensity profile from the surface up to 2.5 km, and finally,
373 the hourly evolution of PM2.5 concentration.

374 The θ_v profiles reveal a strong nighttime inversion on February 22, 2018 (Figure 6a), resulting
375 in a considerable amount of energy required to erode the SBL, a notably shallow ABL, and a late
376 breakup time. The shallow ABL persisted after 14:00 as a direct consequence of the presence of
377 high cloudiness (see Figure 6c) diminishing the incoming short-wave radiation to the surface. The
378 RWP shows relatively strong north-easterly winds ($> 6ms^{-1}$) near the surface and up to approx-
379 imately 400 m throughout the morning. The wind profile shows a reduction of the wind speed
380 during the morning, above 400 m and up to the average depth of the valley (1000-1100 m) where
381 the speed is higher due to the trade winds. In this case, the vertical exchanges within the valley
382 atmosphere are restricted by the fact that no large eddies are being formed. In consequence, under
383 these conditions, pollutants do not mix efficiently, as can be observed both in the relatively high
384 ceilometer backscattering intensity and in the PM2.5 hourly concentration record (Figures 6c and
385 d). On October 12, 2018, the atmospheric environment was diametrically opposite. Clear skies
386 allowed for a swift transition from stable to neutral conditions, with an efficient ABL growth, and
387 low backscattering intensities and PM2.5 concentration. An important feature is that winds within
388 the ABL are considerably weaker in the morning time on October 12 than during the same period
389 on February 22. A similar finding is reported in Halios and Barlow (2018), with a negative rela-
390 tionship between the growth rate of the mixing layer and the wind speed. More important than
391 the magnitude of the wind speed, wind shear at the top of the ABL is higher on October 12 than

392 on February 22. The latter could imply a larger ABL growth rate due to increased mechanical
393 turbulence by shear production, leading to entrainment.

394 *b. Heating efficiency*

395 The previous results show a non-linear relationship between Q_{req} and Q_{prov} , and reflect an in-
396 crease in Q_{prov} spread with the magnitude of Q_{req} , with implications for the breakup time. In a
397 closed system, most sensible heating would be used to raise the lower troposphere temperature,
398 expanding the ABL, and none of the energy would be exported to the free atmosphere. Under
399 these idealized conditions, the relationship between Q_{req} and Q_{prov} would be bijective (one-to-one
400 correspondence). If there are, however, atmospheric conditions that lead to heat being exported
401 out of the valley's atmosphere, the heating efficiency would be diminished and most likely vari-
402 able (e.g. Leukauf et al. 2017). Figure 7a shows, for different Q_{req} intervals, the 10th, 50th, and
403 90th Q_{prov} percentiles, with their corresponding regression functions. For the 10th and 50th per-
404 centiles, and up to approximately the 70th (not shown), Q_{req} and Q_{prov} follow a linear relationship.
405 The latter suggests that, in 70% of the cases, the energy provided is mostly used to warm up the
406 atmosphere within the valley: There is a linear relationship between Q_{req} and Q_{prov} for all the
407 percentiles explored up to the 70th, and the changes in the slopes among different percentiles are
408 not considerable. In contrast, from the 75th percentile onwards, Q_{prov} increases exponentially with
409 Q_{req} .

410 Consequently, there appears to be a variable heating efficiency rate that is more evident for larger
411 values of Q_{req} , with direct effects on the breakup time: observations suggest that, for all days with
412 very low heating efficiency (large values of Q_{req} and Q_{prov} above the 70th percentile), the breakup
413 occurs later than 13:00 LT. Therefore, it is essential to evaluate which mechanisms or atmospheric
414 patterns are associated with low heating efficiency. Previous work (Angevine et al. 2001; Leukauf

415 et al. 2016; Nadeau et al. 2018) link this possible leakage of energy with the valley circulation and
416 the wind speed. The breaking times vary from 07:00 LT to approximately 16:00 LT, depending
417 on the heat efficiency rate. When the heat efficiency is low (the upper part of the scatterplot), the
418 nighttime inversion breaks late in the afternoon (after 14:00 LT), being unfavorable for pollutant
419 dispersion as shown in the previous subsection.

420 ROLE OF LOCAL WIND SHEAR

421 To evaluate the potential influence of wind speed and vertical shear on the heating efficiency
422 during the morning transition, we followed a composite analysis of these variables during three
423 different subsets of dates. The three subsets of dates with contrasting heating efficiency corre-
424 spond, in all cases, to values above the mean Q_{req} , and (i) Q_{prov} values below the 33th percentile,
425 (ii) between the 33 and the 66th percentile, and (iii) above the 66th percentile (see Figure 7b).
426 Figure 8a to c, and d to f, show the wind speed and the vertical wind shear, respectively, for the
427 three subsets of dates. Area I corresponds to the lowest heating efficiency among the three subsets.
428 Conversely, Area III corresponds to the highest heating efficiency (less energy provided to reach
429 neutralization for a similar amount of Q_{req}). The evidence indicates that for lower wind speeds
430 near the surface between 6:00 and 10:00 am LT, and more notably, for higher vertical wind shear,
431 the erosion of the SBL occurs earlier and with less energy provided to the atmosphere in the form
432 of surface H .

433 The observed enhanced shear corresponding to dates in the Area III set compared to the other
434 sets is elevated, being maximum across the top of the SBL at the entrainment zone, rather than
435 near the surface. Even under low wind speeds and with shear differences less than 1 ms^{-1} , the
436 observational evidence suggests that shear production of TKE cannot be neglected. From the

437 observations, the elevated shear appears to play an important role in enhancing the erosion of the
438 SBL, likely by generating TKE.

439 Different authors have studied the role of elevated shear in the evolution of the CBL, most using
440 a modeling approach and some using observations in flat terrains (e.g. Angevine 1999; Fedorovich
441 et al. 2001; Conzemius and Fedorovich 2006; Fedorovich and Conzemius 2008; Halios and Barlow
442 2018). However, there is no consensus on whether a mean elevated shear enhances or suppresses
443 entrainment. Conzemius and Fedorovich (2006) state that the boundary layer begins to grow
444 due to increasing surface H and entrainment, with air from the free atmosphere being engulfed
445 by convective thermals and becoming part of the boundary layer, a process that is modified by
446 the presence of an elevated shear (Fedorovich et al. 2001). Compared to the effect of surface
447 shear, the influence of elevated shear across the inversion on turbulence in the SBL and CBL is
448 much less studied. It is clear that in addition to the often dominant buoyancy forcing, the CBL
449 development is modulated by wind shear, which modifies considerably the internal structure of the
450 lower troposphere. Therefore from this point of view, the timing of the breakup is modulated by
451 the evolution of the surface H and the amount of mechanical turbulence due to wind shear.

452 Very few studies have explored the role of the elevated shear in a setting characterized by com-
453 plex terrain and urbanization. The observational evidence presented here is not in agreement with
454 the results presented in the theoretical work by Hunt and Durbin (1999). They found that the ele-
455 vated shear prevented the entrainment process and the generation of TKE by deforming thermals
456 so that they do not penetrate as effectively into the inversion, interfering with the entrainment, a
457 process referred to as shear sheltering. However, in their work, they did not consider the potential
458 effects of density stratification and the complex terrain setting. In their results, thermals do not
459 overshoot their equilibrium level, and the CBL growth is slower. Fedorovich et al. (2001) and
460 Conzemius and Fedorovich (2006) explore the directional effect of the elevated wind shear on the

461 turbulent exchange across the capping inversion in Large Eddy Simulation (LES) experiments. In
462 the numerical experiments, when the mixed-layer air has a higher momentum than the air above
463 the inversion (negative elevated shear), CBL growth is enhanced contrary to the sheltering pro-
464 cess described by Hunt and Durbin (1999). In contrast, in cases of positive shear, CBL growth is
465 diminished. In contradiction to the mentioned modeling results, the evidence in Figure 8c and d
466 shows a case where higher positive elevated shear (mixed-layer air has less momentum than the
467 air above the inversion) leads to faster erosion of the SBL compared to when the positive elevated
468 shear is weaker. The evidence suggests that the elevated shear does result in thermal damping
469 at the inversion layer inhibiting the entrainment; conversely, it likely favors TKE generation and
470 intensification of vertical transport of air from the mixed layer into the above-inversion region.
471 The coincidence in the modeling studies and the results in Figure 8 lies in that the elevated shear
472 appears to be much more important than the surface shear in the erosion of the SBL.

473 In addition to the vertical wind speed and wind shear, the magnitude of the upslope-downslope
474 winds for the sets of days in Areas I, II, and III was also contrasted. The results do not show
475 considerable and consistent differences among the three sets of dates.

476 ROLE OF SYNOPTIC CONDITIONS

477 The role of synoptic conditions on the ABL evolution over the Aburrá Valley is assessed con-
478 sidering the velocity potential, stream function, and OLR anomalies. The anomalies are computed
479 as the difference between the daily average of the variable of interest and the monthly average
480 of the same variable. The velocity potential, OLR, and stream function average anomalies for the
481 set of days corresponding to Areas I, II, and III in Figure 7b are shown in Figures 9, 10, and 11,
482 respectively. Together, these variables represent the overall regional-scale convective activity.

483 The results suggest that, overall, for a similar Q_{req} values, the erosion of the SBL occurs not
484 only faster but also with lower values of Q_{prov} (higher heating efficiency) for cases when the
485 deep convection is inhibited regionally. Conversely, the SBL erosion is delayed, often until the
486 afternoon, in scenarios when the regional deep convective activity is enhanced. In the latter case,
487 the Q_{prov} values are three-four times larger than in the former conditions.

Figure 9a, b and c show, for Area I, positive velocity potential anomalies at 700 hPa over north-
ern South America, weak anomalies at 500 hPa, and negative anomalies at 200 hPa, respectively.
Such configuration indicates an enhancement of the deep convective activity in the region. In
tropical South America, an enhanced deep convective activity often leads cloud formation. Figure

Fig: OLR_A anomalies show negative OLR anomalies associated with Area I, which agrees with the observed velocity

488 The deep convective activity and OLR contrast for the dominant regional conditions in Area
489 I vs. Area III suggests that the radiative forcing associated with deep convective clouds plays
490 a dominant role in modulating the SBL erosion than the dynamical effect of the regional-scale
491 convection. The average radiation between 06:00 and 12:00 LT for Area I is 401 Wm^{-2} and
492 for Area III is 472 Wm^{-2} . It is expected that the 71 Wm^{-2} difference in radiation reaching the
493 surface lead to a belated erosion of the SBL. However, the radiation difference in itself does not
494 directly explain the larger Q_{prov} required in these cases, considering that Q_{req} at 06:00 LT is similar.
495 Nevertheless, it does suggest that the extended SBL erosion period leads to inefficient heating of
496 the ABL. It is likely that with longer erosion times, different processes such as heat export outside
497 of the valley linked to upslope flow (e.g. Noppel and Fiedler 2002) lead to lower heating efficiency.

498 *c. Variability of the breakup time and consequences for air quality*

499 The absence of a marked top of the atmosphere radiation and surface air temperature seasonality
500 in low-latitudes does not imply an insignificant valley-scale response to the annual climatology. In

501 fact, the annual cycle in the tropics does impose variable large scale forcing, modulating the ABL
502 variability. In the tropics, the seasonality of the Intertropical Convergence Zone (ITCZ) modulates
503 local cloudiness, precipitation, and surface incident radiation, altering the characteristics of the
504 nocturnal inversion and the erosion of the SBL.

505 Figure 12a shows a time-dependent clustering in the Q_{req} vs. Q_{prov} scatterplot around two
506 different seasons (Feb-Jun and Jul-Nov). The latter suggests that the strength of the nighttime
507 inversion, hence the energy required to erode the SBL, changes significantly throughout the year.
508 Consequently, the inversion-breaking time also varies (see Figure 12b). Similar results are seen for
509 the WRF forecasts runs (Figure 5). From February to mid-June, the proxy for Q_{req} is, on average,
510 twice as large as that of the July-November period, and the inversion breakup occurs later in the
511 day, in some cases even after 14:00 LT.

512 The timing of the atmospheric neutralization plays an important role in modulating air pollutant
513 concentration. The PDFs for the daily average PM_{2.5} concentration, conditioned to breaking times
514 occurring within four different hours during the day, suggest that, as the inversion breakup time
515 occurs later in the day, the likelihood of higher PM_{2.5} concentrations increases. The concentration
516 of aerosols in the valley's atmosphere is mainly influenced by the anthropogenic emissions at
517 surface level and by the vertical dispersion of pollutants after the inversion breakup (e.g. Herrera-
518 Mejía and Hoyos 2019). In cases of a late breakup, emissions accumulate within the SBL until
519 thermal turbulence is activated, after which pollutants are lifted out of the valley. For the specific
520 case of the Aburrá Valley, where atmospheric pollutant dispersion out of the valley is almost
521 entirely thermodynamically driven, the magnitude of the turbulent exchange must be large enough
522 for the ABL to reach the mountain peaks (1200 m.a.s.l.), where the pollutants are advected away
523 far from the valley.

524 4. Conclusions

525 The variability and implications of the timing of the stable boundary layer breakup have been
526 examined for a narrow valley located in the tropical Andes Cordillera using a combination of
527 in-situ turbulent scale observations and remotely sensed macroscopic features of the local atmo-
528 sphere. Given the topographic features of the region, it is imperative to understanding when and
529 under which conditions the nocturnal inversion breaks up because it corresponds to the time when
530 the exchanges between the surface and the free atmosphere intensify and reach their maximum,
531 resulting in a more efficient pollutant vertical transport.

532 The assessment is based on an observational diagnostic framework developed to study the
533 breakup time using proxies for the energy required to erode the atmospheric inversion (Q_{req})
534 and the amount of energy provided to the atmosphere via sensible heating (Q_{prov}), combining
535 high frequency measurements near the surface with macroscopic observations of the atmosphere
536 in the vertical profile. In this framework, the inversion breakup occurs when Q_{prov} via surface
537 sensible heat flux (H) is equal to Q_{req} . Different thermodynamic indices were considered as
538 proxies for Q_{req} , including changes of virtual potential temperature in the vertical at the lower-
539 troposphere ($\Delta_z \theta_v$) and $CINE$. The inversion breakup assessment using the proposed framework
540 involves challenges associated with the spatial representativeness of each of the variables consid-
541 ered in the study. However, the high covariability between the hourly H , $\Delta_{200} \theta_v$, $\Delta_{800} \theta_v$, $\Delta_{Total} \theta_v$
542 ($\theta_v(1200) - \theta_v(50)$), $\Delta_{Sup} \theta_v$ ($\theta_v(800) - \theta_v(50)$), $CINE$, $CINE_{1200}$, and $CINE_{1500}$ indicates a co-
543 herent response among different atmospheric scales considered, hence serving as validation of the
544 proposed methodology, regardless the limitations.

545 The relationship between Q_{req} and Q_{prov} allows concluding that Q_{prov} indeed is higher when the
546 magnitude of Q_{req} is large, regardless of the proxy used. However, the observations indicate that

547 the Q_{req} - Q_{prov} relationship is by no means simple. The evidence suggests the existence of non-
548 constant heating efficiency for large values of Q_{req} , similar to the findings of Leukauf et al. (2017).
549 In approximately 70% of the cases, the energy provided is mostly used to warm up the valley's
550 atmosphere. In contrast, for approximately 25% of the cases, Q_{prov} increases exponentially with
551 Q_{req} .

552 The vertical wind shear appears to be an important factor modulating the breakup time, hence
553 the apparent heating efficiency. Higher vertical wind shear is linked to the earlier erosion of the
554 SBL, with less energy provided to the atmosphere in the form of surface H . Moreover, the higher
555 vertical wind shear does not occur near the surface. Instead, it is elevated, and it is maximum across
556 the top of the SBL at the entrainment zone, suggesting that shear production of TKE cannot be
557 neglected. The elevated shear, regardless of directional considerations, appears to play an essential
558 role in enhancing the erosion of the SBL, likely by generating TKE. The evidence suggests that
559 the timing of the breakup depends not only on the surface H but also on the amount of mechanical
560 turbulence due to the elevated wind shear. The observational evidence presented here is important
561 since there is no consensus on whether a mean elevated shear enhances or suppresses entrainment.

562 The synoptic conditions also play a role in the ABL evolution over the Aburrá Valley and
563 breakup time. Velocity potential and OLR anomalies indicate that the erosion of the SBL oc-
564 curs faster and with lower values of Q_{prov} (higher heating efficiency) when the deep convection is
565 inhibited regionally. Conversely, the SBL erosion is delayed in scenarios when the regional deep
566 convective activity is enhanced. The contrast in deep convective activity and OLR linked to vari-
567 able heating efficiency suggests that the radiative forcing associated with deep convective clouds
568 plays a dominant role in modulating the SBL erosion than the dynamical effect of the regional-
569 scale convection. The difference in average radiation between 06:00 and 12:00 LT between cases
570 with high and low heating efficiency was found to be approximately 70 Wm^{-2} . This difference

571 is considerable and translates into considerably different breakup time, and with longer erosion
572 times, various processes such as heat export outside of the valley through upslope flow reduce the
573 heating efficiency.

574 The results suggest a large breakup time variability as a function of heating efficiency. In addi-
575 tion, the breakup time variability has been shown to have a profound impact on local air quality
576 within the valley. The evidence indicates that, for later breakup times, the likelihood of higher
577 PM_{2.5} concentrations increases considerably. In cases of a late breakup in complex terrains, an-
578 thropogenic emissions accumulate within the SBL until thermal turbulence is activated.

579 *Acknowledgments.* PENDING: OEAD and U Innsbruck, SIATA-AMVA, UNAL..., data... ORL,
580 ERA... Interpolated OLR data provided by the NOAA/OAR/ESRL PSL, Boulder, Colorado, USA,
581 from their Web site at
582 , all part of the meteorological and air-quality monitoring network of the Medellín and Aburrá
583 Valley Early Warning System (SIATA, www.siata.gov.co)

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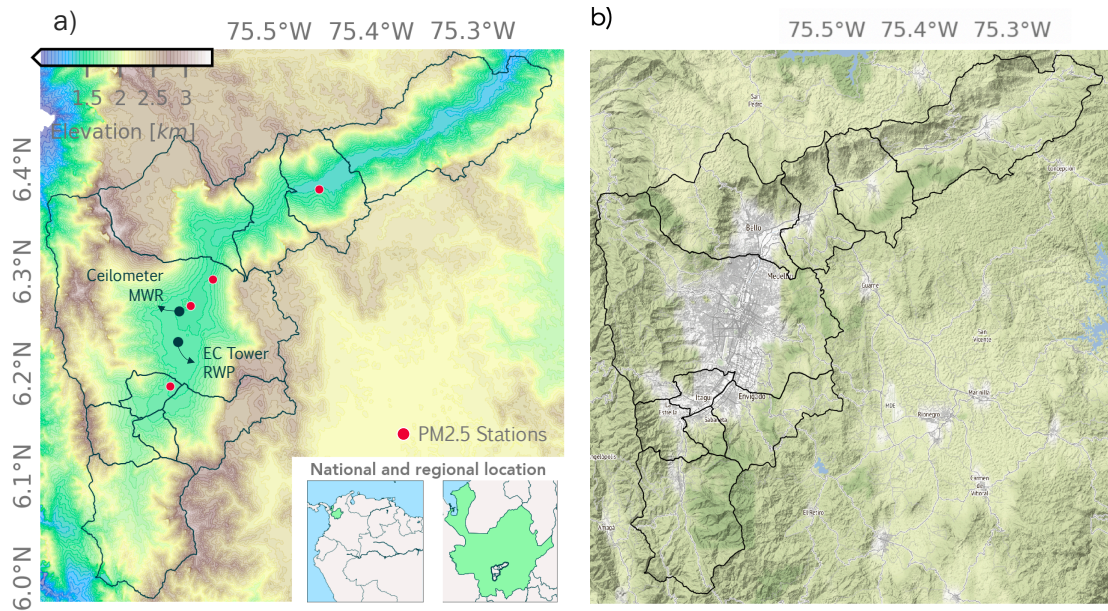
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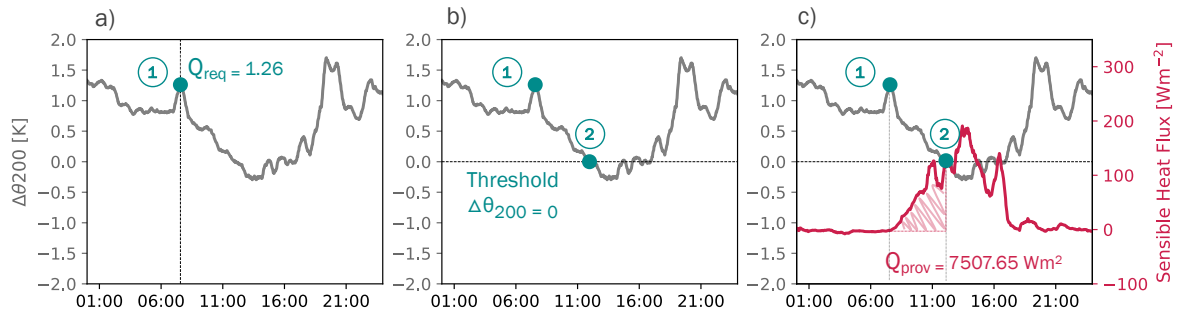
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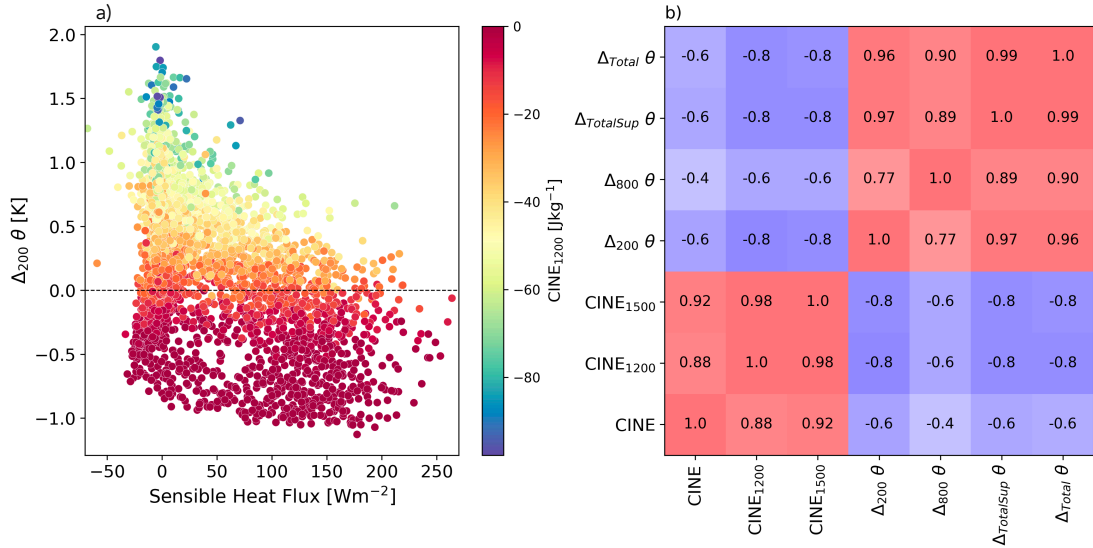
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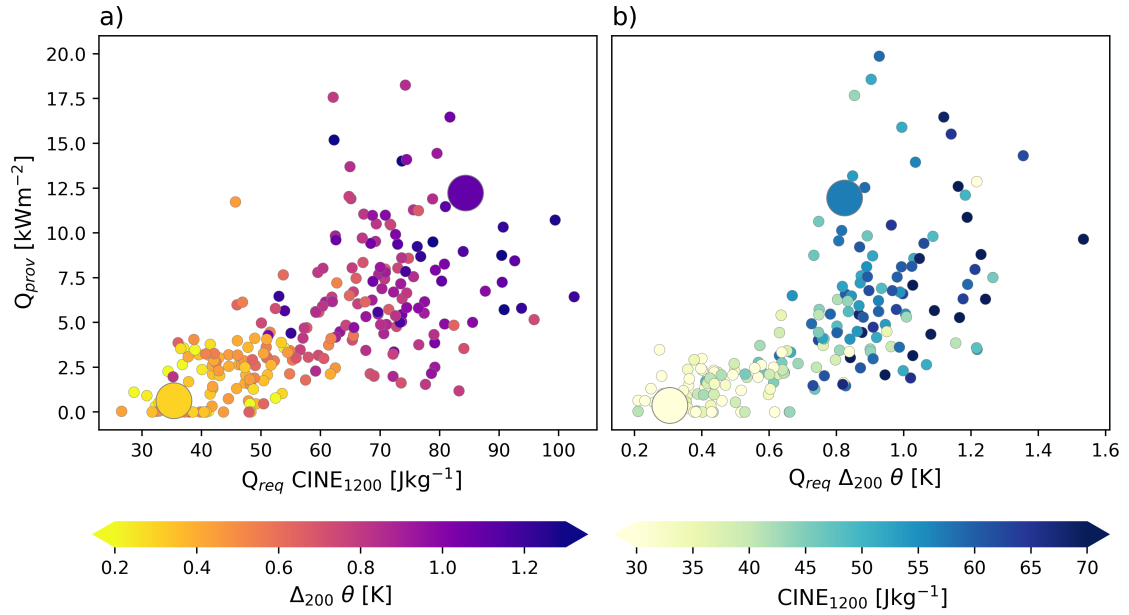
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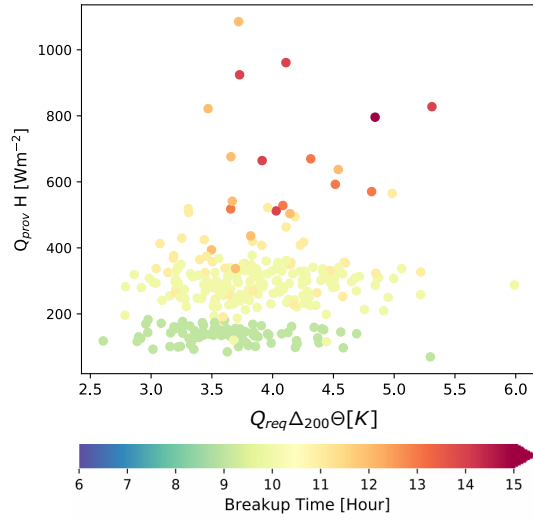
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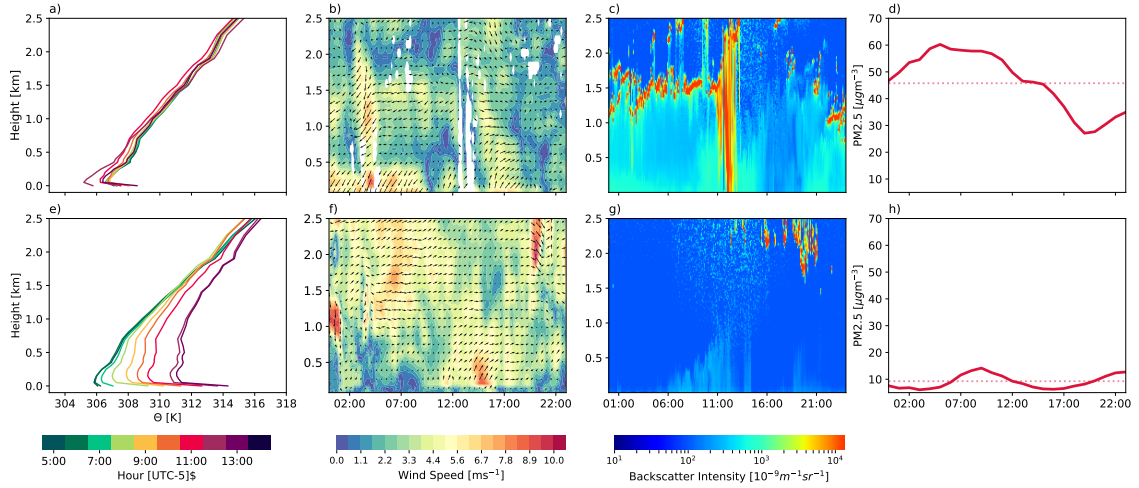
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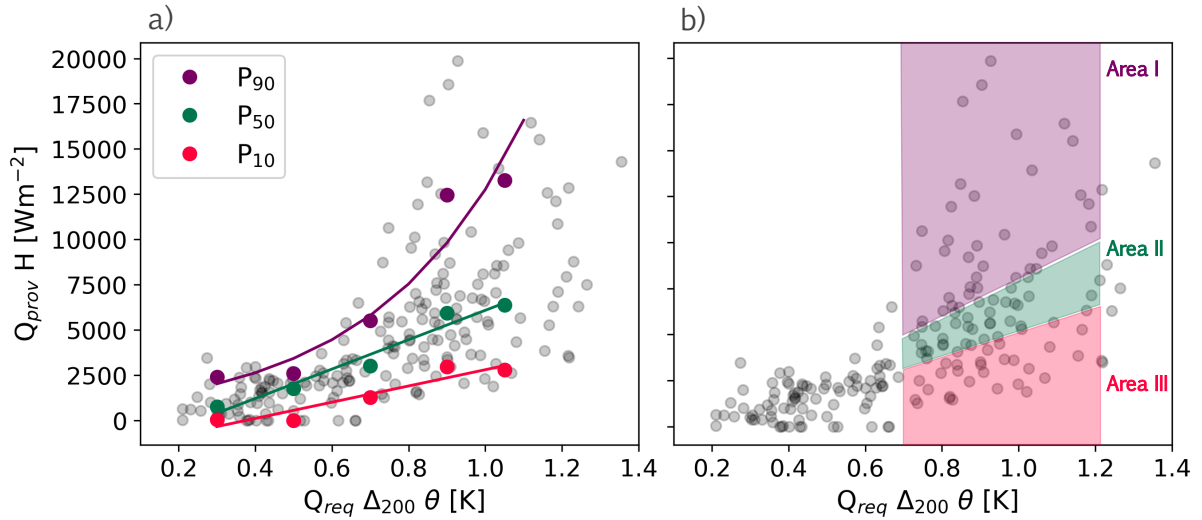
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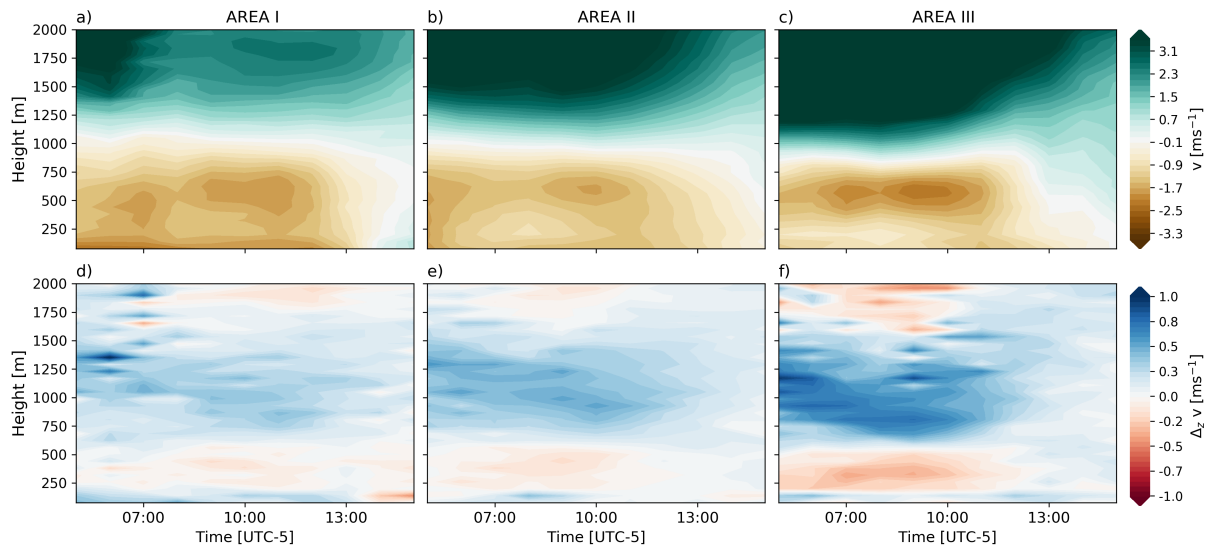
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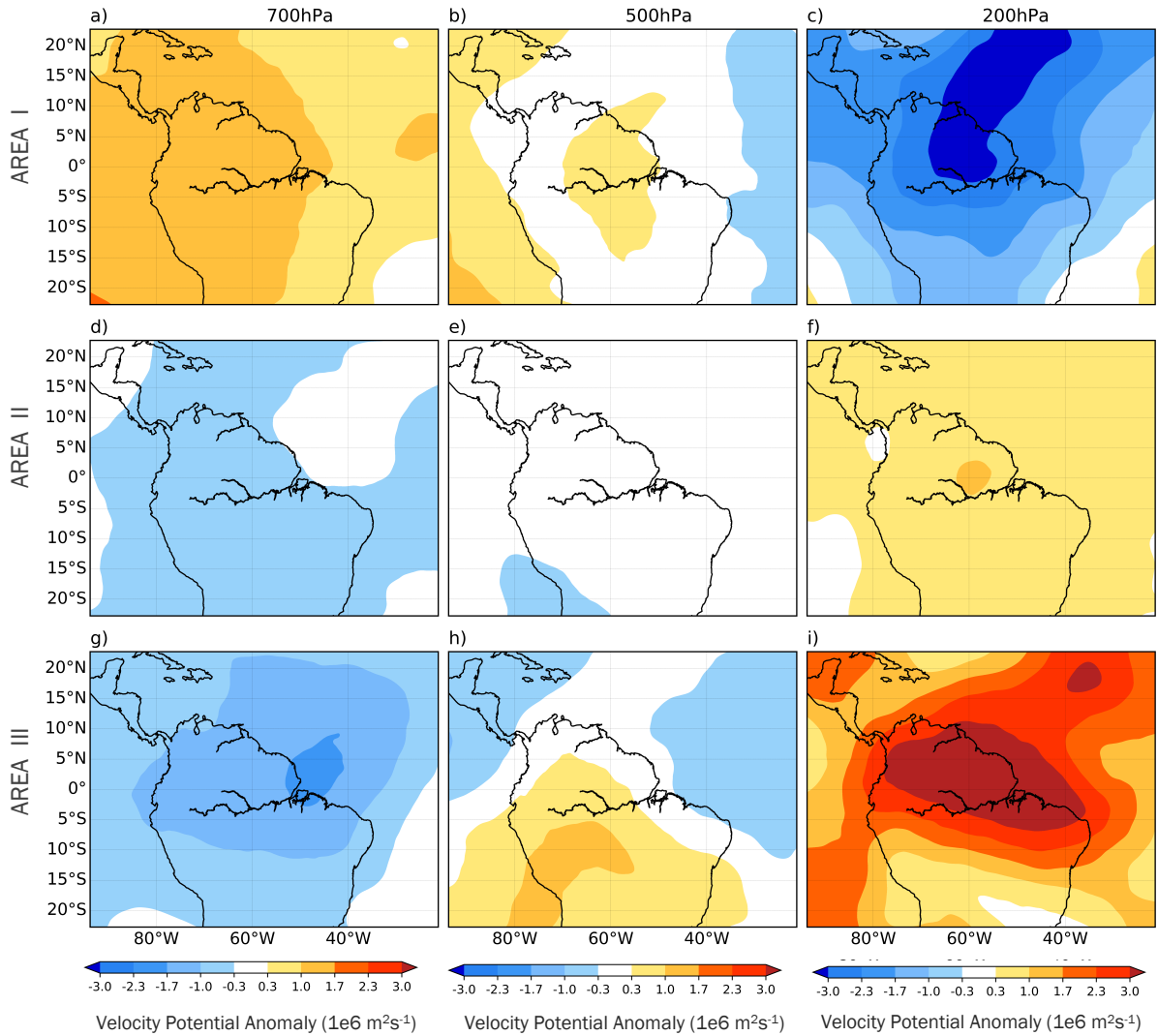
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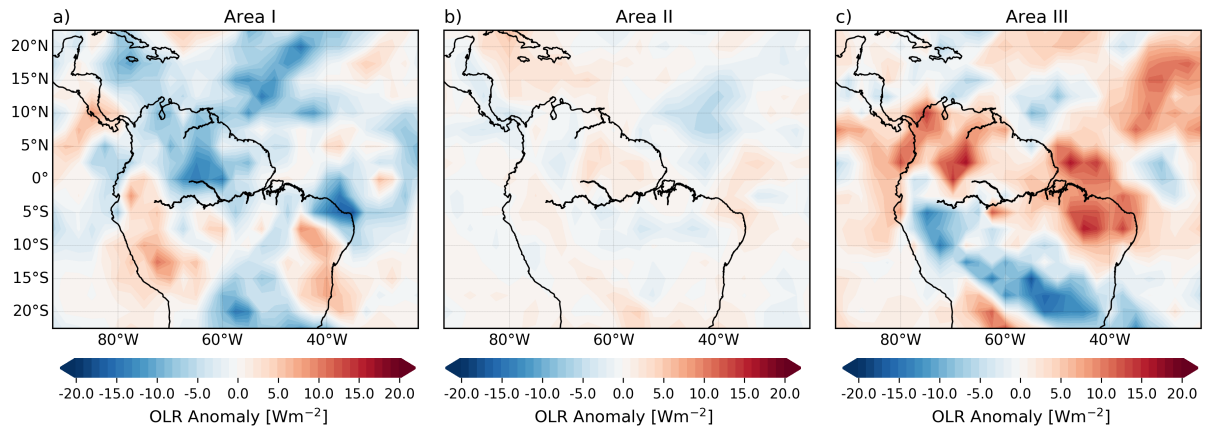


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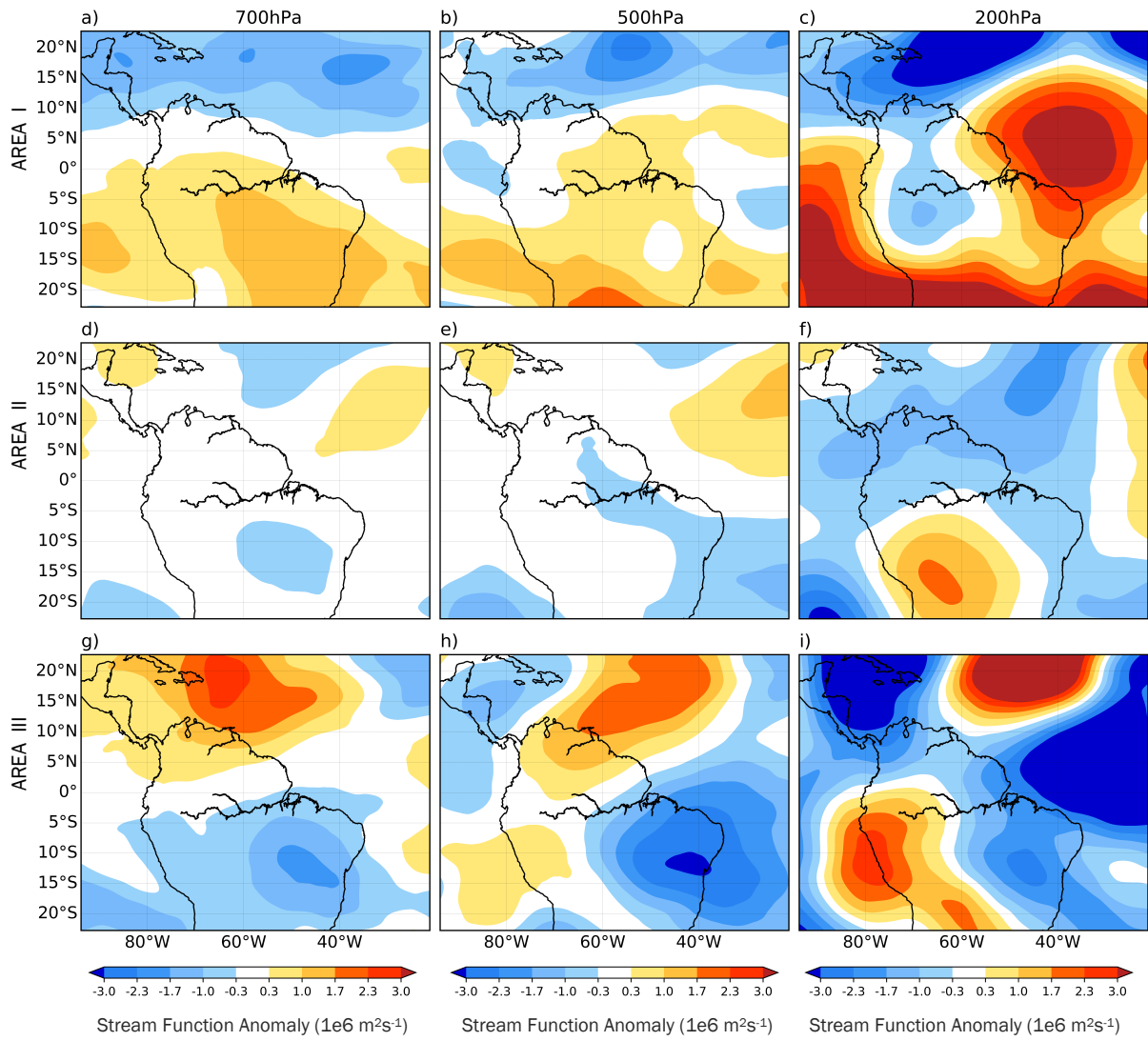
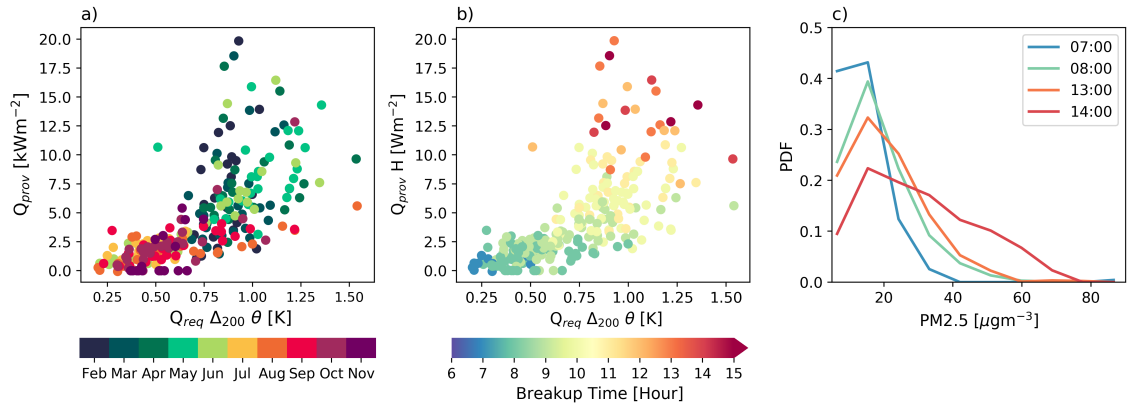


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