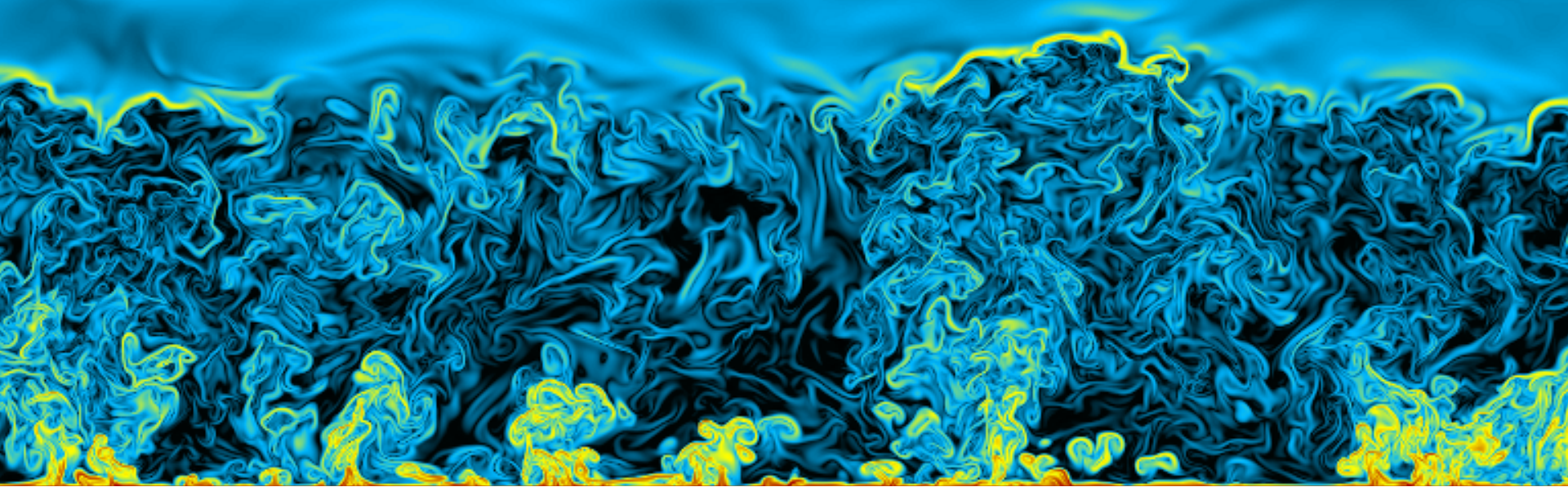


BOUNDARY LAYER METEOROLOGY



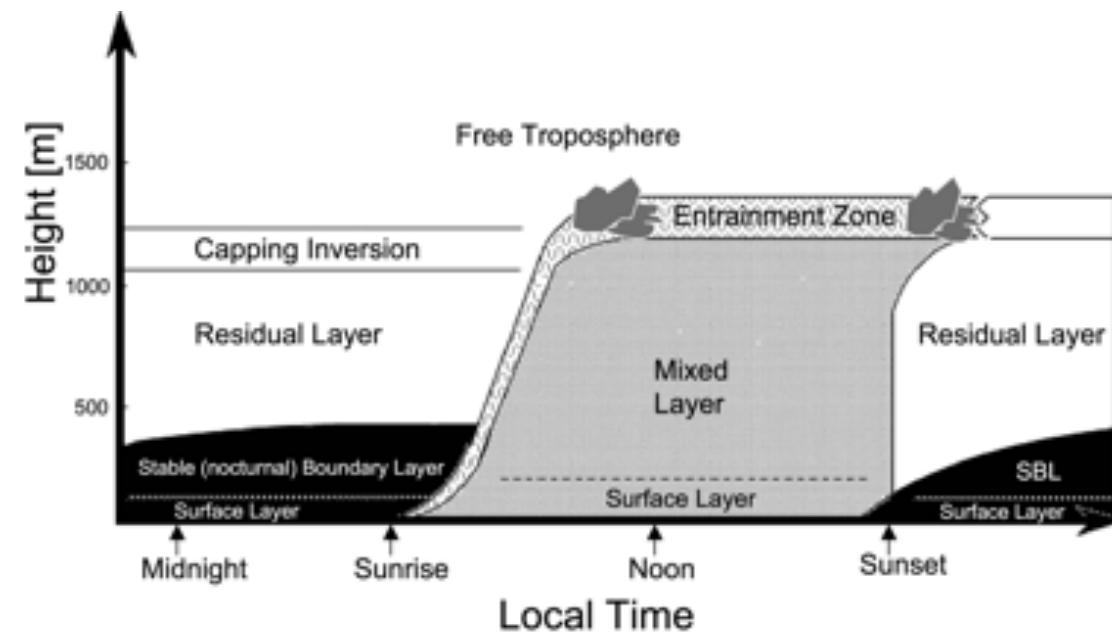
Prof. Ivana Stiperski, Dr. Manuela Lehner
Department of Atmospheric and Cryospheric Sciences

Chapter 4

Similarity Theory

Similarity Theory

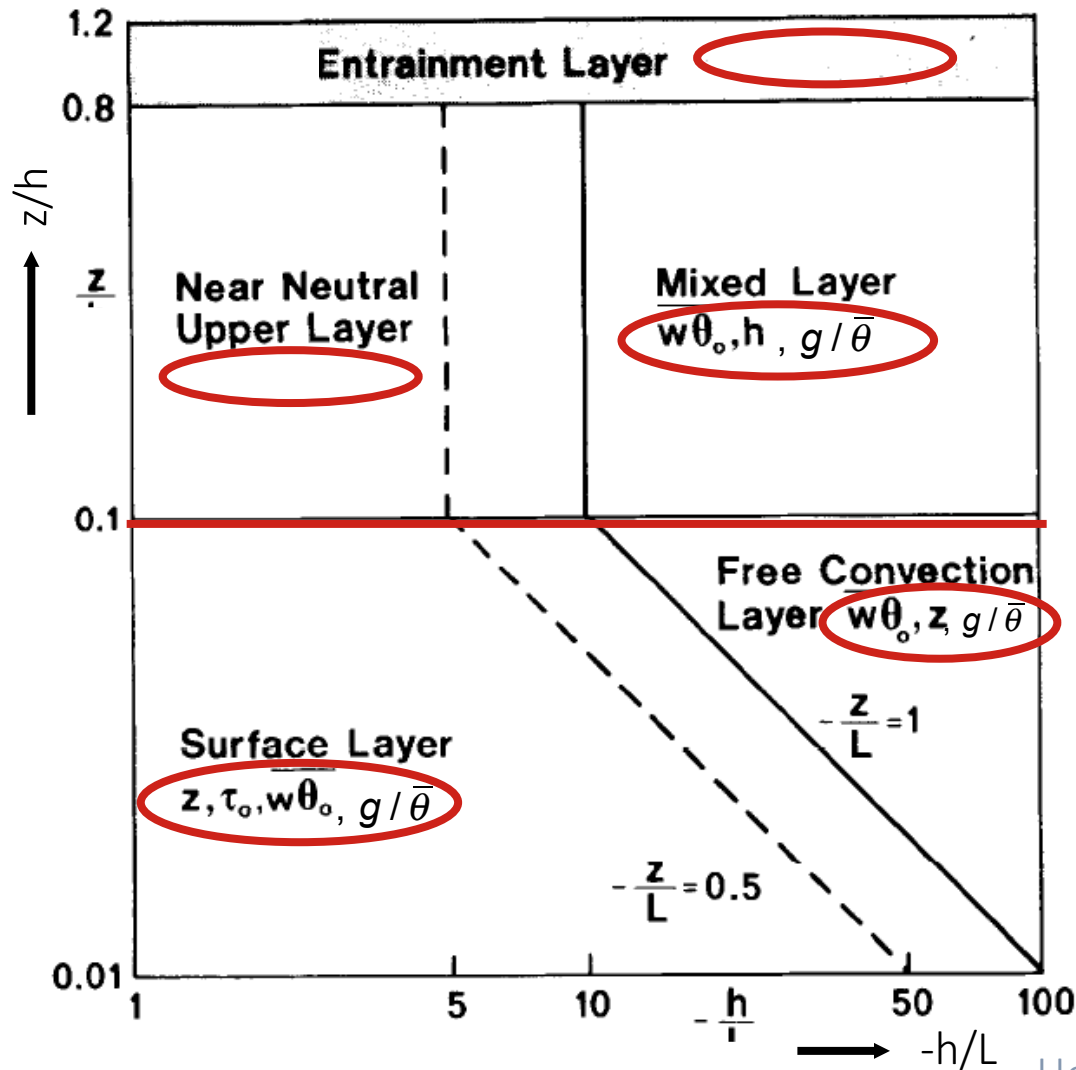
- one example shown: MOST (Surface Layer)
- in the boundary layer: different *Scaling Regimes* exist



- Boundary Layer in 'phase space'
- horizontally homogeneous (presumed)
- non-dimensional height vs. non-dimensional stability

Holtslag and Nieuwstadt (1986)

Scaling regimes: Unstable

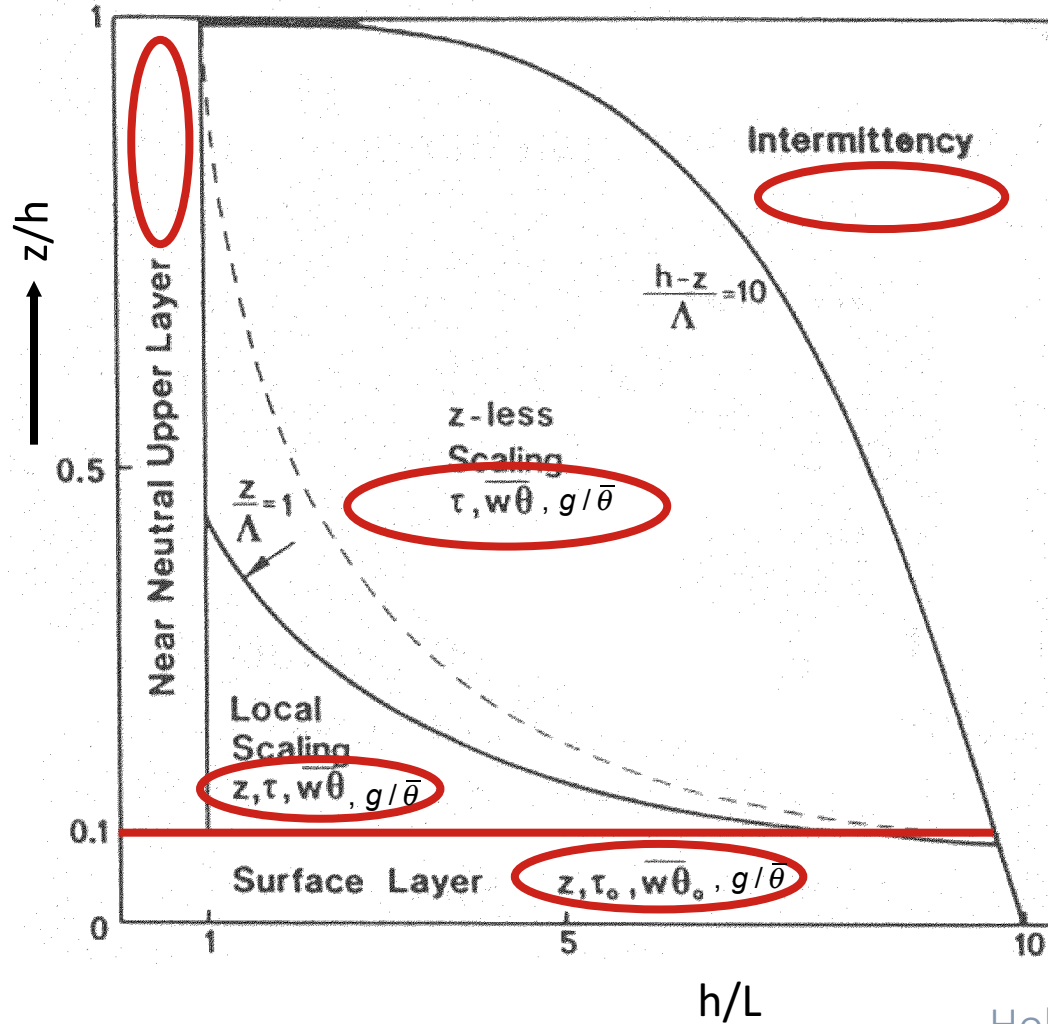


$$h \hat{=} z_i$$

$$L \equiv -\frac{1}{k} \frac{u_*^3}{w'\theta'_0} \left(\frac{g}{\bar{\theta}}\right)^{-1}$$

Holtslag and Nieuwstadt (1986)

Scaling regimes: Stable



$$h \hat{=} z_i$$

$$L \equiv -\frac{1}{k} \frac{u_*^3}{\overline{w'\theta'_0}} \left(\frac{g}{\overline{\theta}}\right)^{-1}$$

Holtslag and Nieuwstadt (1986)

Scaling regimes

- 'always' Surface Layer (MOST) at the bottom
- regimes *with* variables: 'successful'
- regimes *without*: 'not successful'

'successful' :


- T_f = 'forcing time scale'
- T_m = Time scale 'mean profile'

- $T_f \gg T_m$: quasi stationary

quasi stationary \leftrightarrow successful

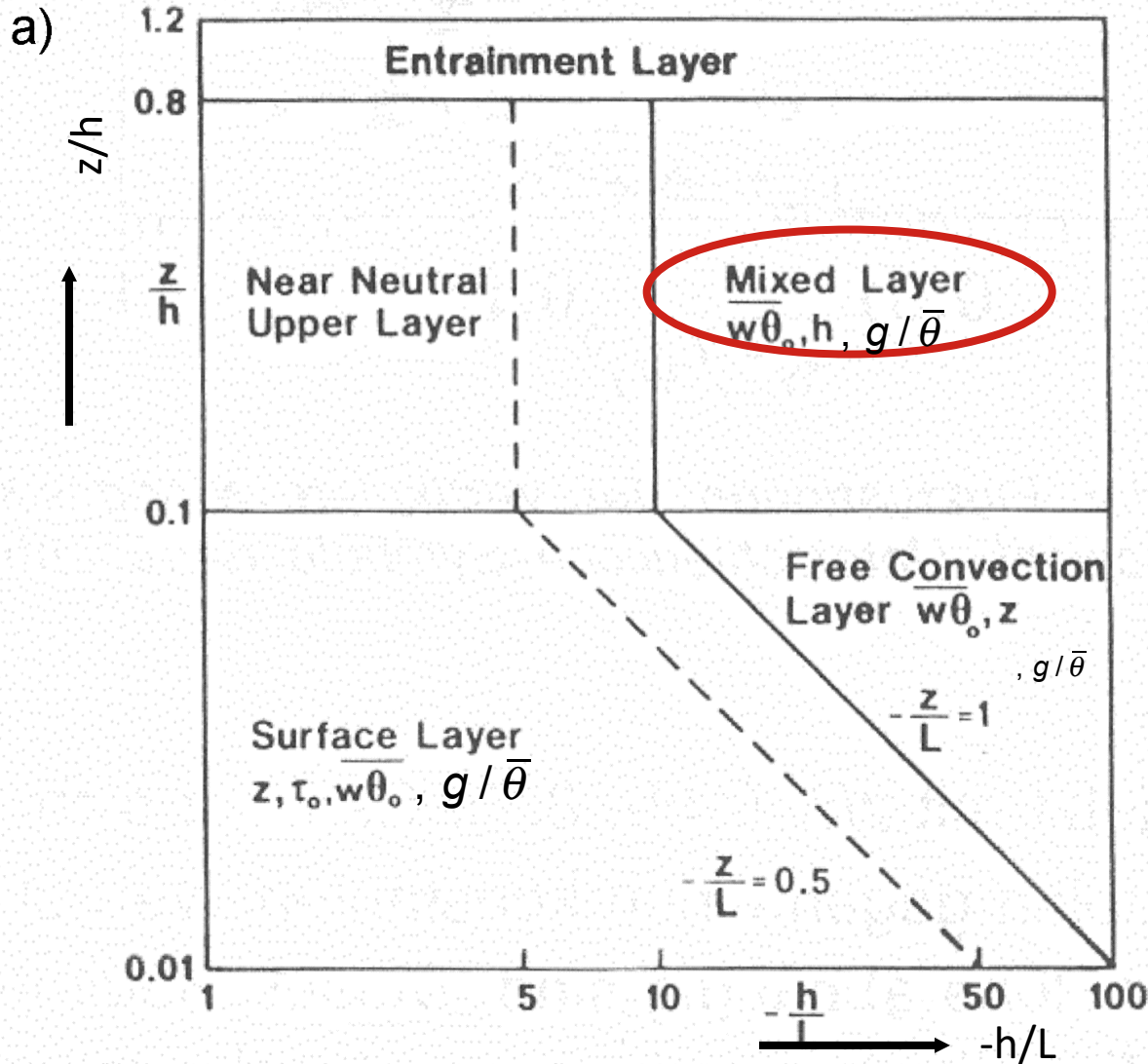
Scaling regimes: The successful Three

The successful three

→ Surface Layer (MOST) 
(including Free Convection limit)

→ Mixed Layer

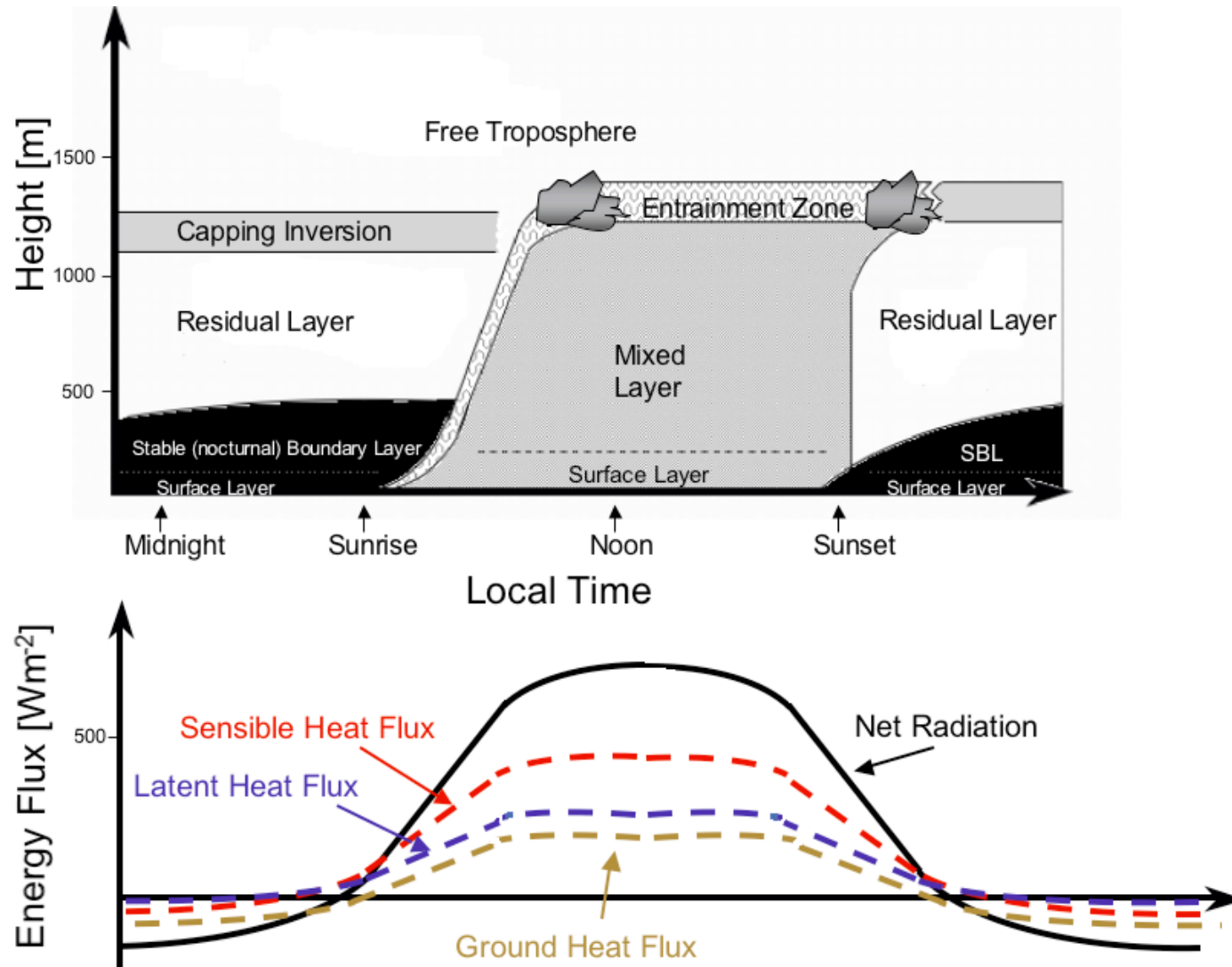
Scaling regimes: Mixed Layer



$$h \hat{=} z_i$$

$$L \equiv -\frac{1}{k} \frac{u_*^3}{\overline{w'\theta'_0}} \left(\frac{g}{\bar{\theta}}\right)^{-1}$$

Scaling regimes: Mixed Layer



Scaling regimes: Mixed Layer

Processes:

$\overline{w'\theta'}$ heating at the ground

z_i entrainment

$g/\bar{\theta}$ buoyancy

z length scale



$N=1$

$$\pi_1 = z / z_i$$

$$\frac{\bar{a}}{a_*} = f_a\left(\frac{z}{z_i}\right)$$

Scaling regimes: Mixed Layer

velocity scale:

→ dimensional analysis:

$$w_* = \left(\frac{g}{\theta} \overline{w' \theta'_{oZ_i}} \right)^{1/3}$$

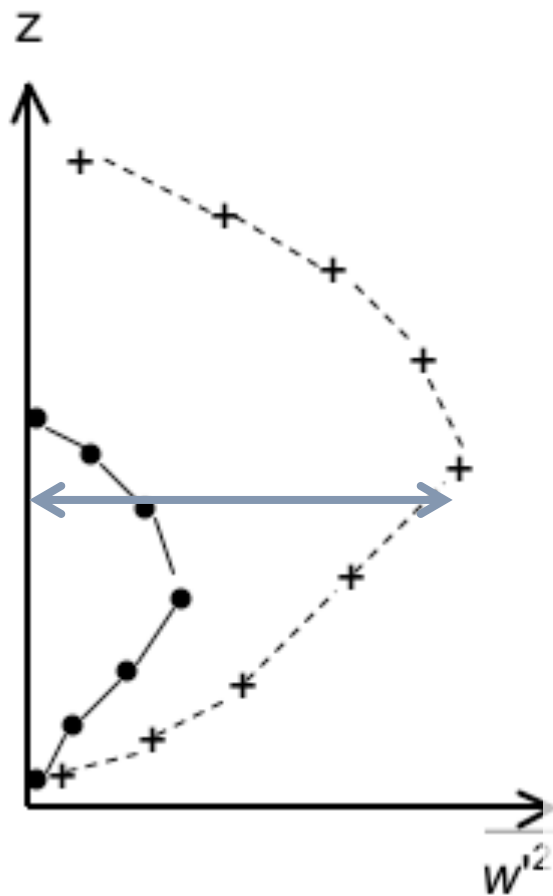
temperature scale:

→ dimensional analysis:

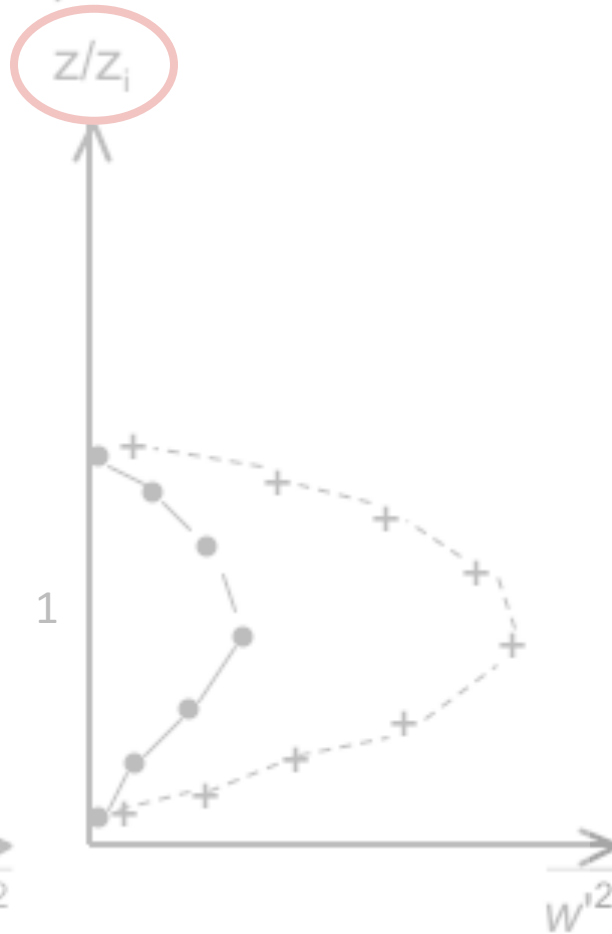
$$\theta_{*CBL} = \overline{w' \theta'_{oZ_i}} / w_*$$

Scaling regimes: Mixed Layer

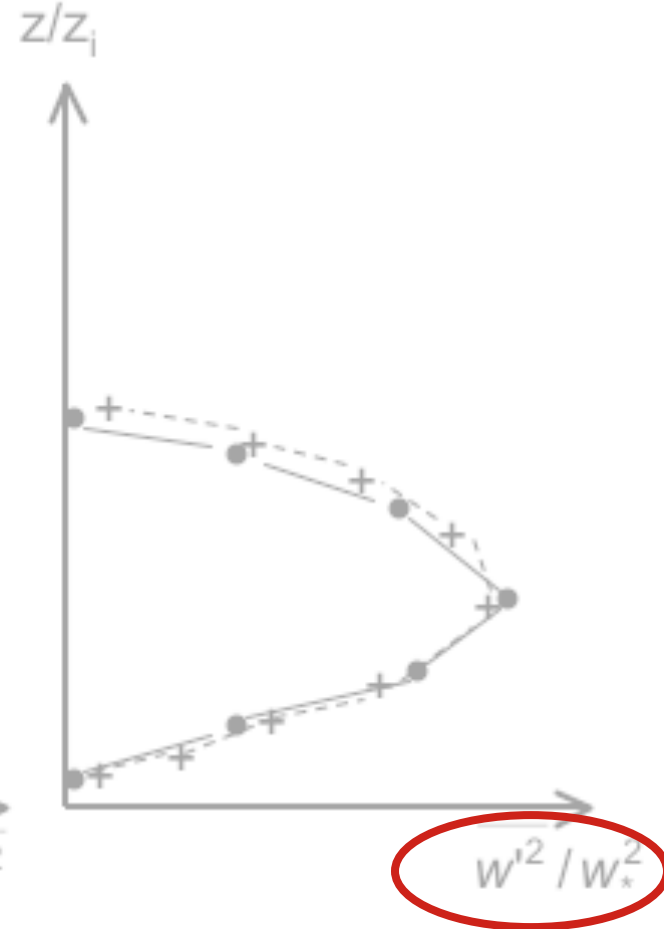
a)



b)



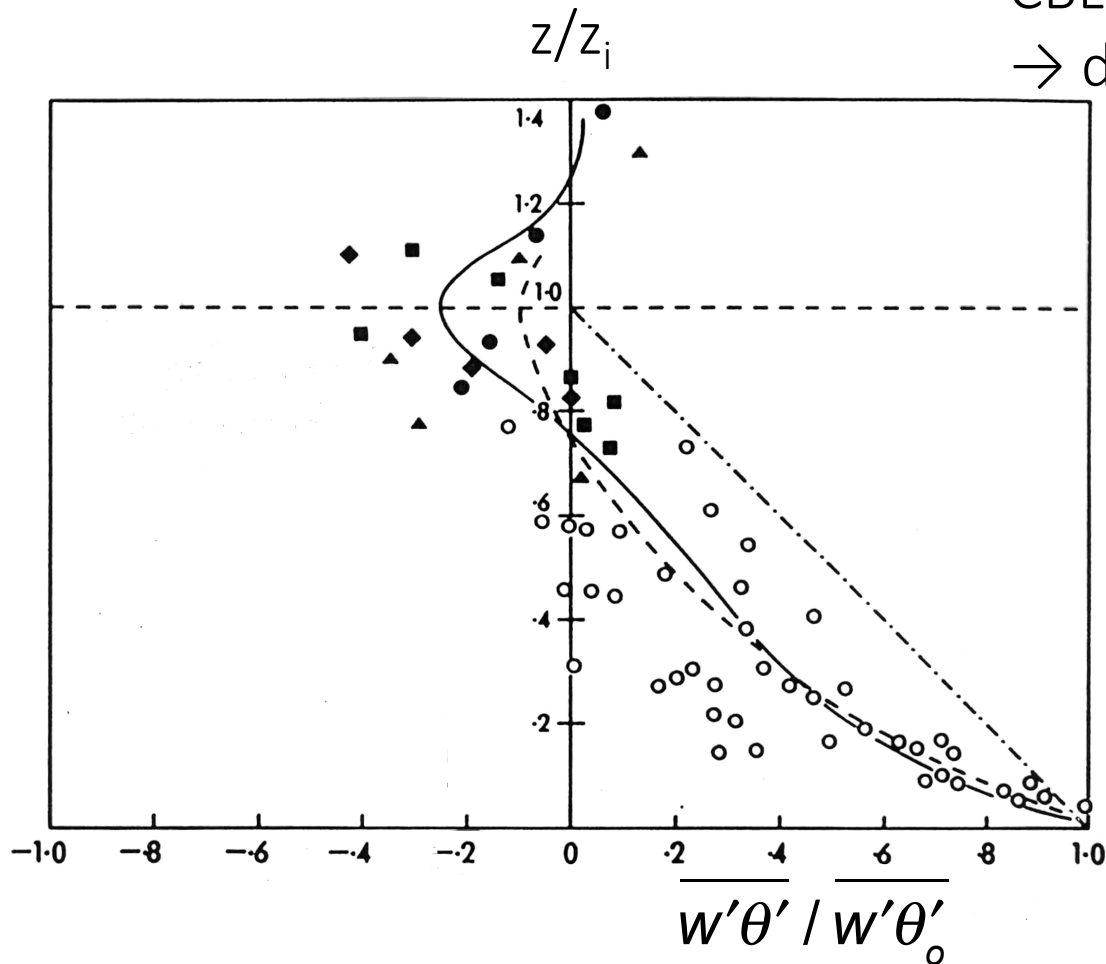
c)



Scaling regimes: Mixed Layer

CBL

→ different measurements



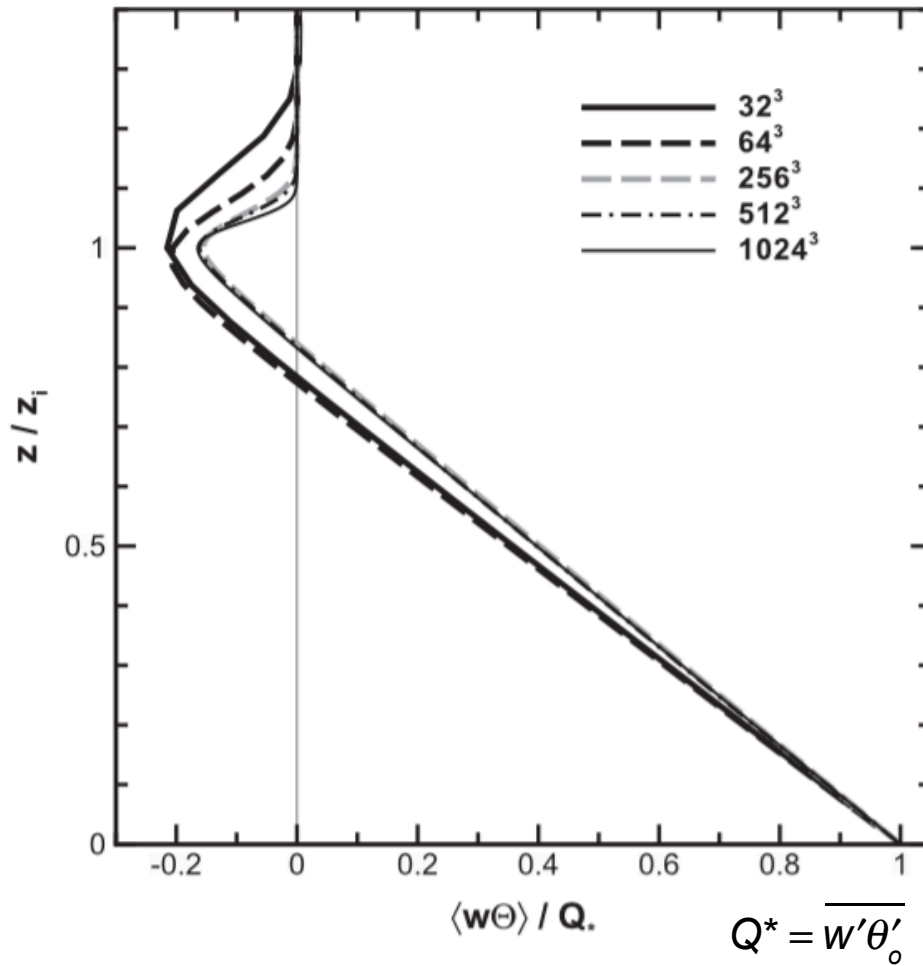
symbols:

different flight campaigns

→ difficult measurements

→ 'large' scatter

Scaling regimes: Mixed Layer



LES of CBL

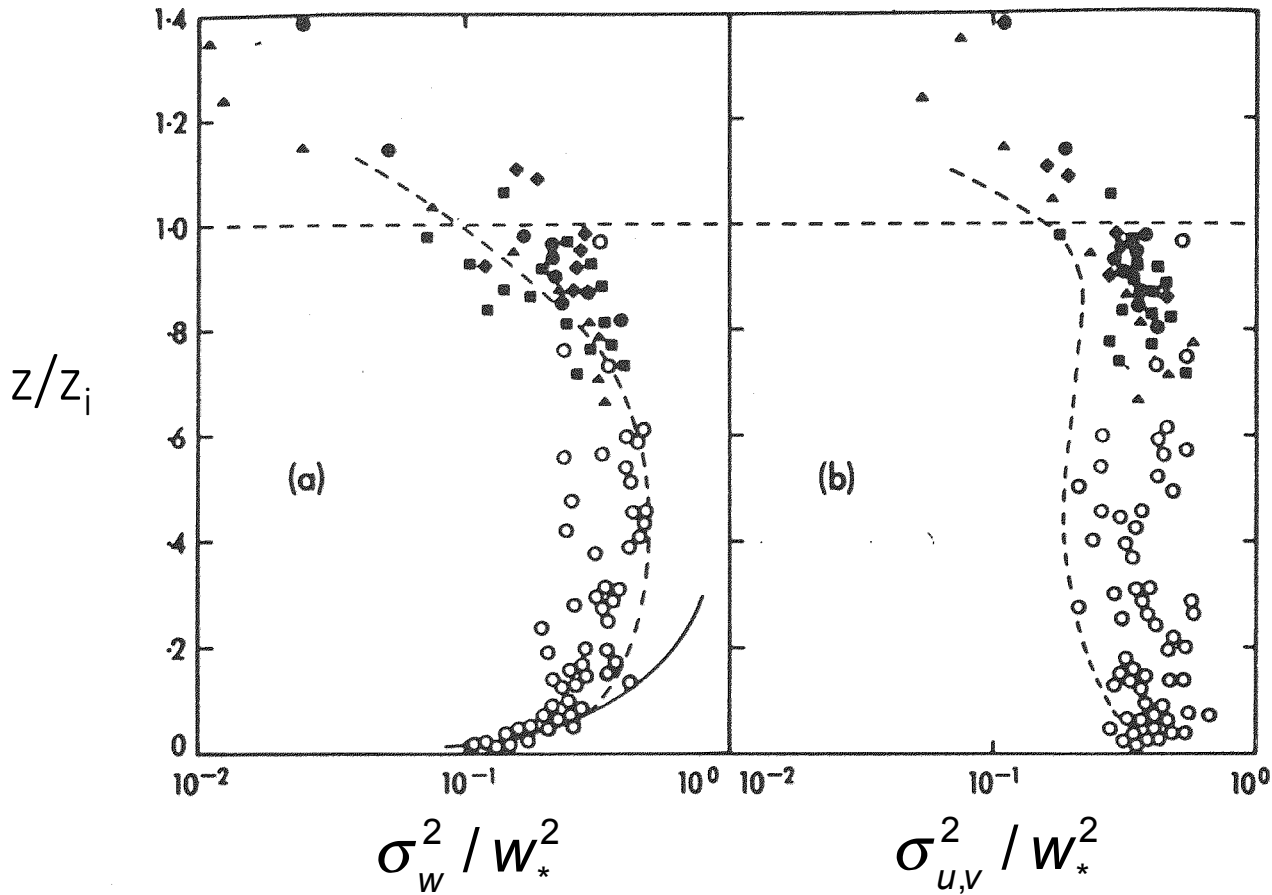
→ different resolutions

→ 'no' scatter under
ideal conditions

Sullivan and Patton (2011)

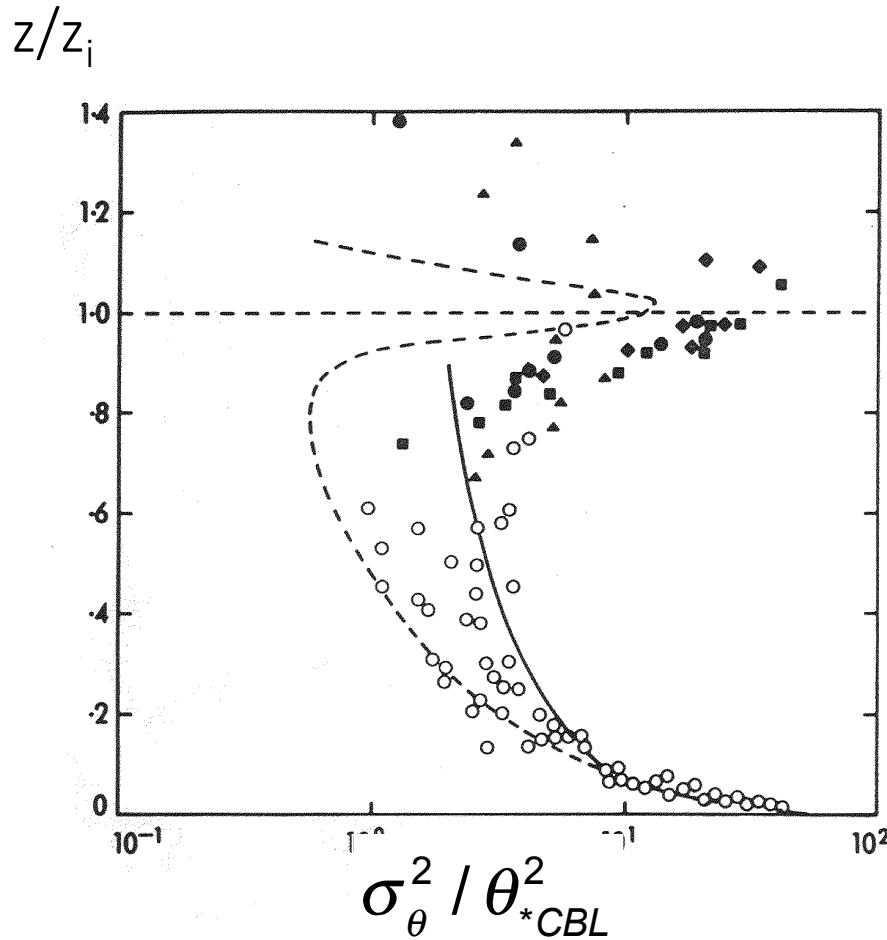
DOI: [10.1175/JAS-D-10-05010.1](https://doi.org/10.1175/JAS-D-10-05010.1)

Scaling regimes: Mixed Layer



CBL
→ different
measurements
(symbols)

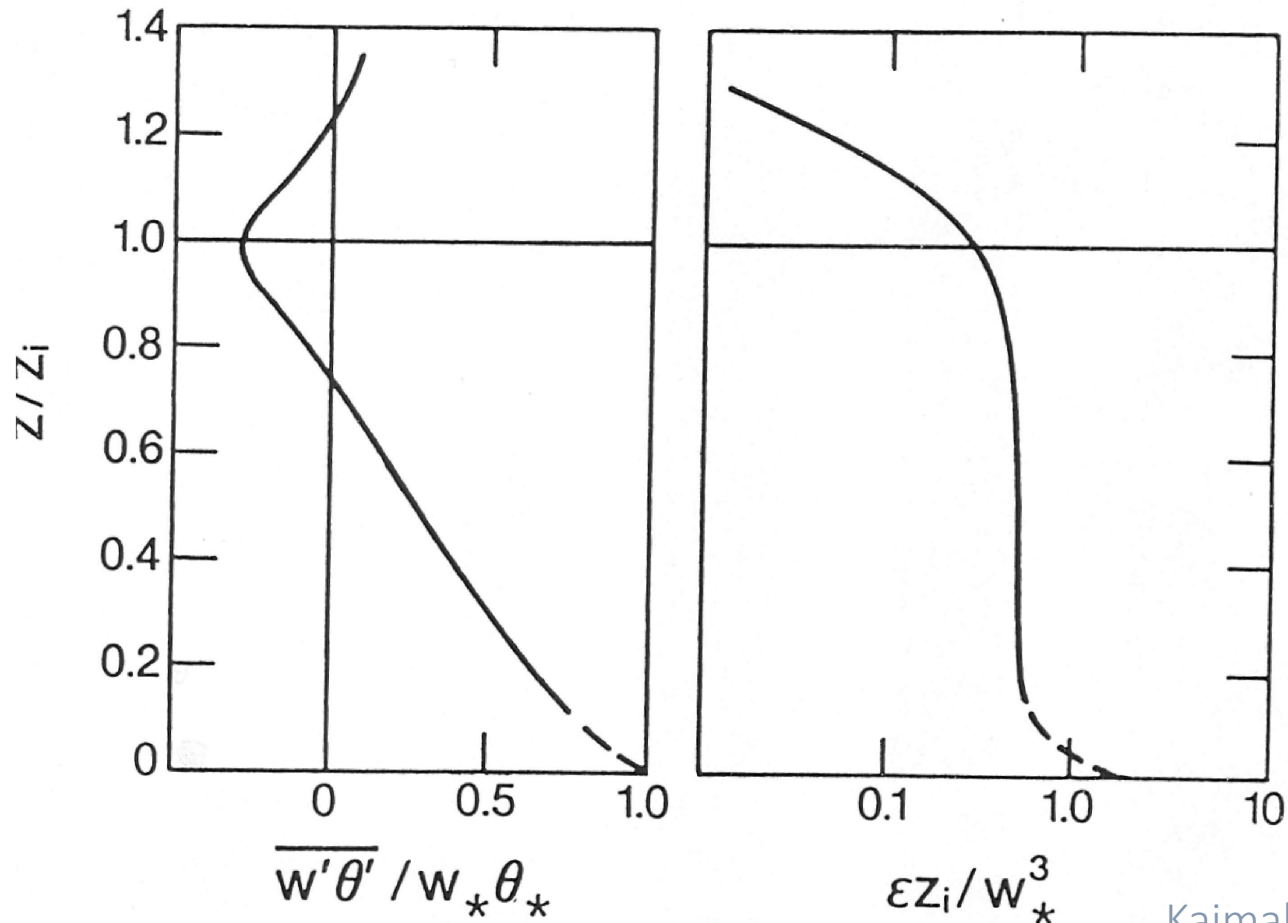
Scaling regimes: Mixed Layer



CBL

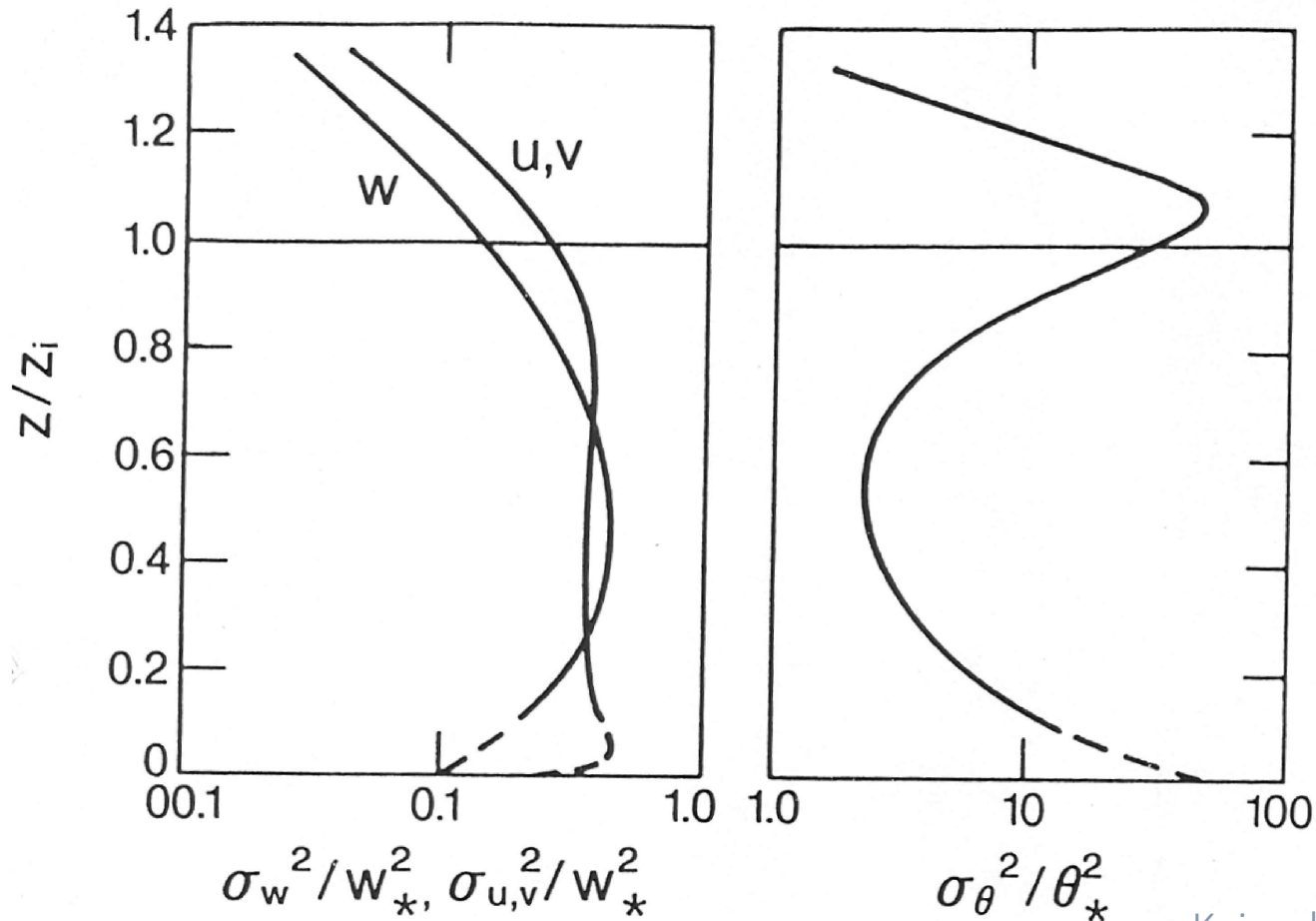
→ different
measurements
(symbols)

Scaling regimes: Mixed Layer



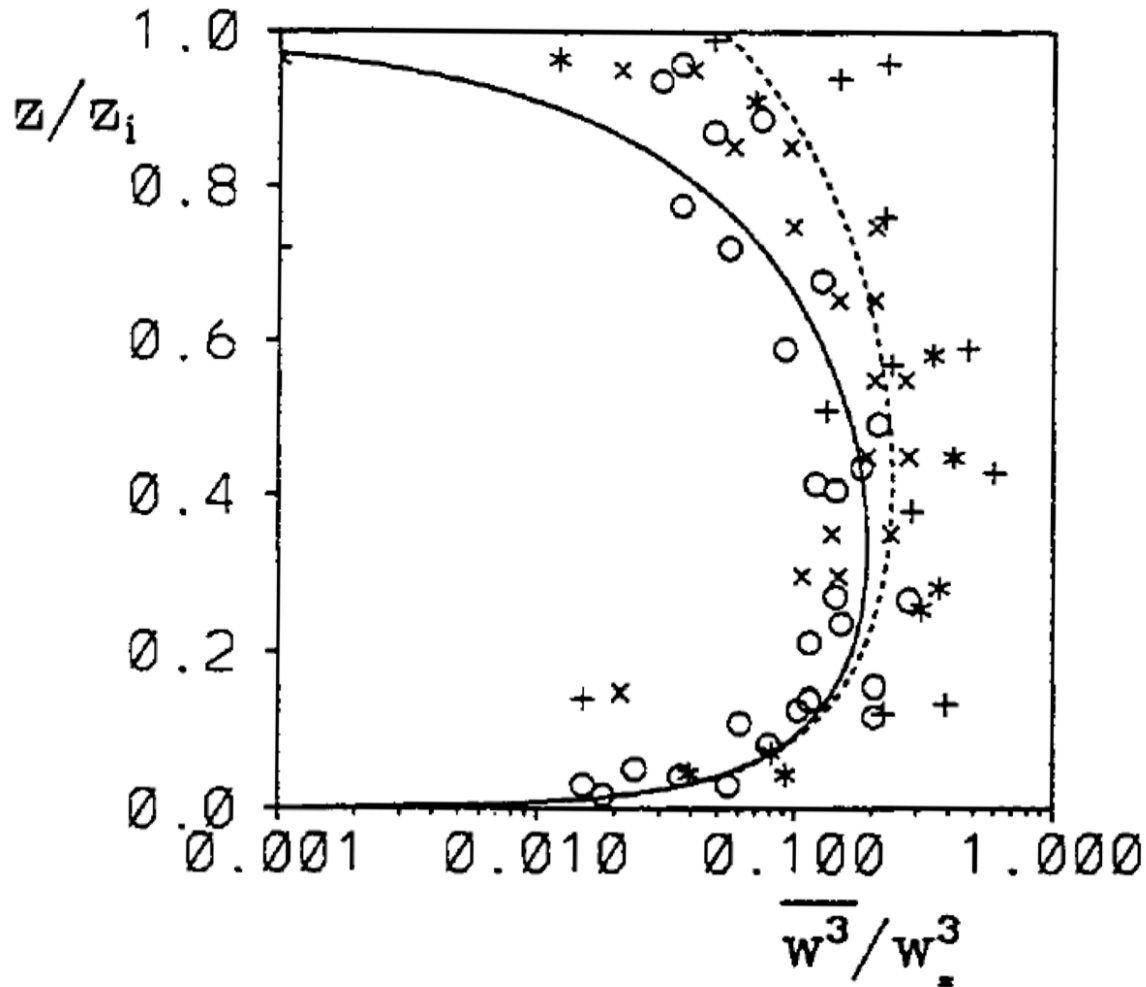
Kaimal and Finnigan (1994)

Scaling regimes: Mixed Layer

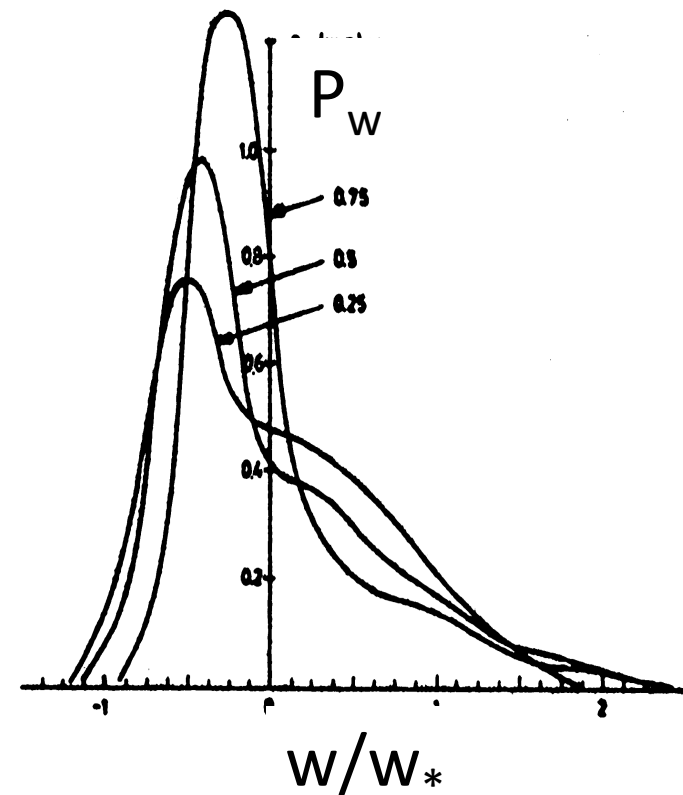


Kaimal and Finnigan (1994)

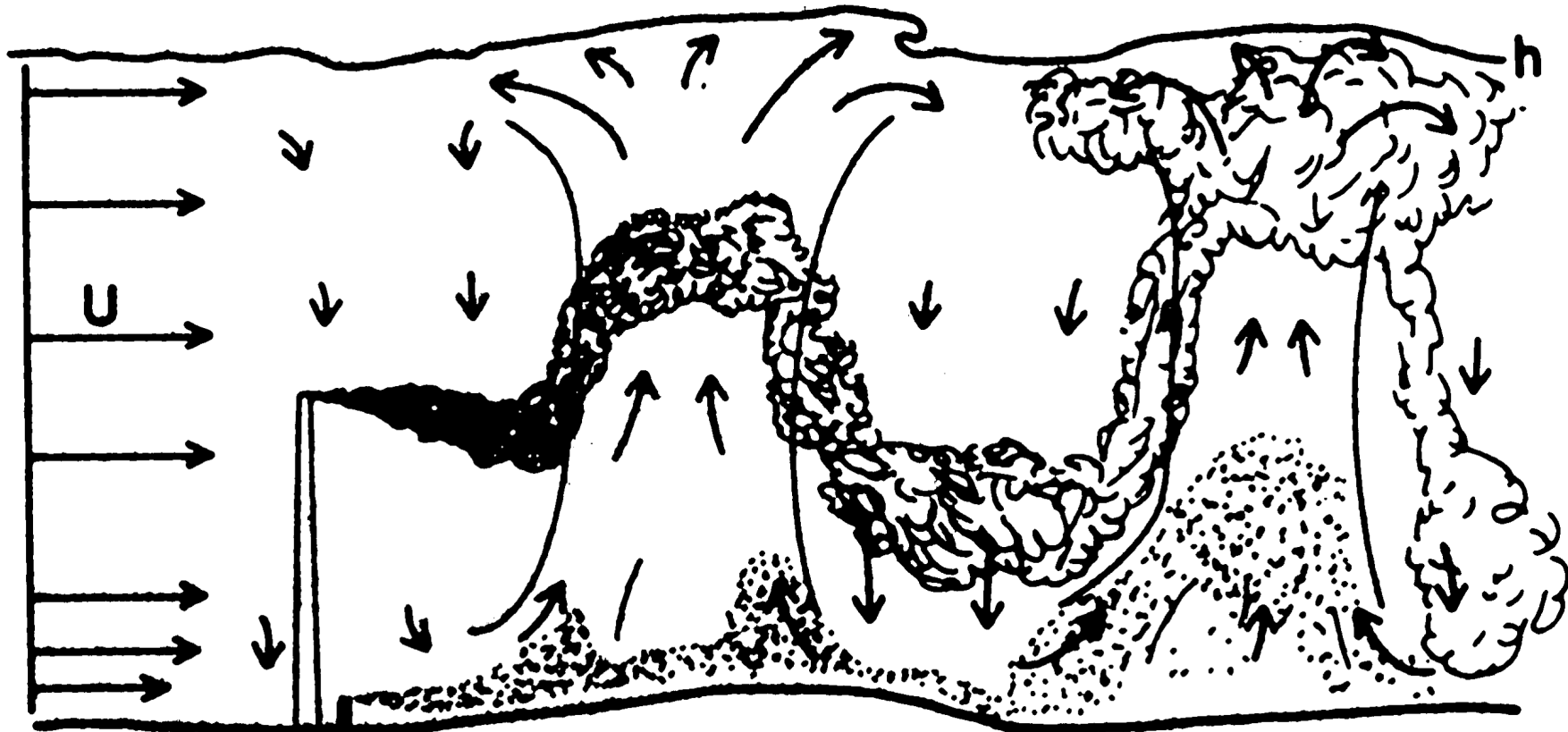
Skewness w: Mixed Layer

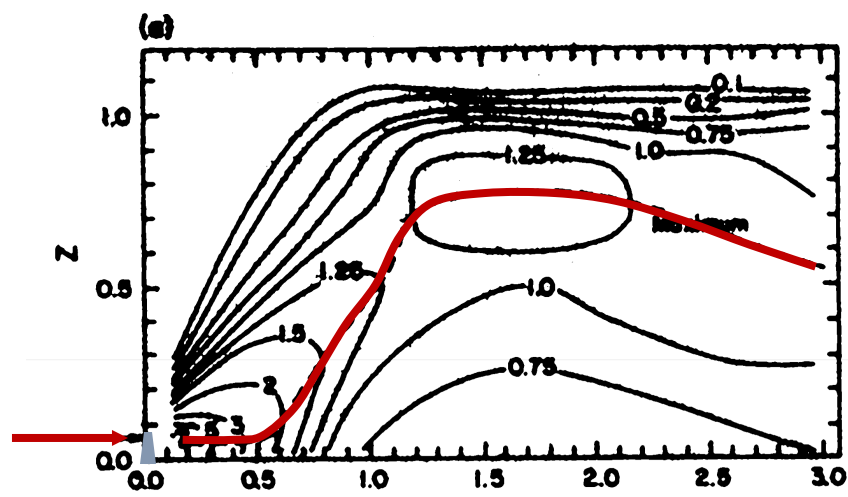


$$\frac{\overline{w^3}}{w_*^3} = f_{w^3}\left(\frac{z}{z_i}\right)$$

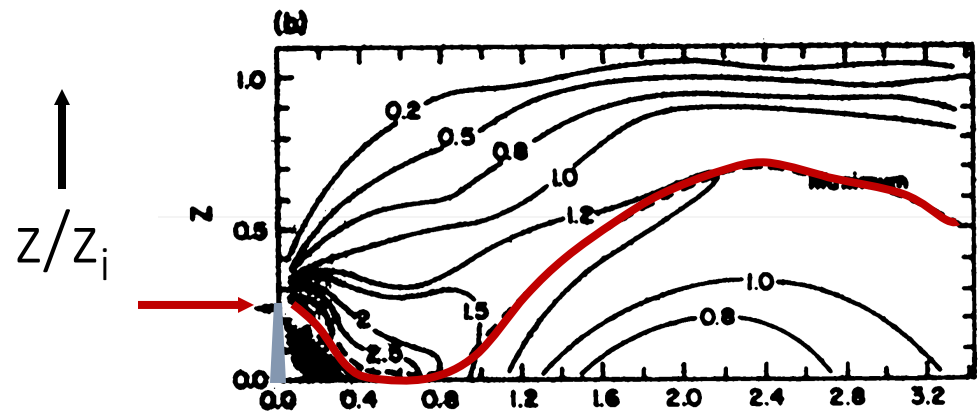


Mixed layer

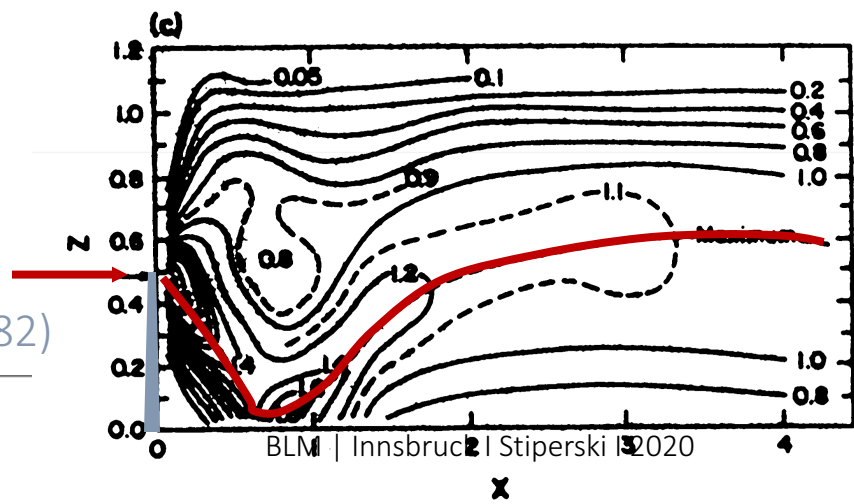




— maximum concentration



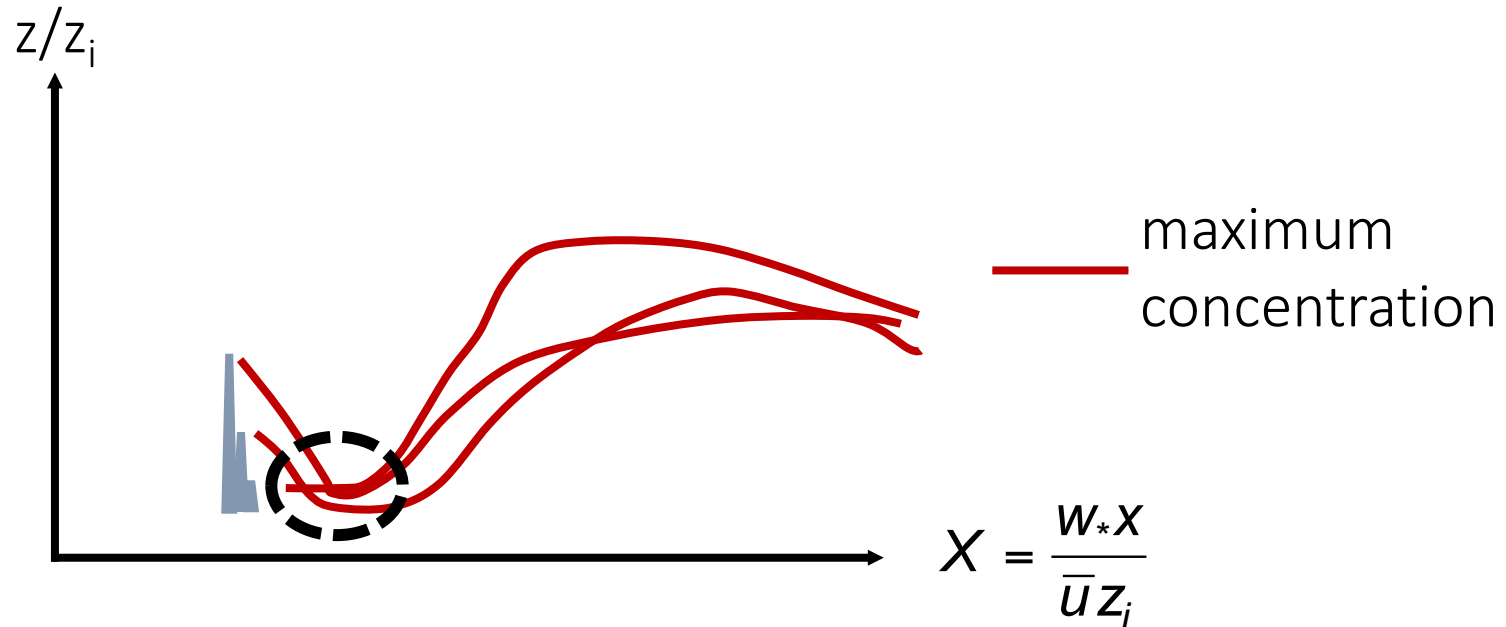
↑
 z/z_i



$$X = \frac{W_* X}{\bar{u} z_i} \longrightarrow$$

Deardorff and Willis (1982)

Scaling regimes: Mixed Layer



- scaling of concentration in ML: not 'similar' for different emission heights
- not only local dispersion process
- non-local (large eddies) mixing

Deardorff and Willis (1982)

Summary: Mixed Layer Scaling

- one π -group: z/z_i
- every scaled mean variable:
$$\frac{\bar{a}}{a_*} = f_a\left(\frac{z}{z_i}\right)$$
- convective velocity scale: w_*
- works for wind & scalar variances, turbulent fluxes, skewness w , ..
- works for spectra (chapter 7)
- ‘not interesting’ for profiles of mean variables (uniform)
- does not work for mean concentrations (non-local influence)

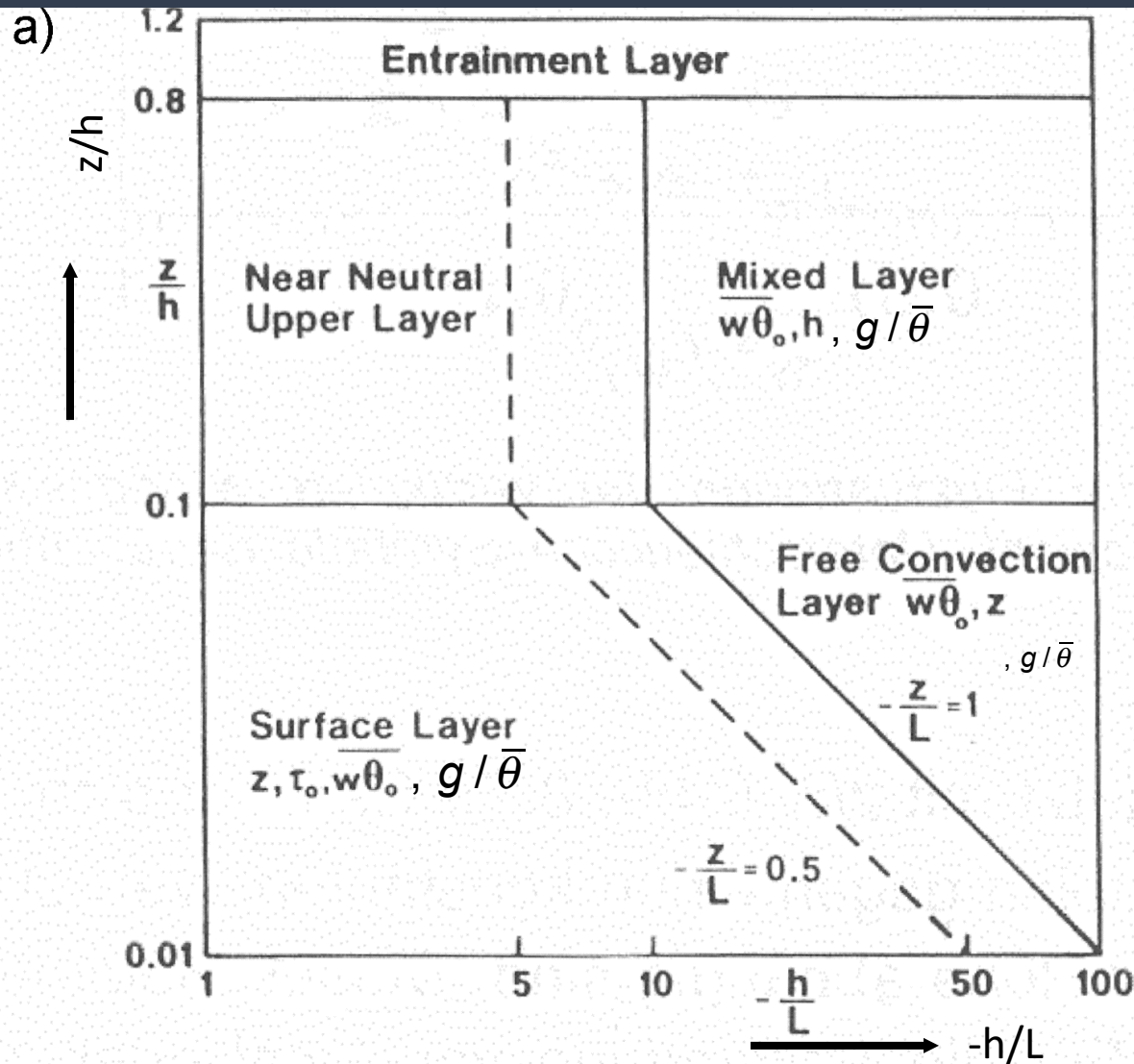
Scaling regimes: The successful Three

→ Surface Layer (MOST) ✓

→ Mixed Layer ✓

→ Local Scaling layer (incl. z-less scaling)

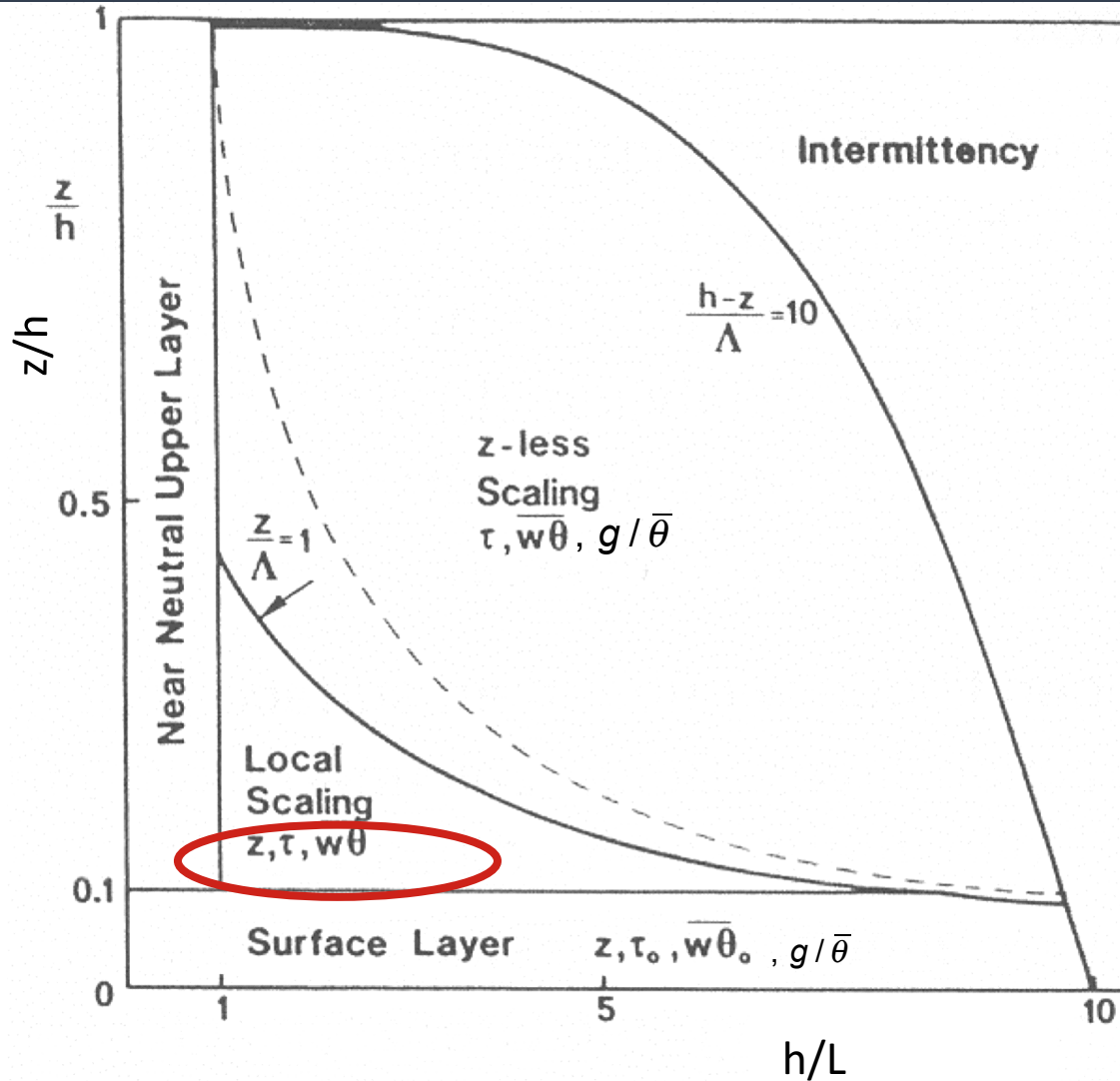
Scaling regimes: Unstable



$$h \hat{=} z_i$$

$$L \equiv -\frac{1}{k} \frac{u_*^3}{\overline{w'\theta'_o}} \left(\frac{g}{\bar{\theta}}\right)^{-1}$$

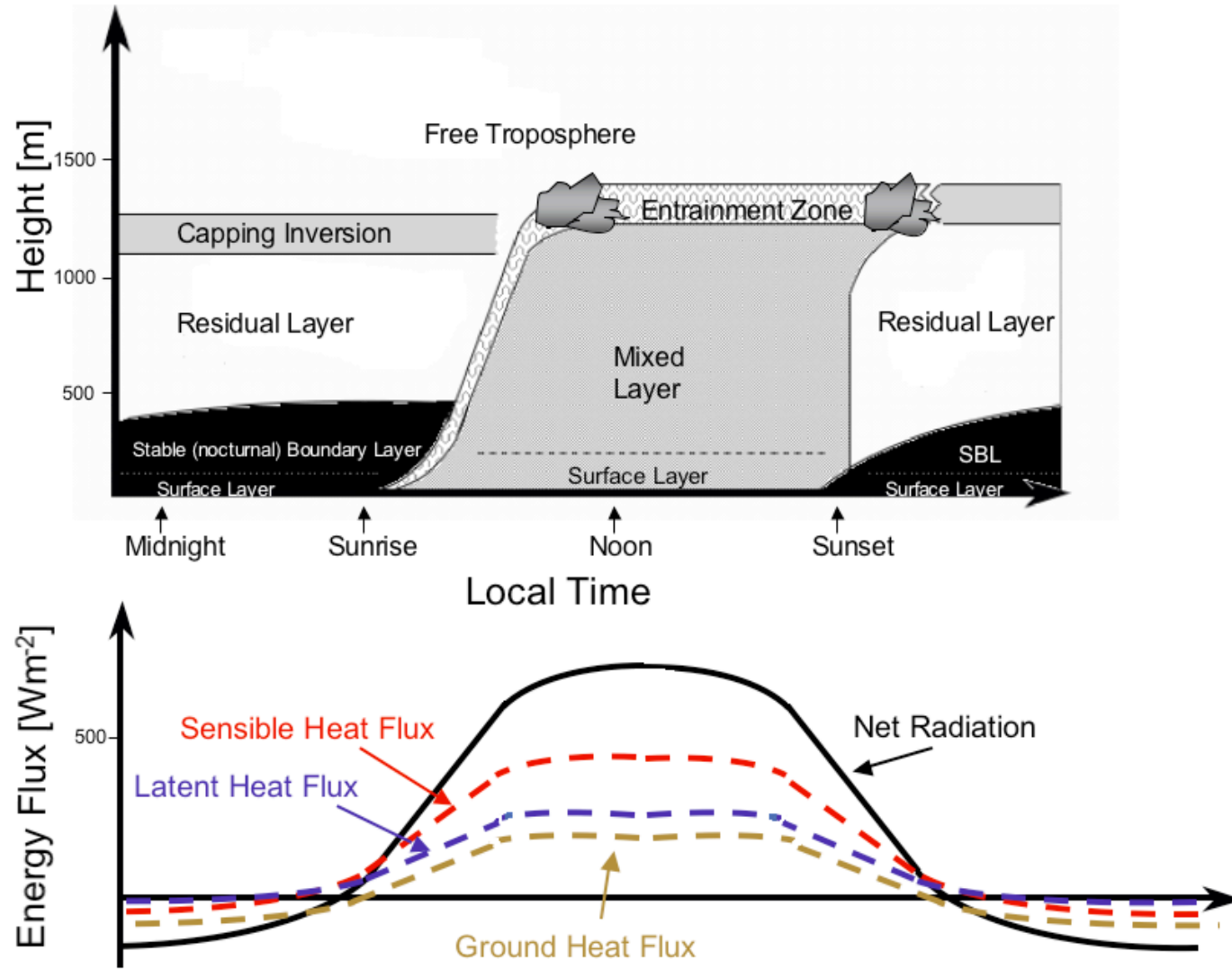
Scaling regimes: Stable



h = SBL height

$$L \equiv -\frac{1}{k} \frac{u_*^3}{w'\theta'_o} \left(\frac{g}{\theta}\right)^{-1}$$

Scaling regimes



Scaling regimes: Local Scaling

- stable stratification (large z/L)
- weak turbulence
- decoupled from surface
- not the surface fluxes for scaling

<u>MOST</u>		<u>LS</u>
$\overline{u'w'_o} \left[= \tau_o \right]$	→	$\overline{u'w'}(z) \left[= \tau(z) \right]$
$\overline{w'\theta'_o}$	→	$\overline{w'\theta'}(z)$
L	→	$\Lambda = \frac{\tau^{3/2}}{k \frac{g}{\theta} \overline{w'\theta'}}$

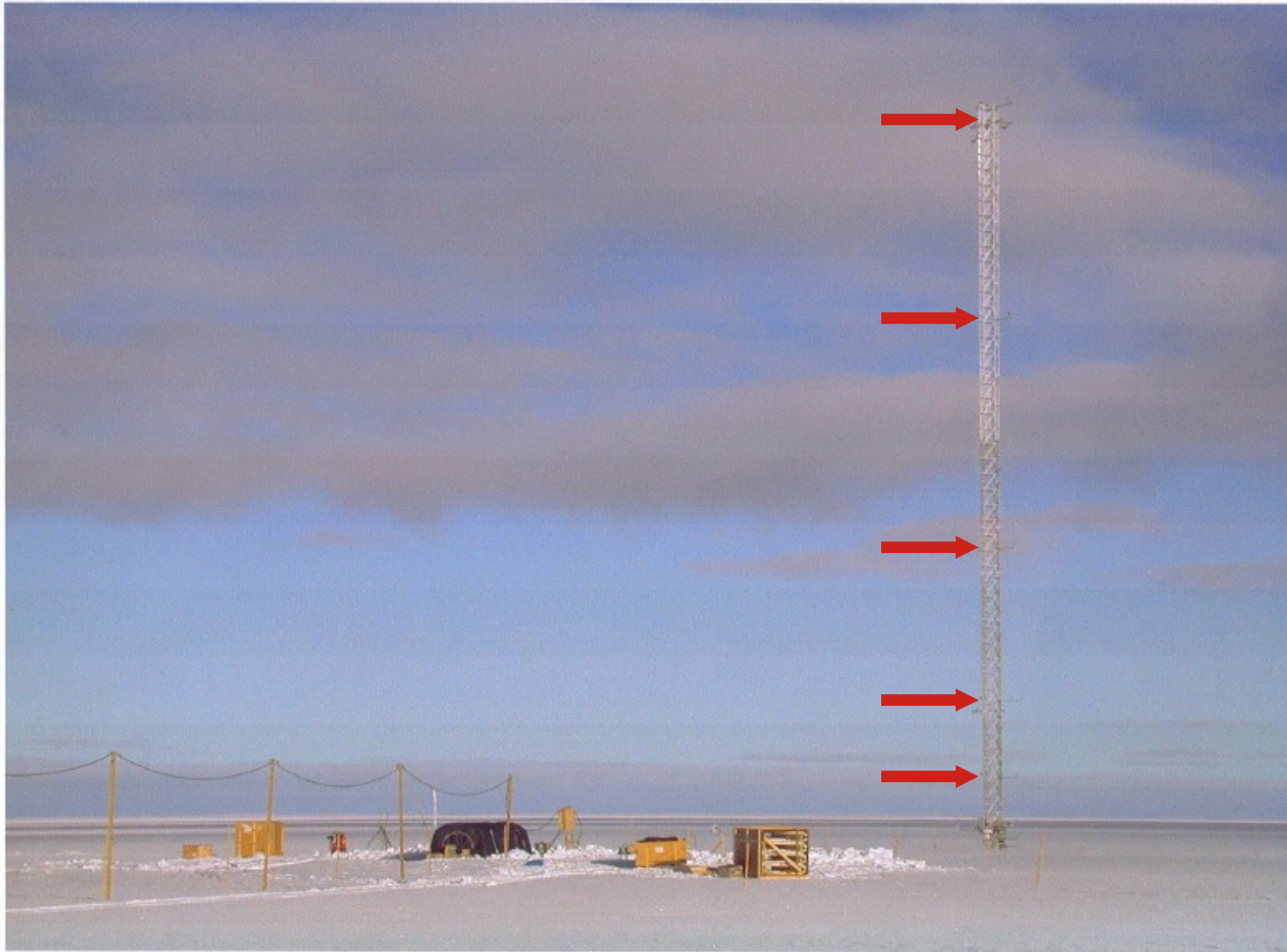
$$\frac{\bar{a}}{a_*} = f_a\left(\frac{z}{\Lambda}\right)$$



local scaling variables

Scaling regimes: Local Scaling

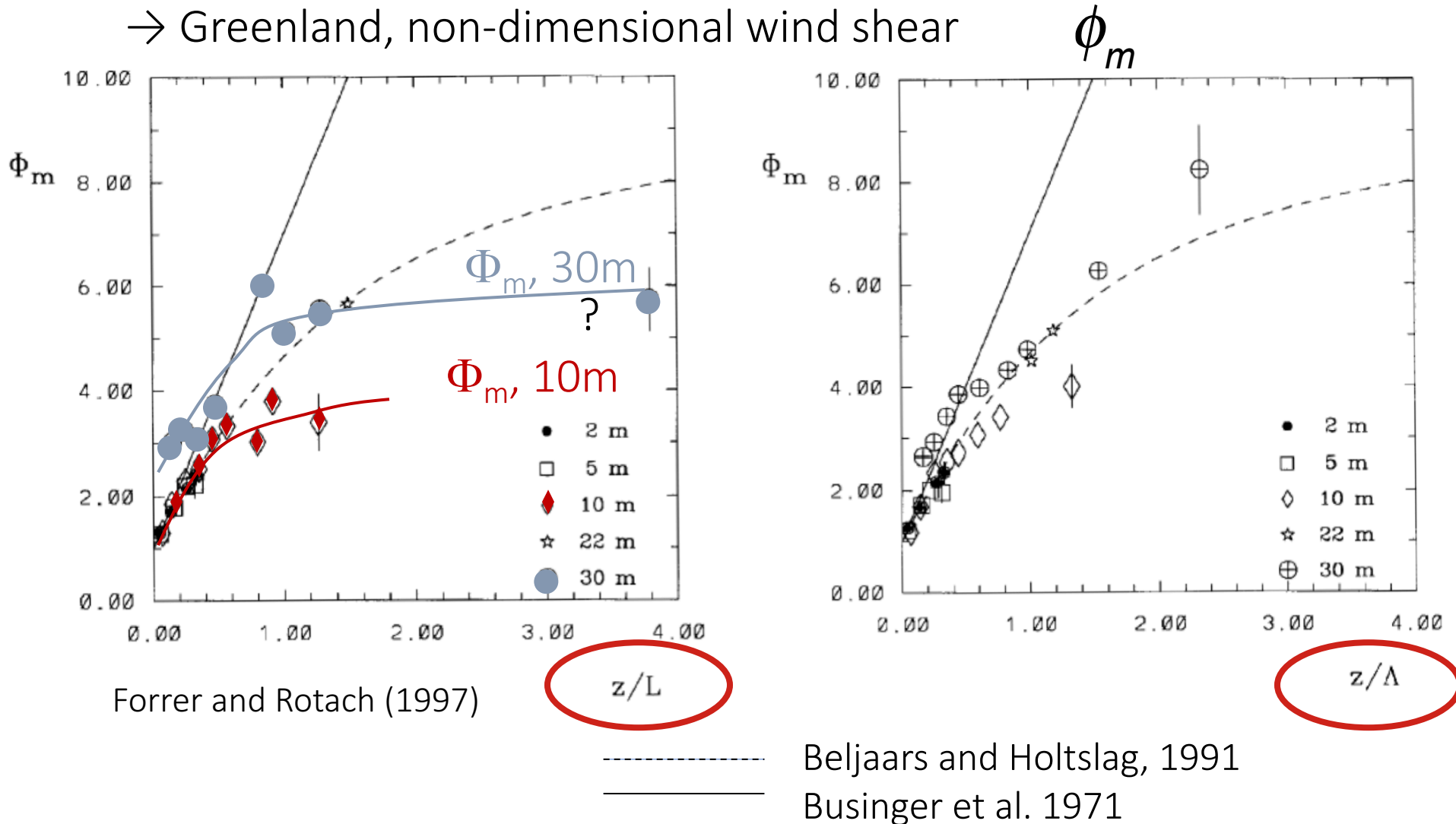
→ Greenland, non-dimensional wind shear



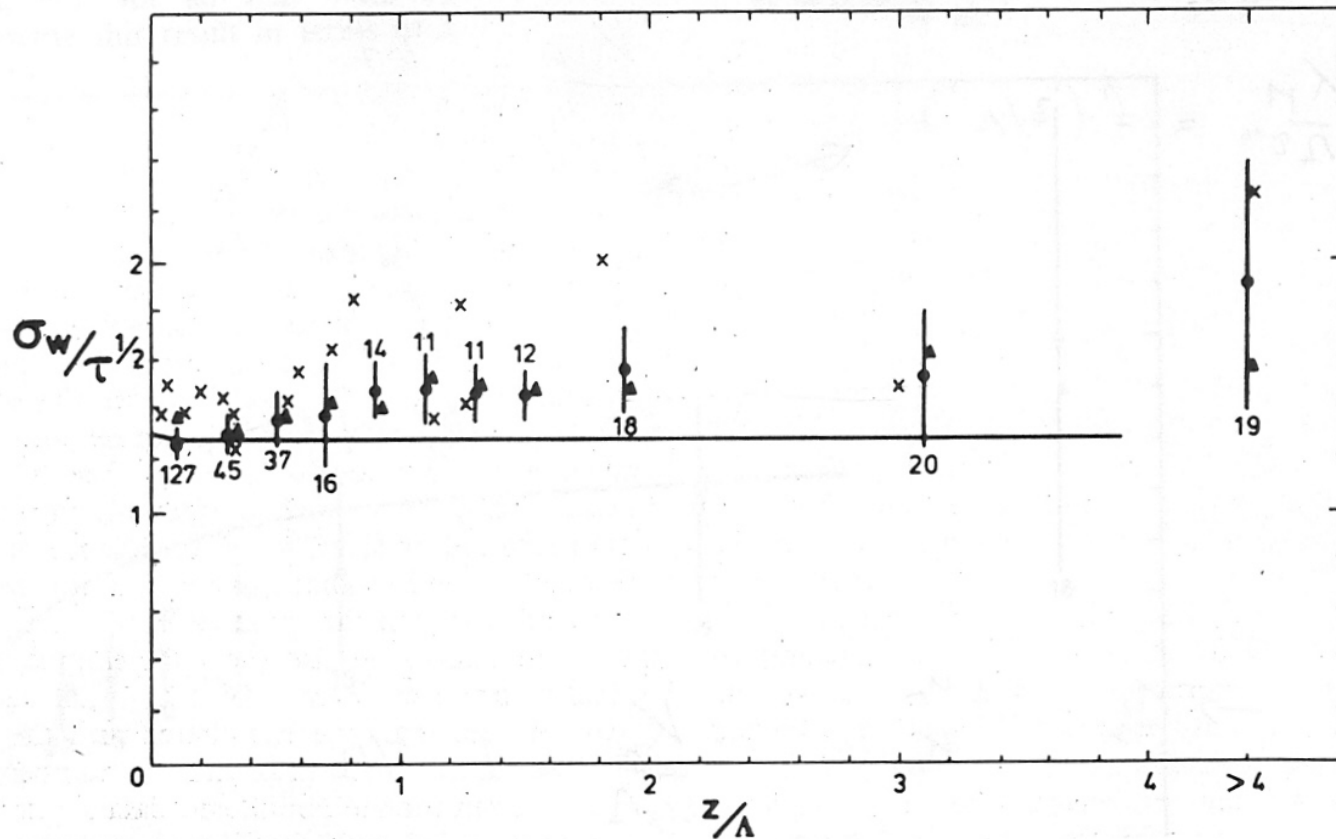
$$\frac{kz}{u_*} \frac{\partial \bar{u}}{\partial z} = \phi_m$$

Scaling regimes: Local Scaling

→ Greenland, non-dimensional wind shear



Scaling regimes: Local Scaling

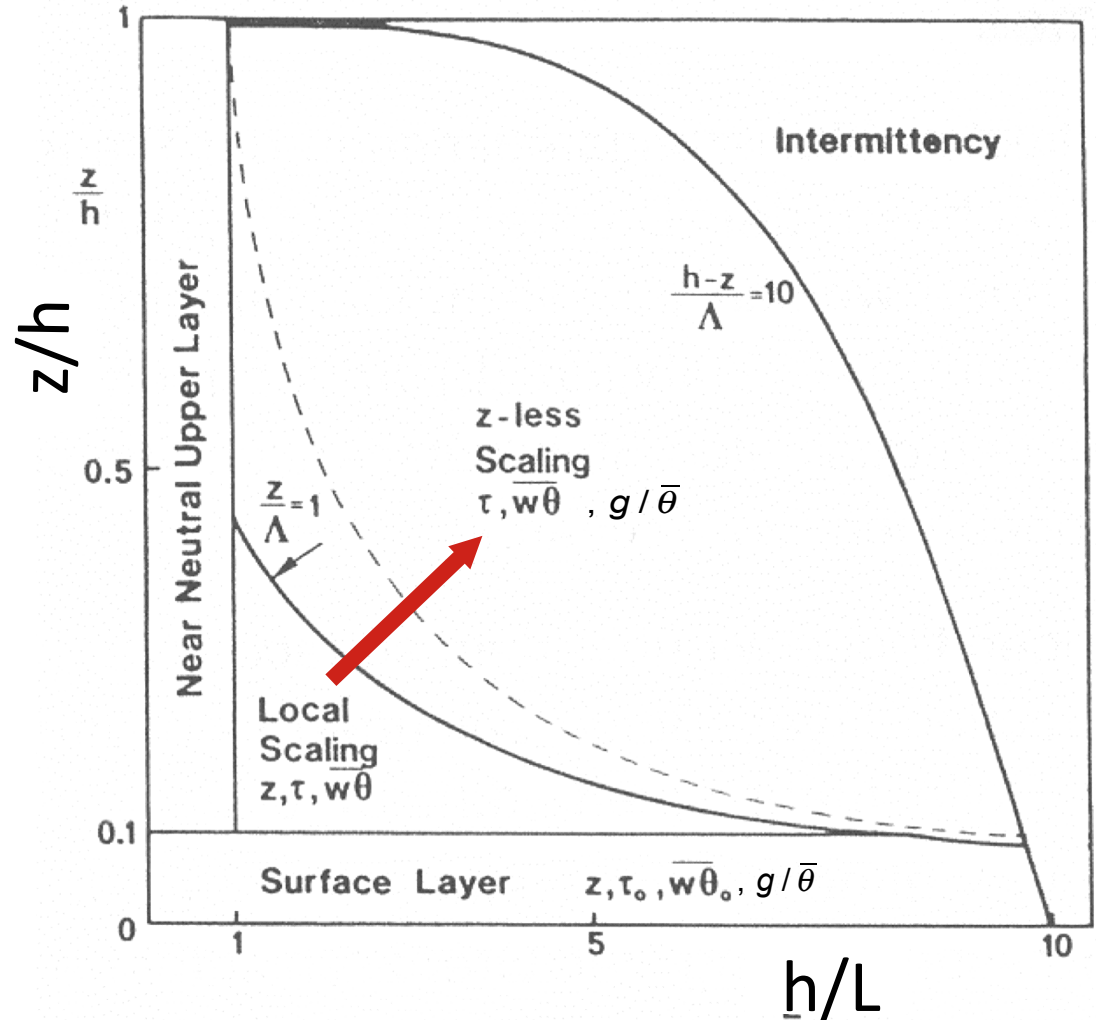


Nieuwstadt (1984)

Local Scaling: z-less limit

$h = \text{SBL height}$

$$L \equiv -\frac{1}{k} \frac{u_*^3}{w' \theta'_o} \left(\frac{g}{\bar{\theta}} \right)^{-1}$$



Scaling regimes: z-less Scaling

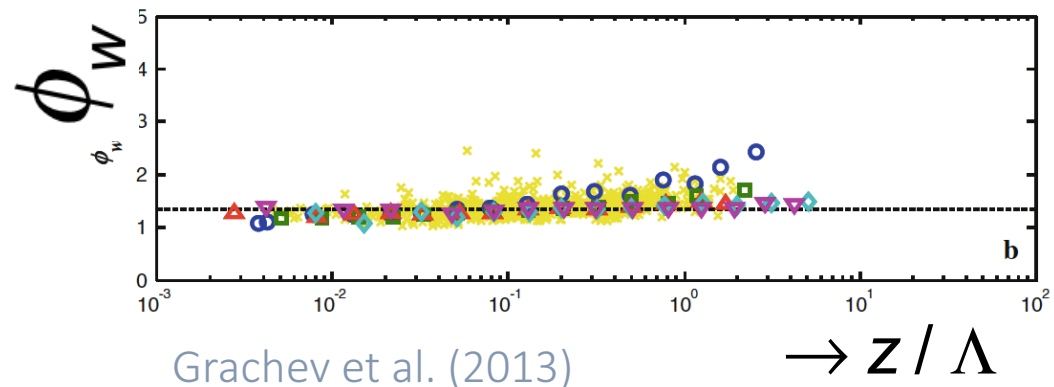
z-less limit: height z no longer important

→ stability so strong that surface influence vanishes

→ Consequences:

Flux-variance relations = const

$$\frac{\sigma_w}{u_{*l}} (= \phi_w(z/\Lambda) = \text{const})$$



Scaling regimes: z-less Scaling

z-less limit: height z no longer important

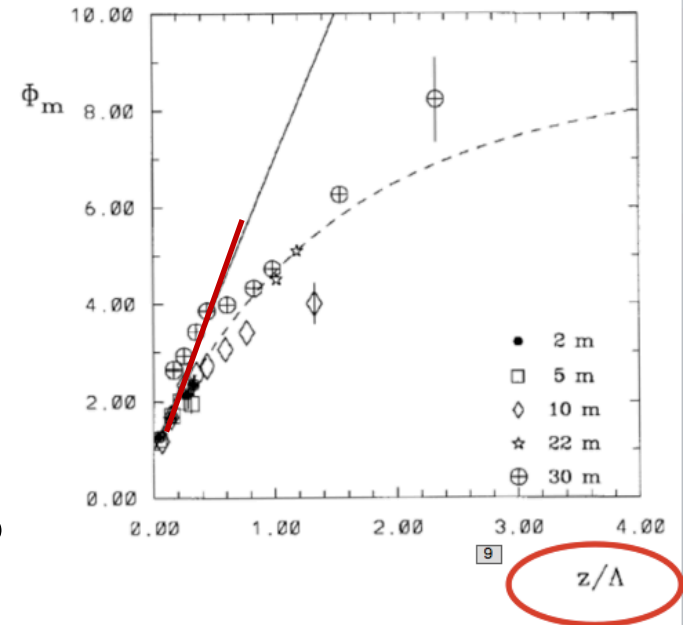
→ stability so strong that surface influence vanishes

→ Consequences:

Flux-gradient relations = linear

$$\frac{\partial \bar{u}}{\partial z} \frac{kz}{u_{*l}} (= \phi_m(z/\Lambda)) = \text{const.} \cdot \frac{z}{\Lambda}$$

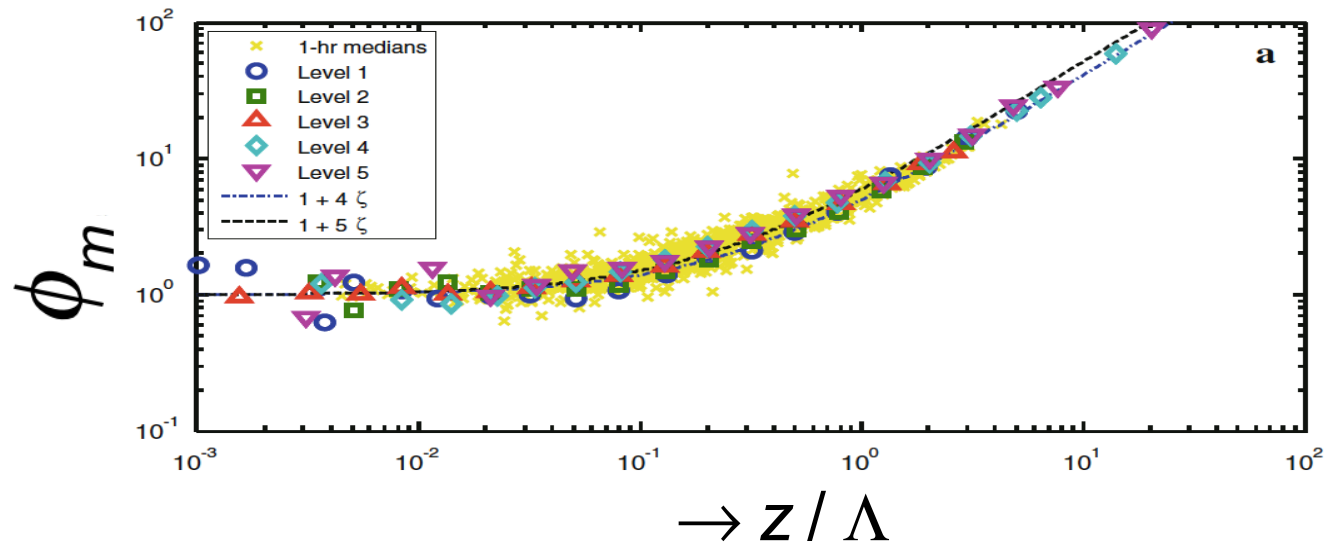
?



Scaling regimes: z-less Scaling

z-less limit: height z no longer important

$$\frac{\partial \bar{u}}{\partial z} \frac{kz}{u_{*l}} = 1 + \beta \frac{z}{\Lambda}$$



→ z-less: only if $R_f < 0.25$ and $Ri < 0.25$

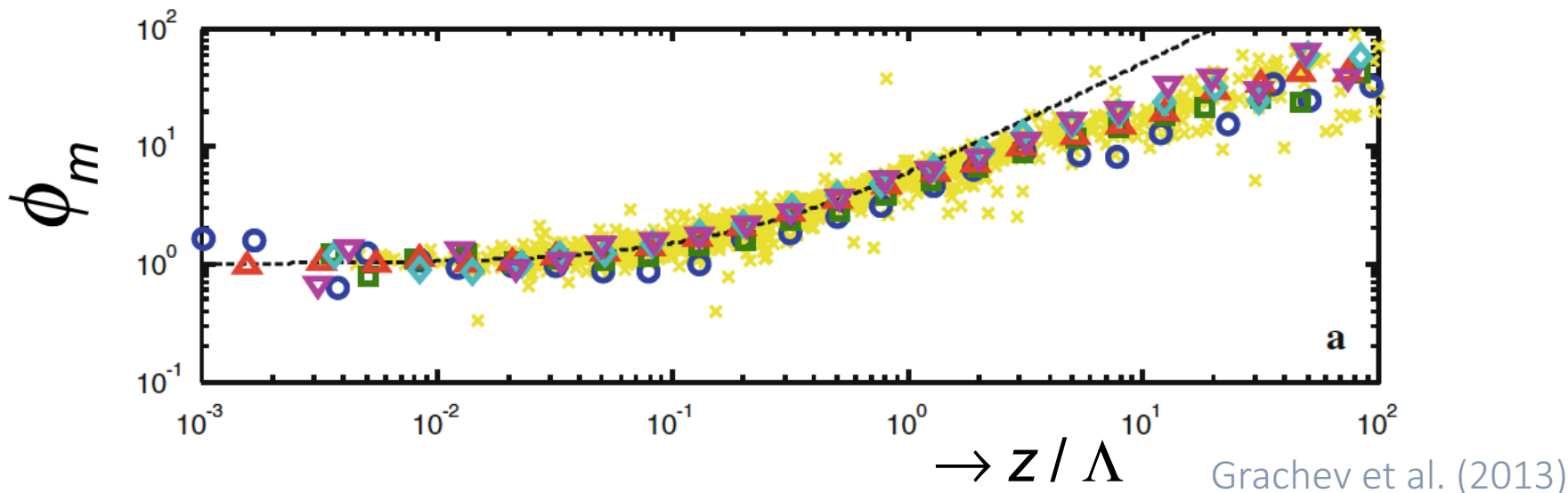
→ 'well-behaved' (fully developed) turbulence

Grachev et al. (2013)

Scaling regimes: z-less Scaling

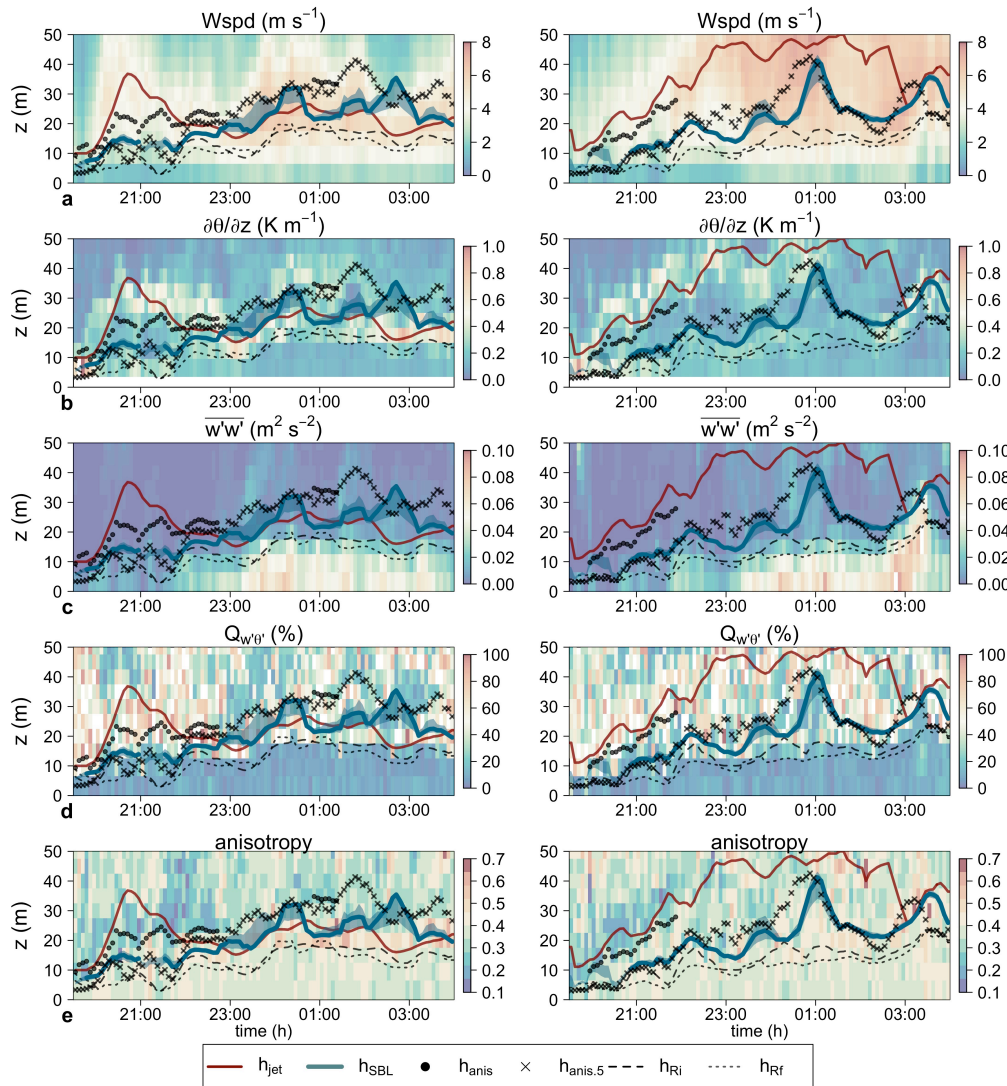
z-less limit: height z no longer important

$$\frac{\partial \bar{u}}{\partial z} \frac{kz}{u_{*r}} = 1 + \beta \frac{z}{\Lambda} \quad \longleftrightarrow \quad \text{fully developed turbulence}$$



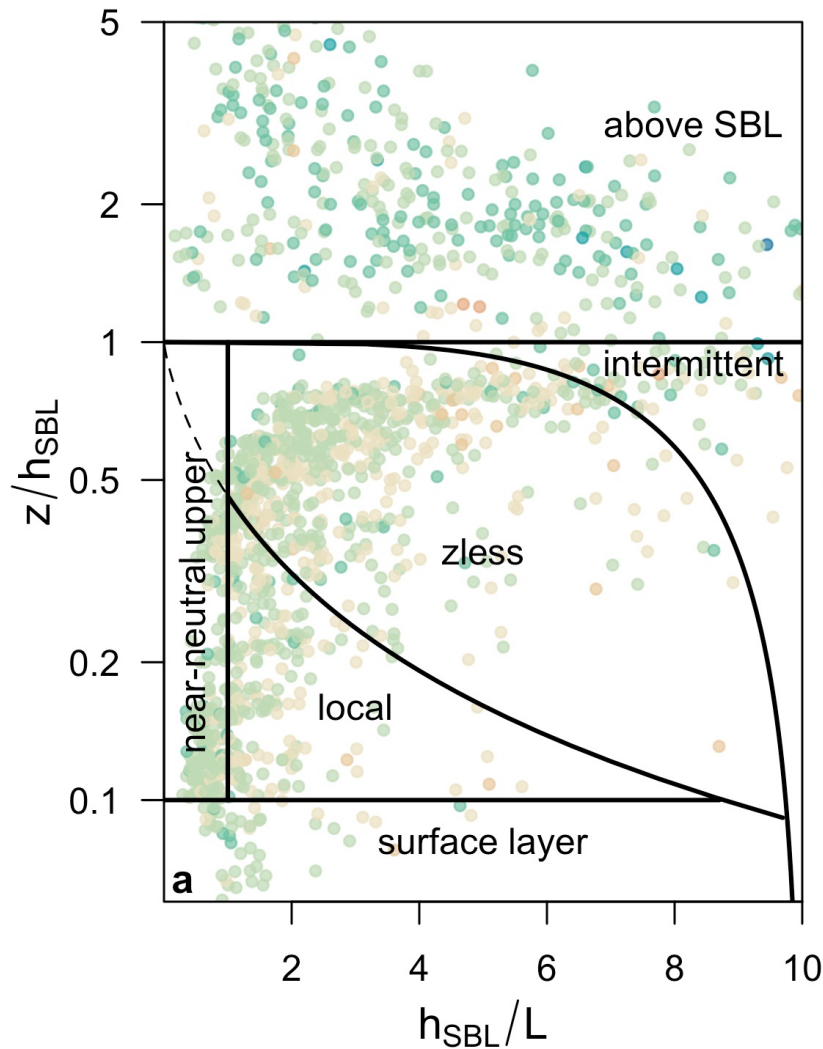
- for 'all conditions' (Ri , R_f): closer to Beljaars and Holtslag formulation
- deviation from 'z-less' behavior

Stable scaling regimes: shallow slope



- 50 m tower on a gentle slope (1°)
- Katabatic flows form
- Top of PBL can be detected from tower measurements
- Ambiguity of PBL top estimate
- But we can test scaling regimes

Stable scaling regimes: shallow slope

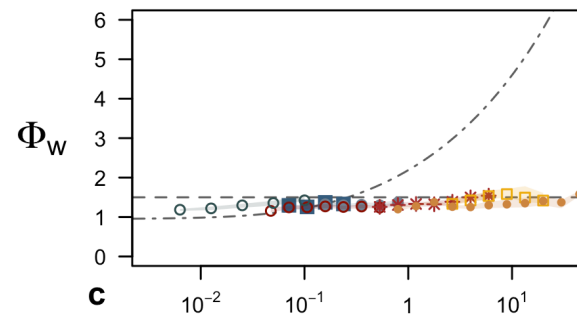
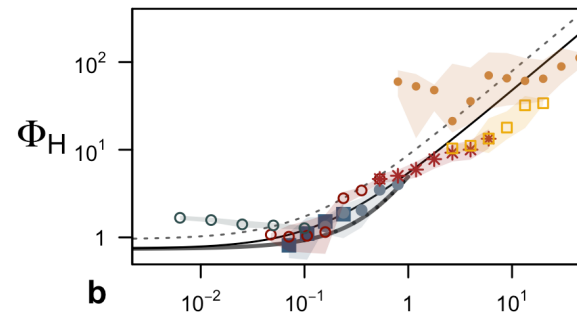
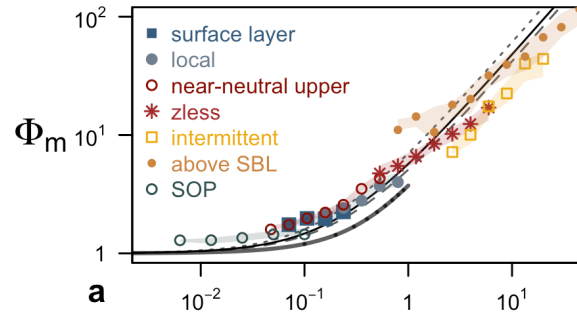


Data from all scaling regimes:

- Surface layer (1st level) – for high h
- Local scaling
- z-less scaling
- Intermittent
- Above SBL

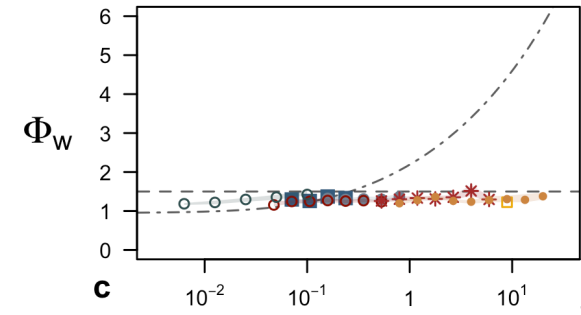
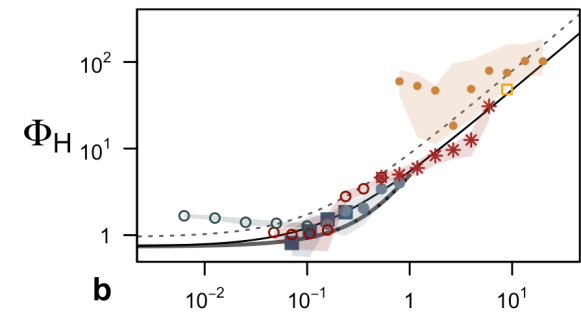
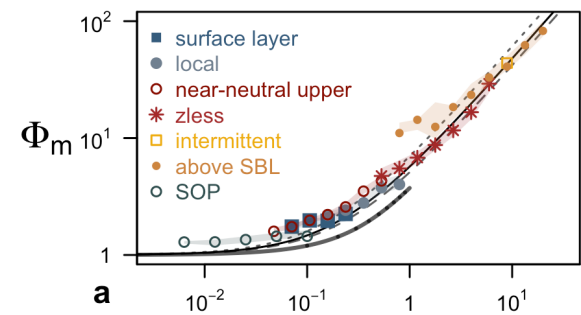
Stable scaling regimes: shallow slope

- Stationary data from each regime



z/Λ

- Stationary & $Ri < 0.21$



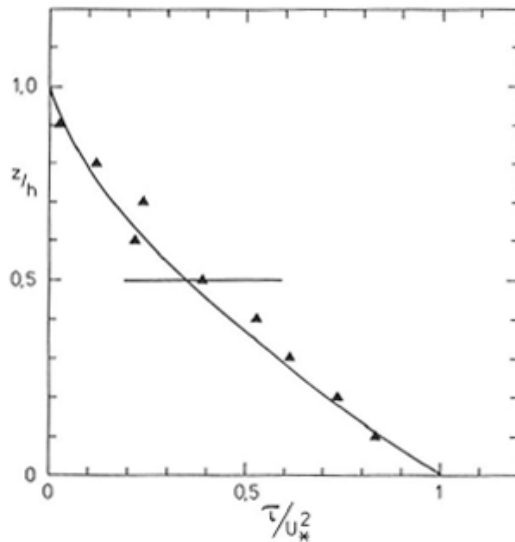
z/Λ

Scaling regimes: Local Scaling

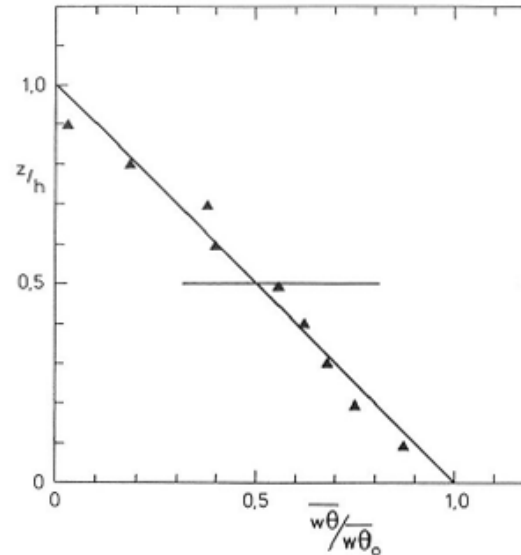
For **Local scaling** to be successful:

Surface fluxes → we need **local flux profiles** (h = SBL height)

[theoretical derivation (local scaling), Nieuwstadt (1984)]



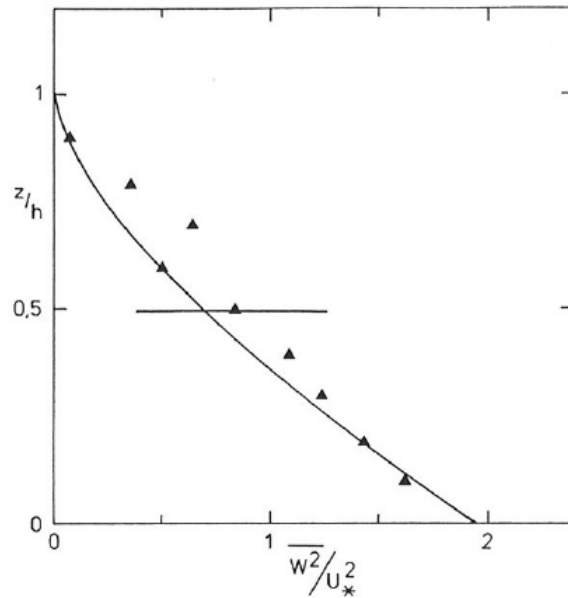
$$\frac{\tau}{\tau_o} = \left(1 - \frac{z}{h}\right)^{3/2}$$



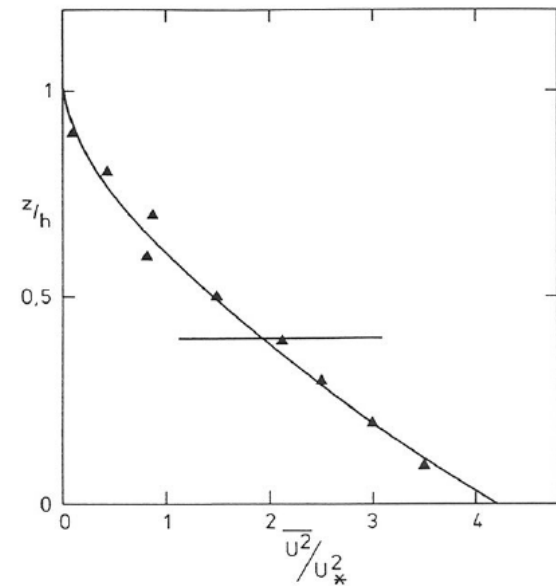
$$\frac{\overline{w'\theta'}}{\overline{w'\theta'_o}} = \left(1 - \frac{z}{h}\right) \text{ Nieuwstadt (1984)}$$

Scaling regimes: Local Scaling

→ profiles of variances



$$\frac{\sigma_w}{u_*} = 1.4 \left(1 - \frac{z}{h}\right)^{3/4}$$

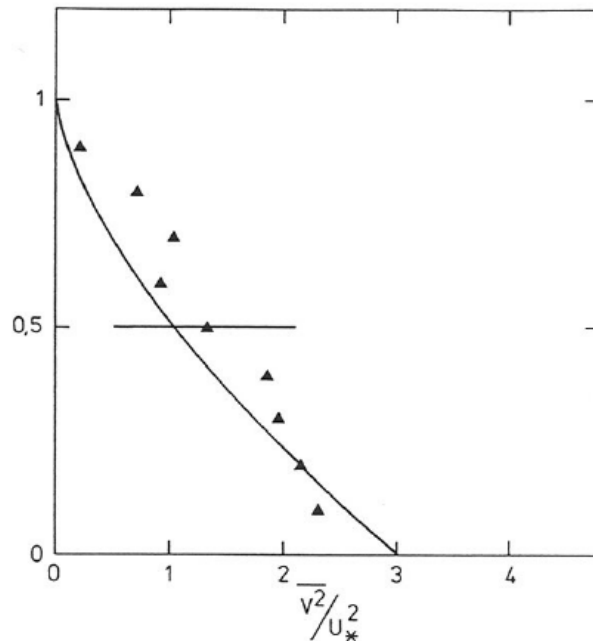


$$\frac{\sigma_u}{u_*} = 2.04 \left(1 - \frac{z}{h}\right)^{3/4}$$

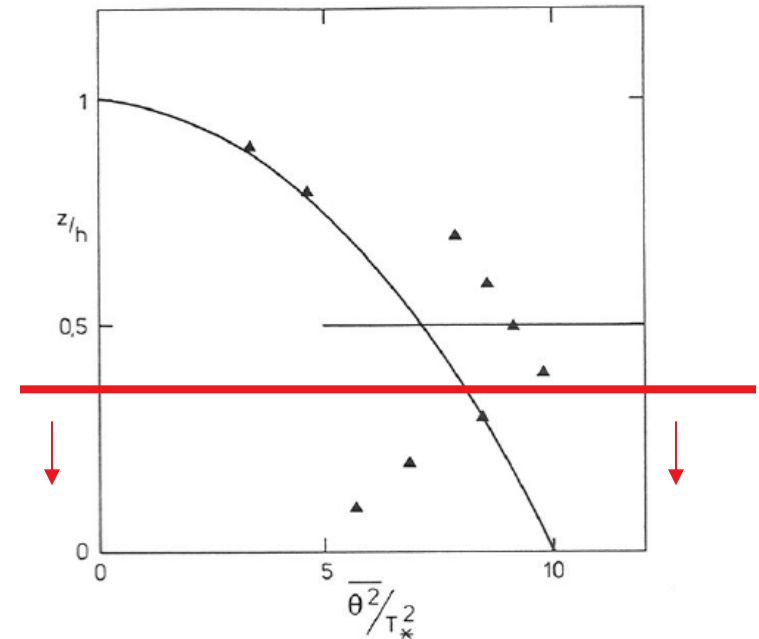
Nieuwstadt (1984)

Scaling regimes: Local Scaling

→ profiles of variances



$$\frac{\sigma_v}{u_*} = 1.7 \left(1 - \frac{z}{h}\right)^{3/4}$$



$$\frac{\sigma_\theta}{\theta_*} = 3.2 \left(1 - \frac{z}{h}\right)^{1/4}$$

Nieuwstadt (1984)

Summary: Local Scaling

→ one π -group: z / Λ

→ every scaled mean variable:

$$\frac{\bar{a}}{a_*} = f_a\left(\frac{z}{\Lambda}\right)$$

→ works for wind & turbulent fluxes

→ works for spectra (Chapter 7)

→ apparently difficult to assess (weak turbulence)

→ does not seem to work for temperature variance

→ subtle border between strongly stable (→ intermittency)
and local / z-less scaling

Scaling regimes

- successful \leftrightarrow quasi-stationary
- Surface Layer, Mixed Layer, Local Scaling
- 'mean' = characteristic profiles
- dependent on a few characteristic variables

→ turbulent fluxes: $\overline{u' w'_{[o]}}$ $\overline{w' \theta'_{[o]}}$

→ length scales: z_i, L, Λ

Scaling in complex terrain

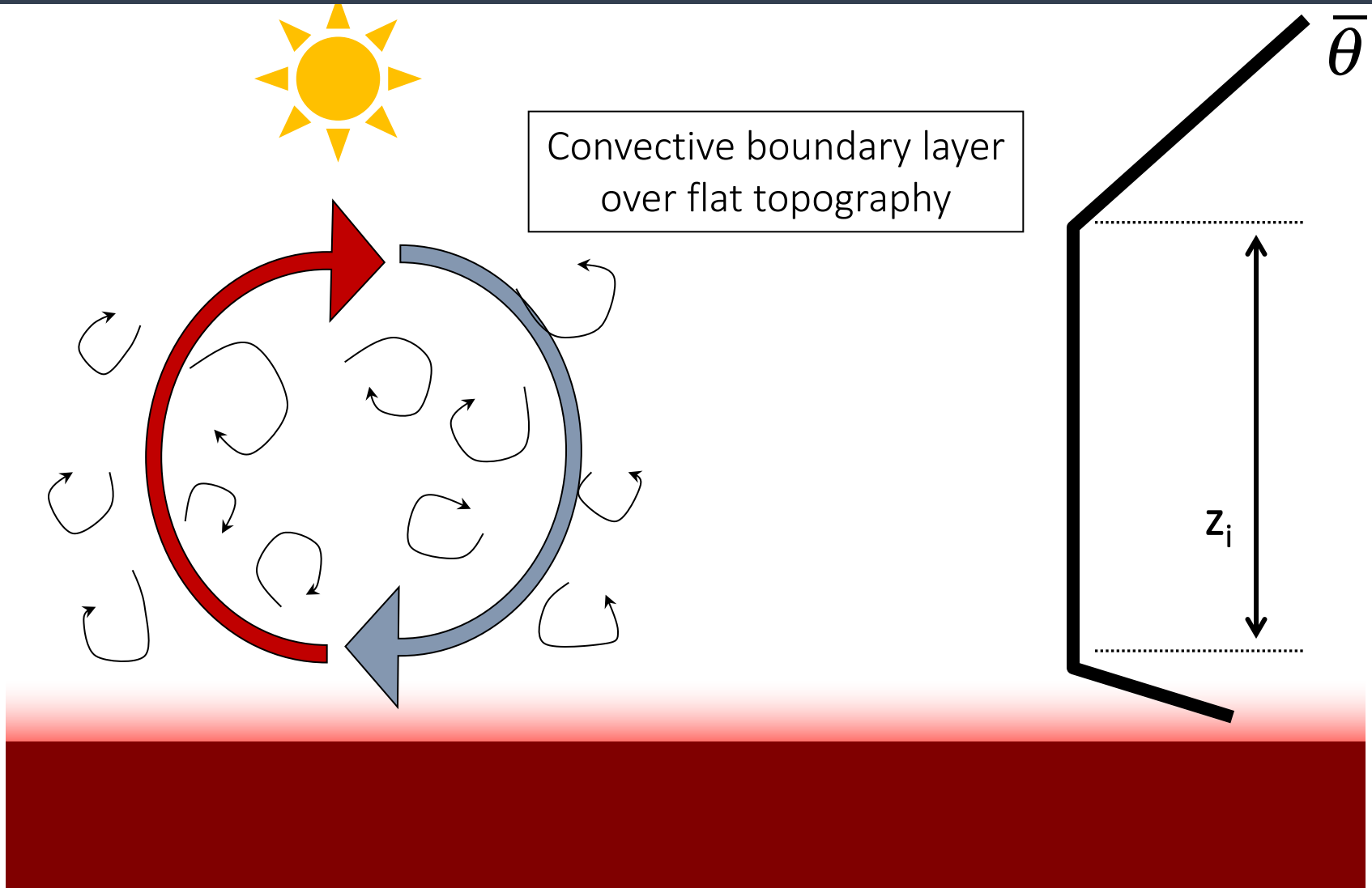
- now that scaling regimes are established in homogeneous terrain...
- how about completely different conditions?
 - not horizontally homogeneous
 - not flat
- do we consider all relevant processes?

Scaling in complex terrain

Application Examples:

1. **Mixed-layer TKE scaling** (MAP Riviera)
2. Wind directional change with height
3. Local scaling flux-variance relationships (12 Datasets)

1. Mixed layer scaling in complex terrain



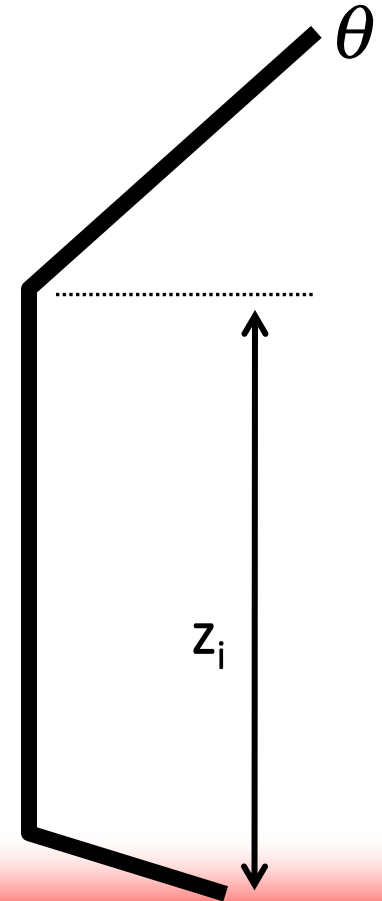
1. Mixed layer scaling in complex terrain

Mixed Layer Scaling

Deardorff (1970,1972)

- example:
$$\frac{TKE}{w_*^2} = f_{TKE} \left(\frac{z}{z_i} \right)$$

- velocity scale
$$w_* = \left(\overline{w'\theta'}_o \frac{g}{\theta} z_i \right)^{1/3}$$

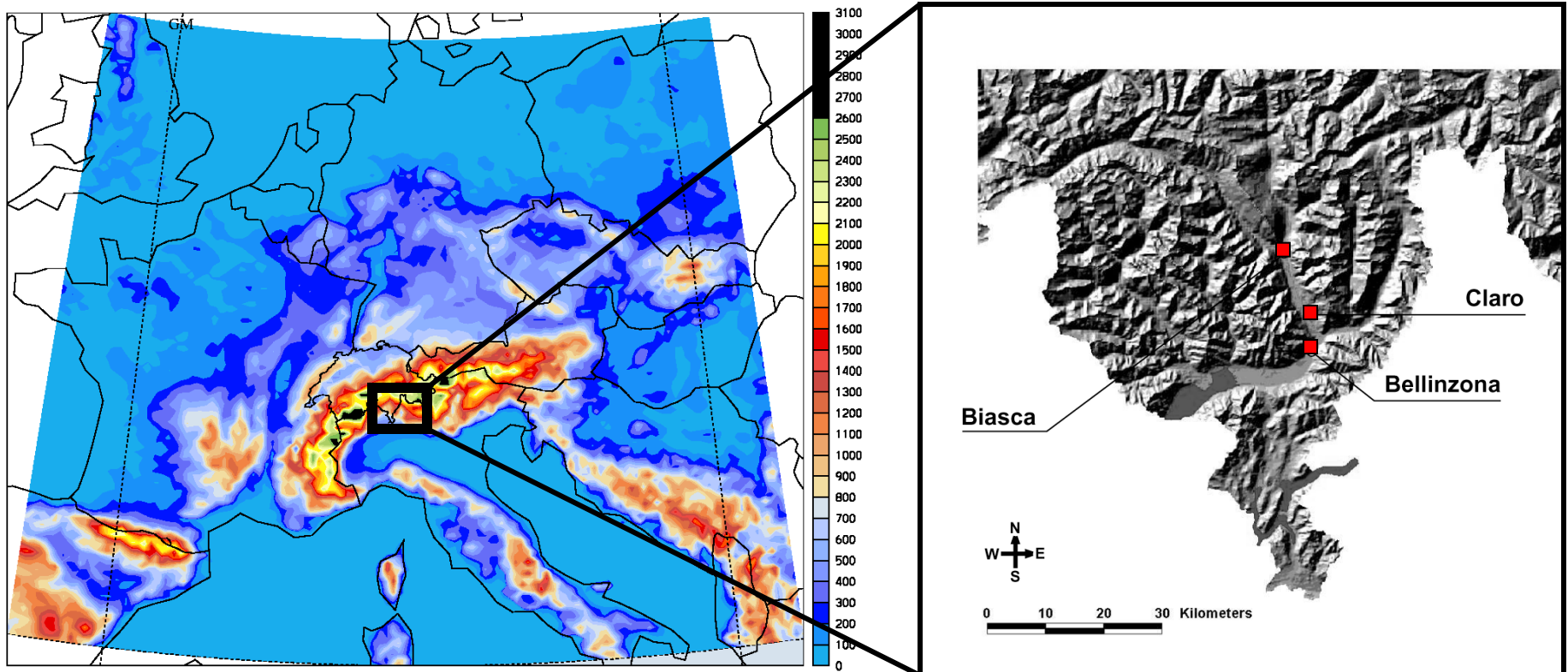


Sensible surface heat flux: H_o

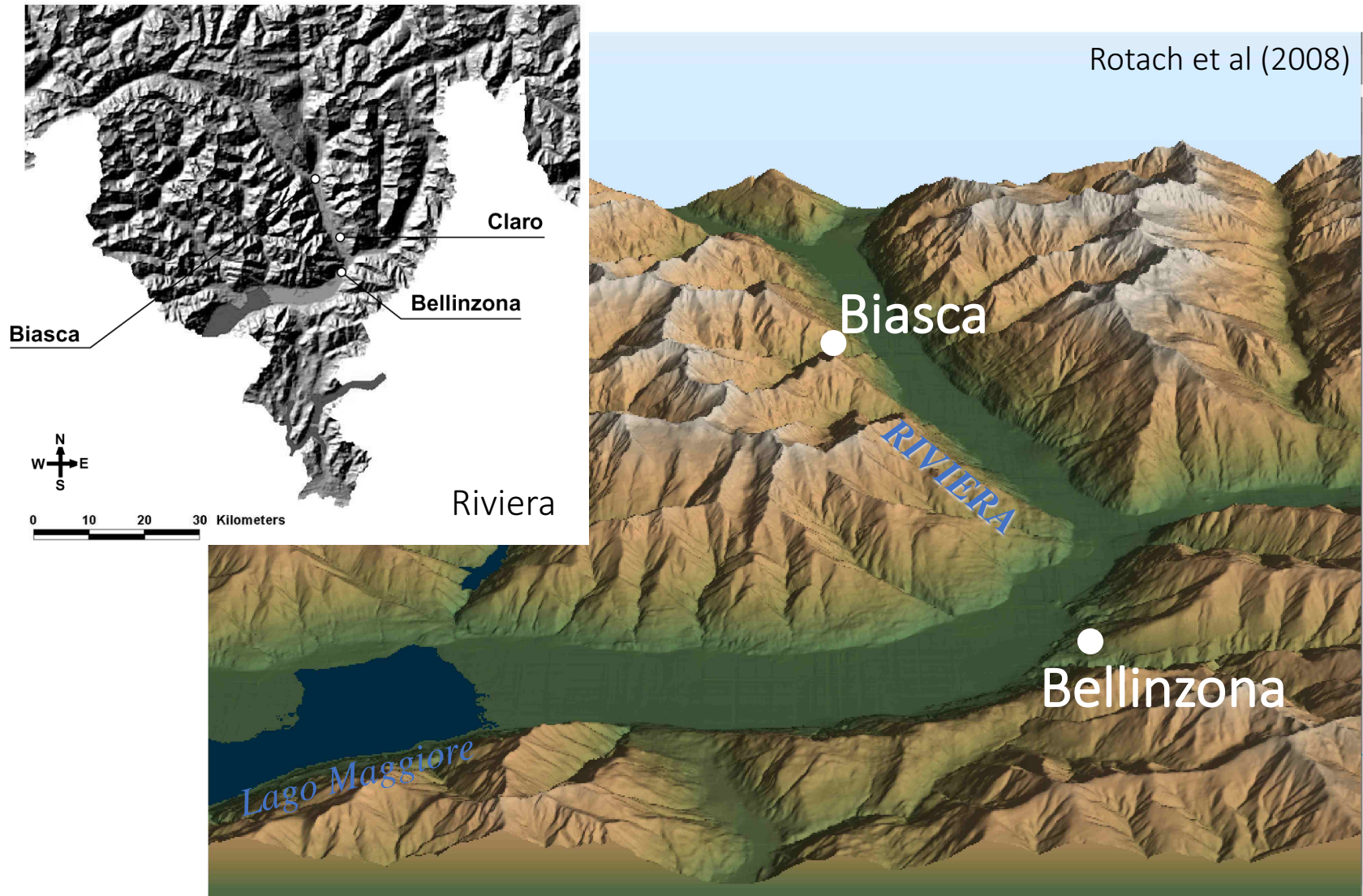
Buoyancy parameter: $\frac{g}{\theta}$

MAP Riviera Project

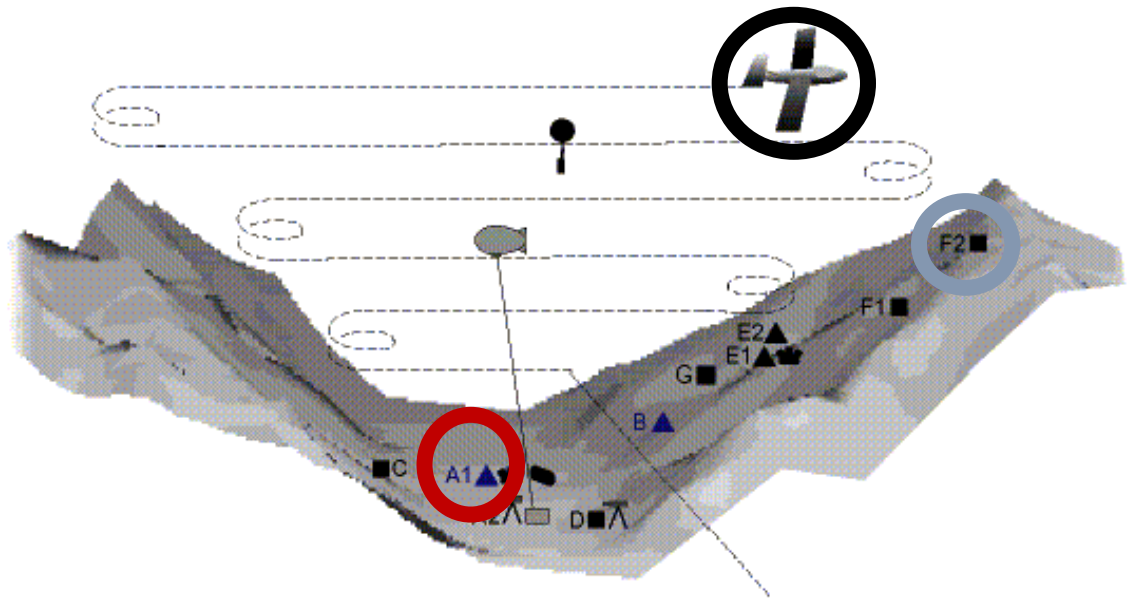
Rotach et al (2004)



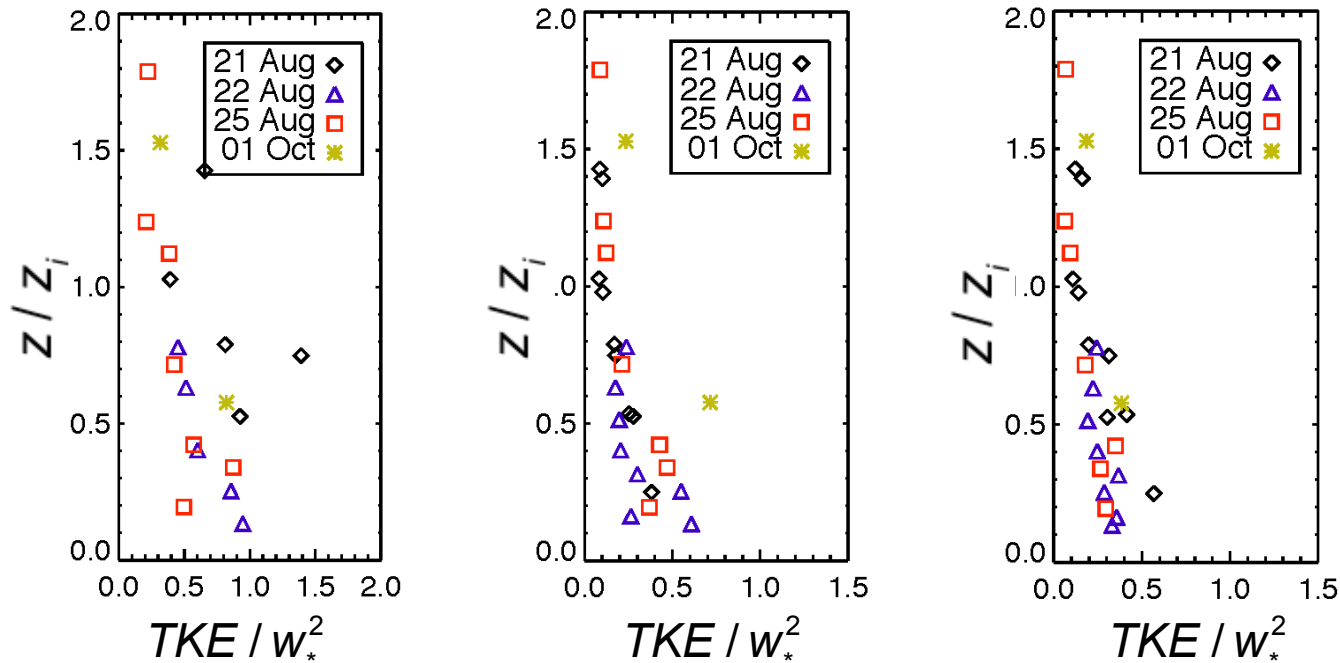
MAP Riviera Project



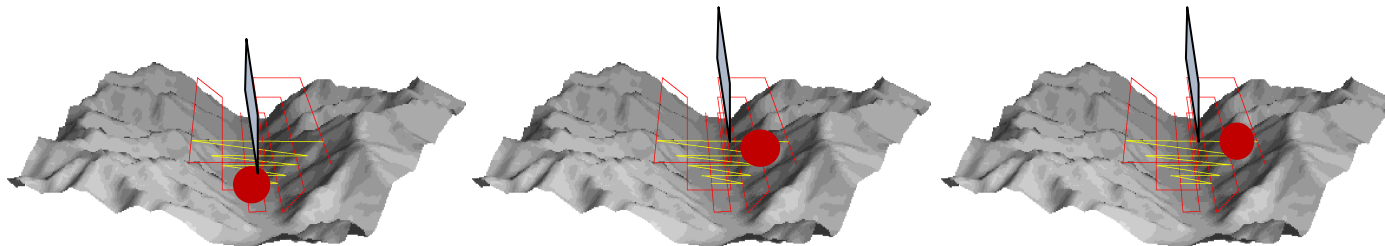
MAP Riviera Project: summer/fall 1999



MAP: Measured Scaled TKE profiles

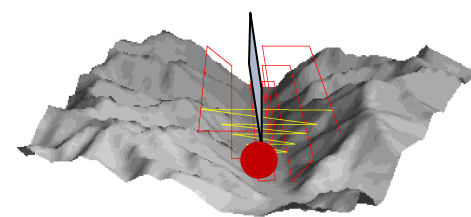
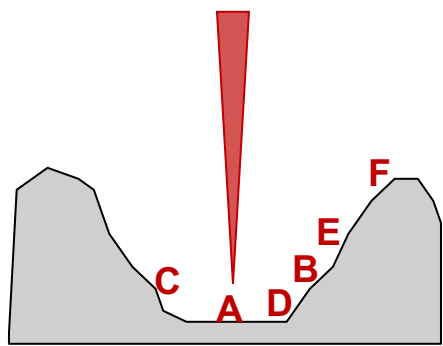
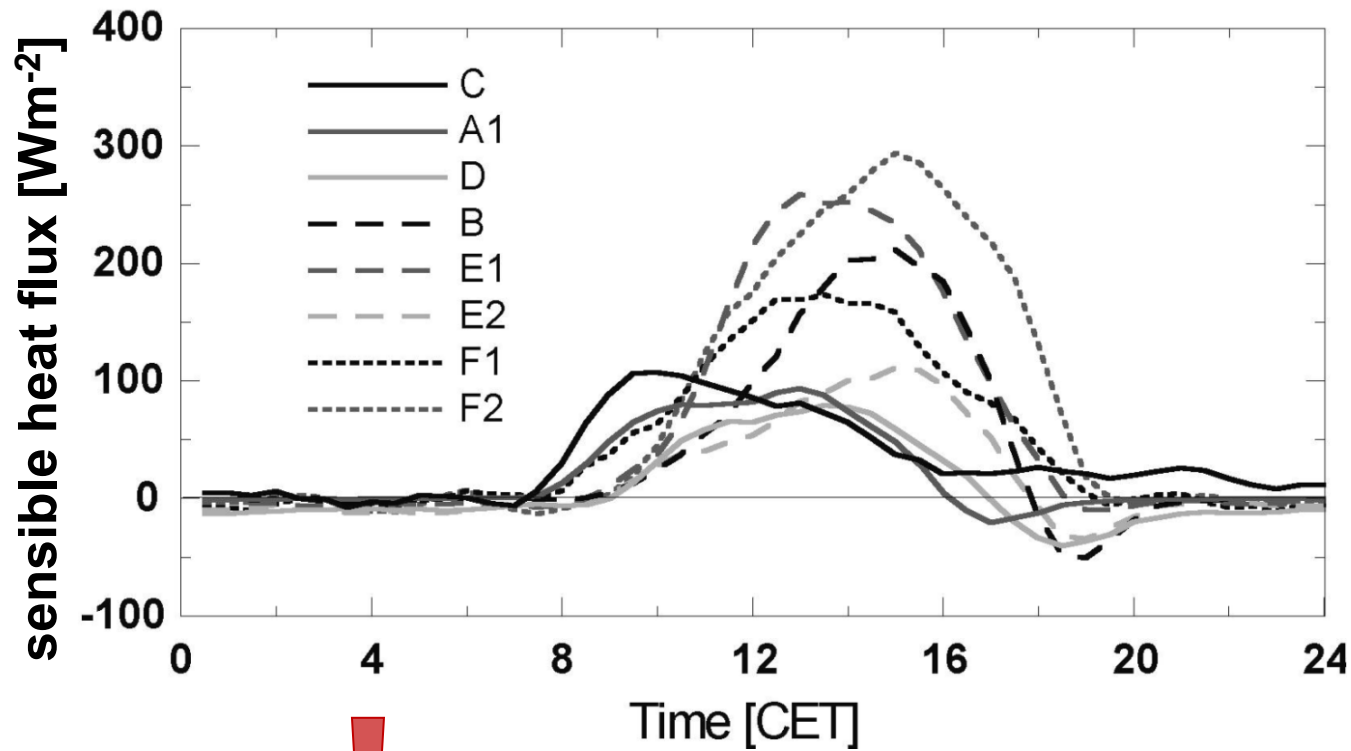


$$w_* = \left(\overline{w'\theta'_o} \frac{g}{\bar{\theta}} z_i \right)^{1/3}$$



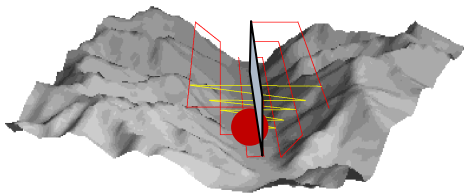
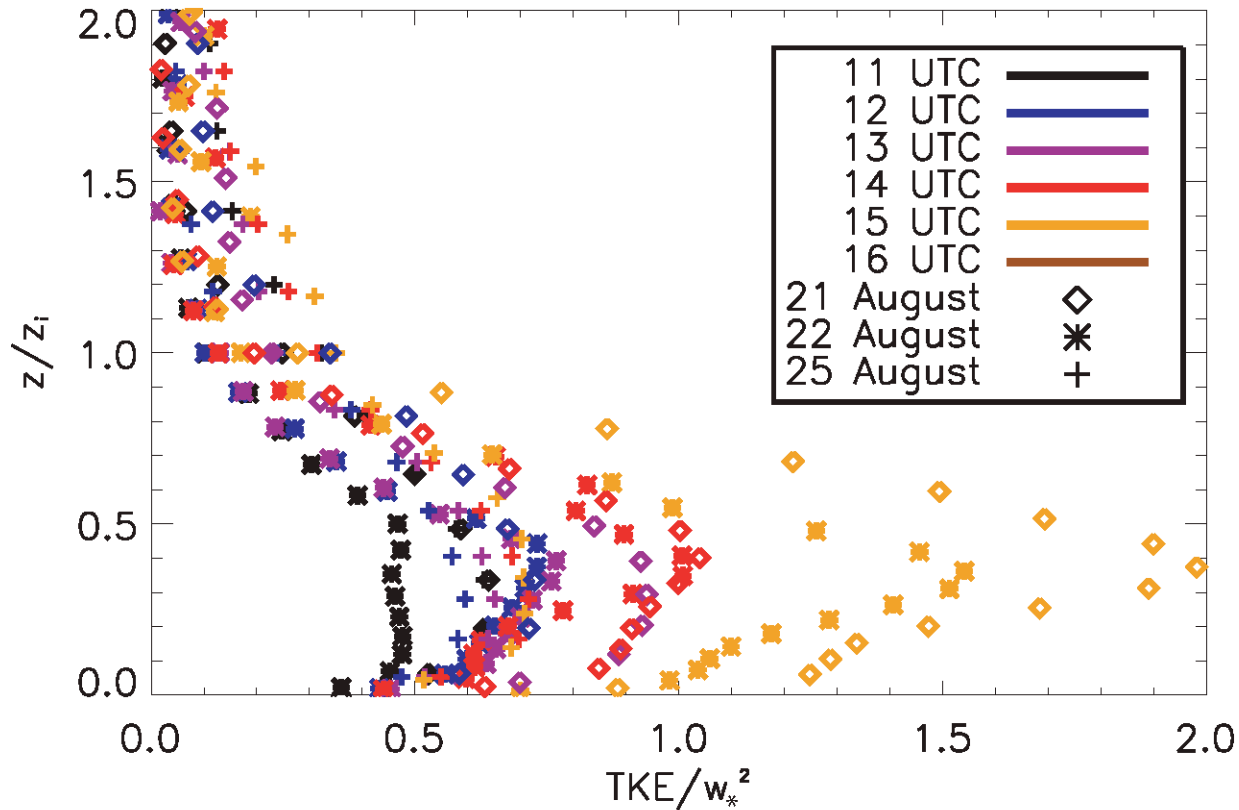
Weigel and Rotach (2004)

MAP: Measured surface fluxes



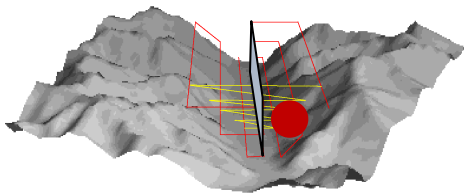
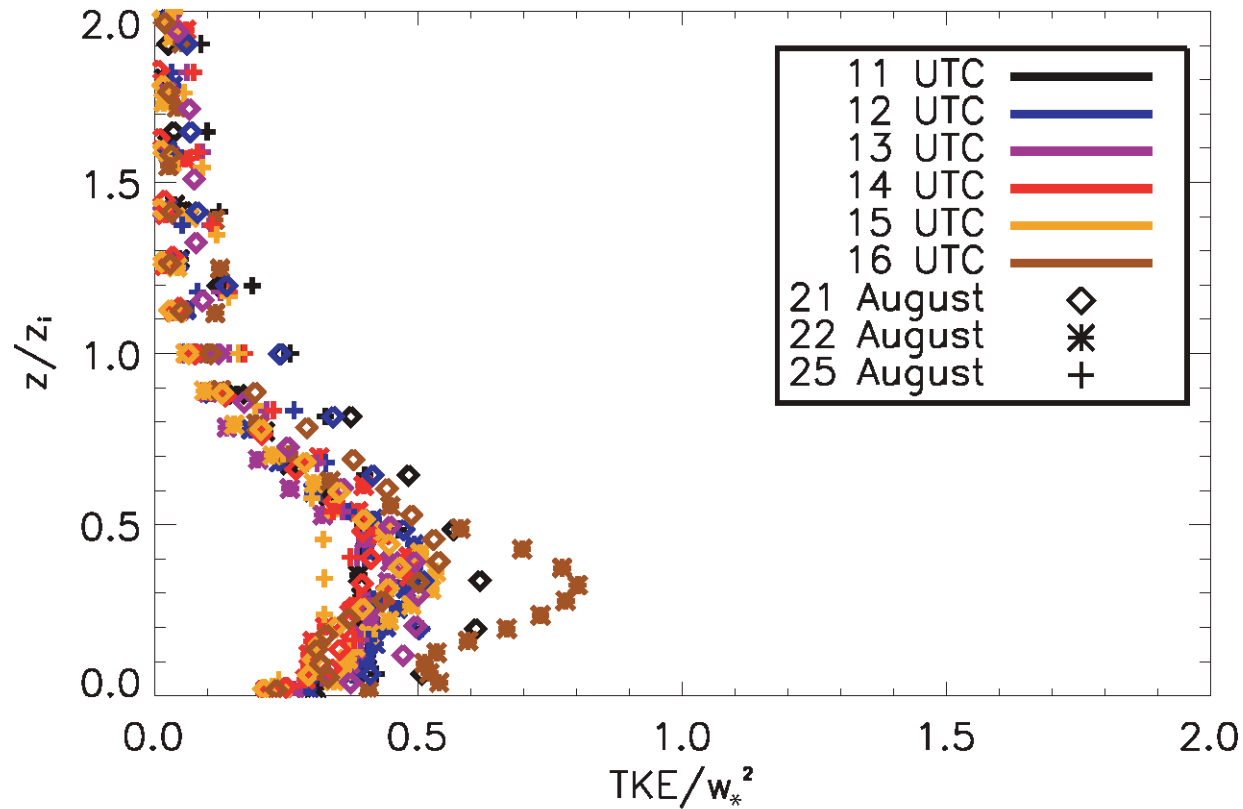
Rotach et al (2008)

MAP: Simulated Scaled TKE profiles



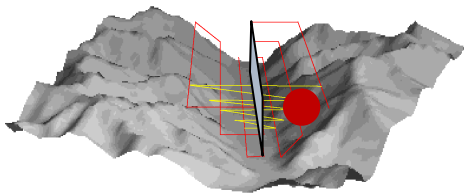
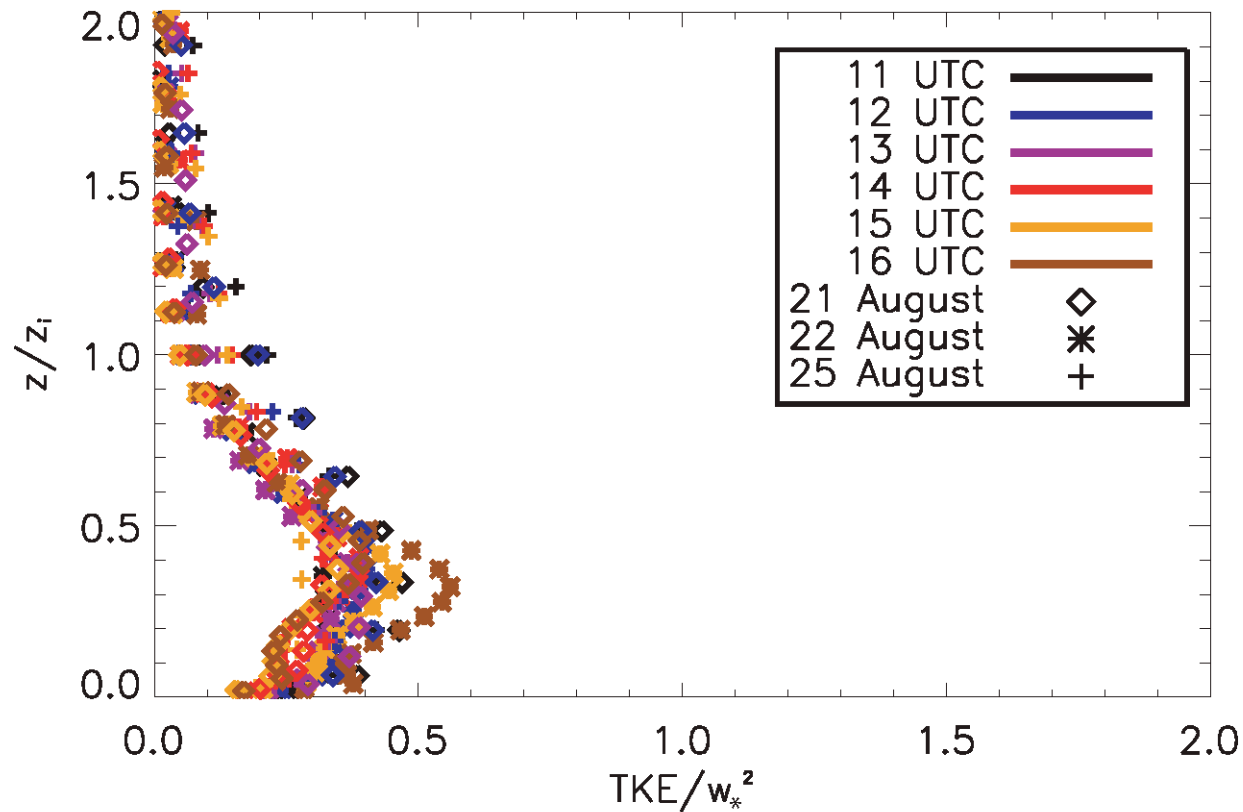
Weigel et al (BLM, 2007)

MAP: Simulated Scaled TKE profiles



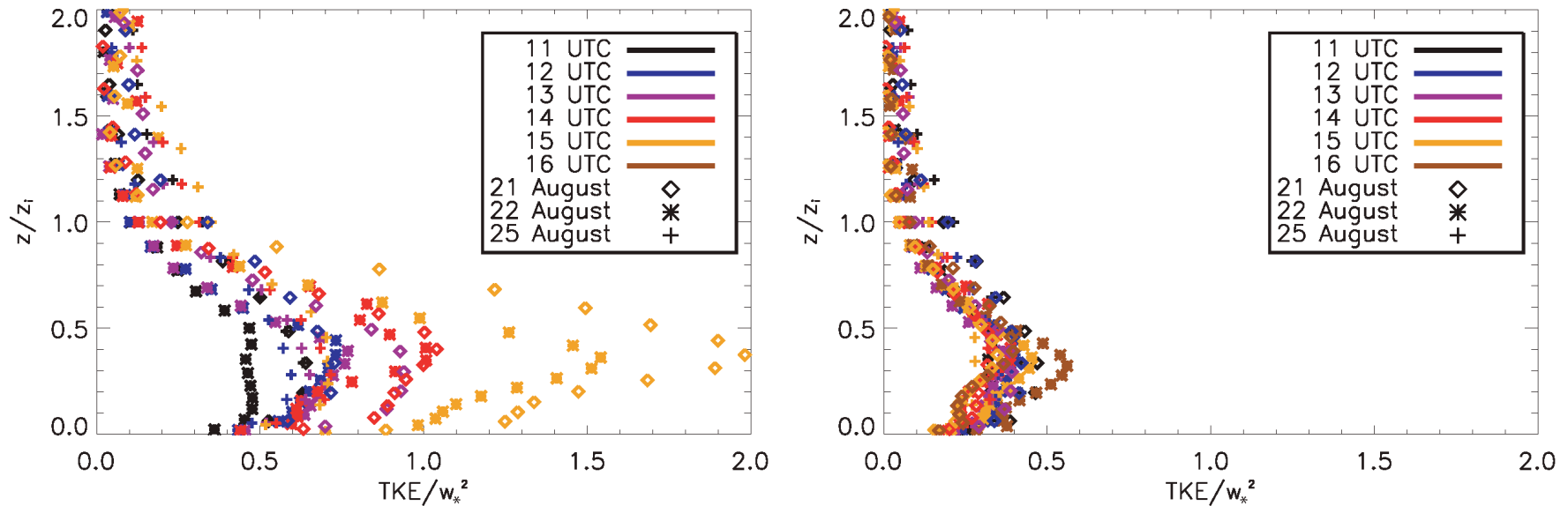
Weigel et al (BLM, 2007)

MAP: Simulated Scaled TKE profiles



Weigel et al (BLM, 2007)

MAP: Simulated Scaled TKE profiles



How can this scaling behavior be understood?
→ Look into simulated TKE budget

MAP: TKE production mechanisms

TKE budget

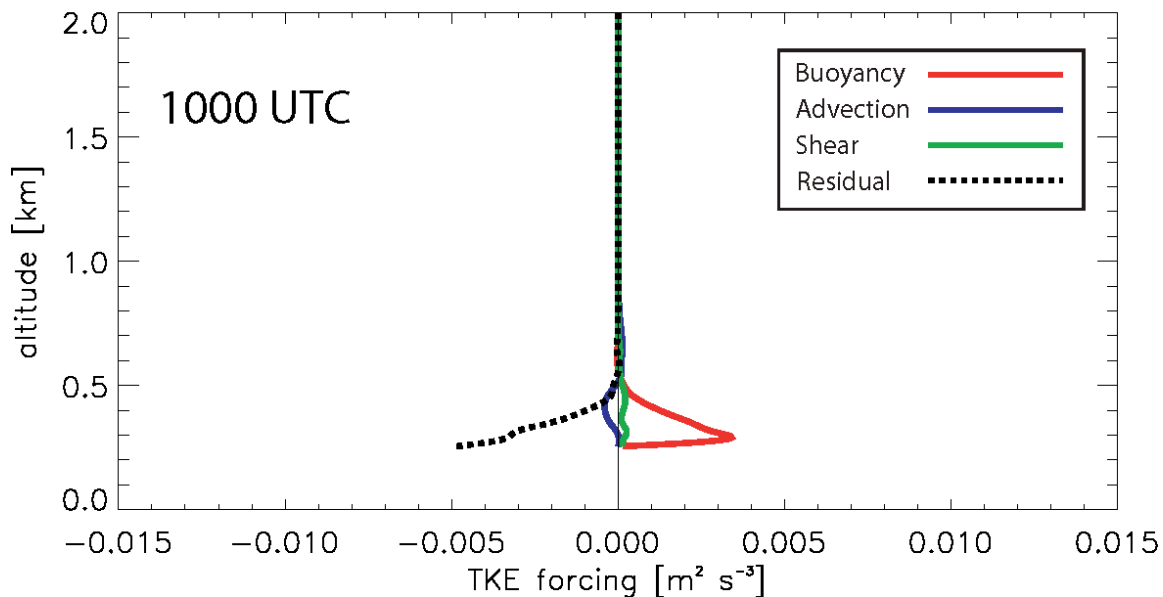
(solved in ARPS's 1.5 order TKE closure)

$$\begin{aligned}
 \frac{\partial \bar{e}}{\partial t} = & \underbrace{-\bar{u} \left(\frac{\partial \bar{e}}{\partial x} \right) - \bar{v} \left(\frac{\partial \bar{e}}{\partial y} \right) - \bar{w} \left(\frac{\partial \bar{e}}{\partial z} \right)}_{\text{ADVECTION}} - \underbrace{(\overline{u'w'}) \left(\frac{\partial \bar{u}}{\partial z} \right) - (\overline{v'w'}) \left(\frac{\partial \bar{v}}{\partial z} \right)}_{\text{SHEAR PRODUCTION}} + \\
 & + \underbrace{\frac{g}{\theta} (\overline{w'\theta'})}_{\text{BUOYANCY PRODUCTION}} - \underbrace{\frac{1}{\rho} \frac{\partial}{\partial z} (\overline{w'p'}) - \frac{\partial}{\partial z} (\overline{ew'}) - \epsilon}_{\text{PRESSURE CORRELATION, TURB. TRANSPORT, DISSIPATION}}.
 \end{aligned}$$

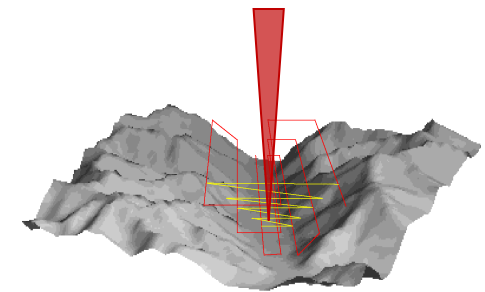
MAP: TKE production mechanisms

Evaluation of the TKE budget equation with ARPS

Late morning



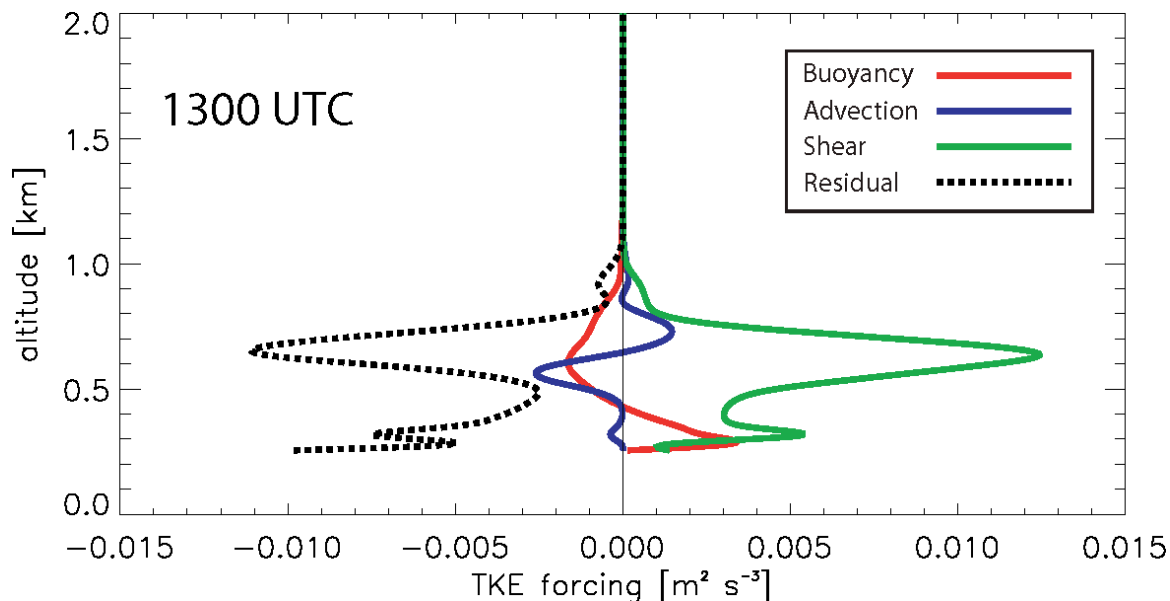
Example 22 August



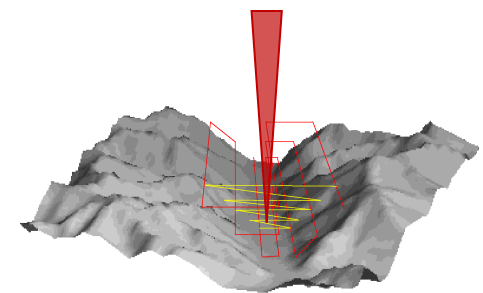
MAP: TKE production mechanisms

Evaluation of the TKE budget equation with ARPS

Early Afternoon

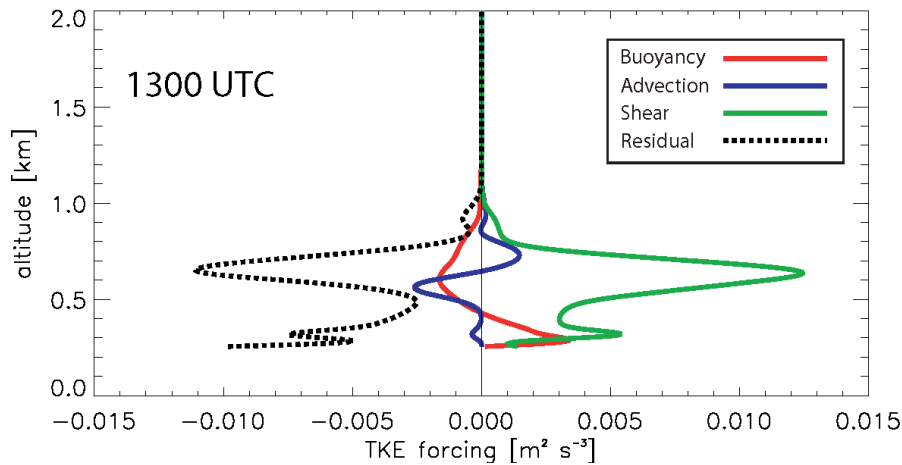


Example 22 August

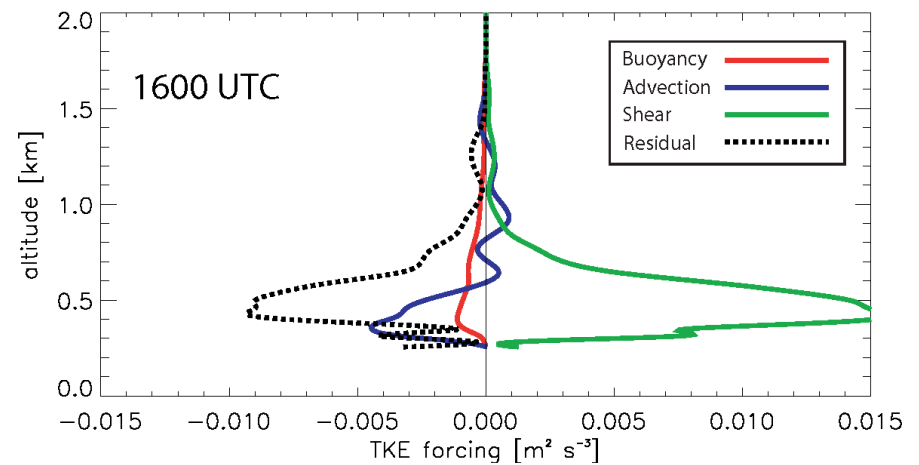


MAP: TKE production mechanisms

Early Afternoon



Late Afternoon



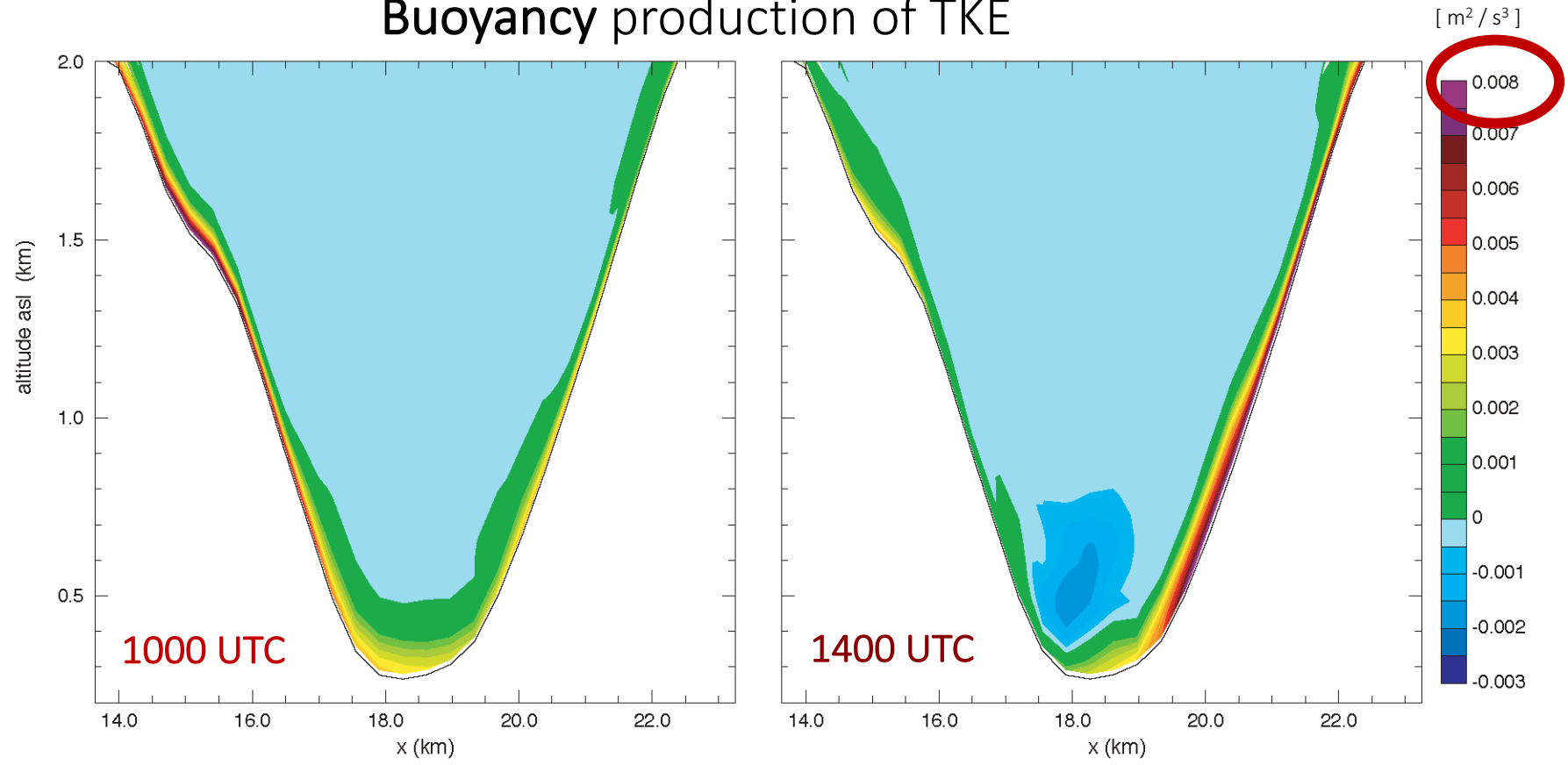
- **Shear** is dominant production mechanism

Turbulence is determined:

- by interaction of thermally driven valley flows
- **not** by **buoyant processes** from heated surface

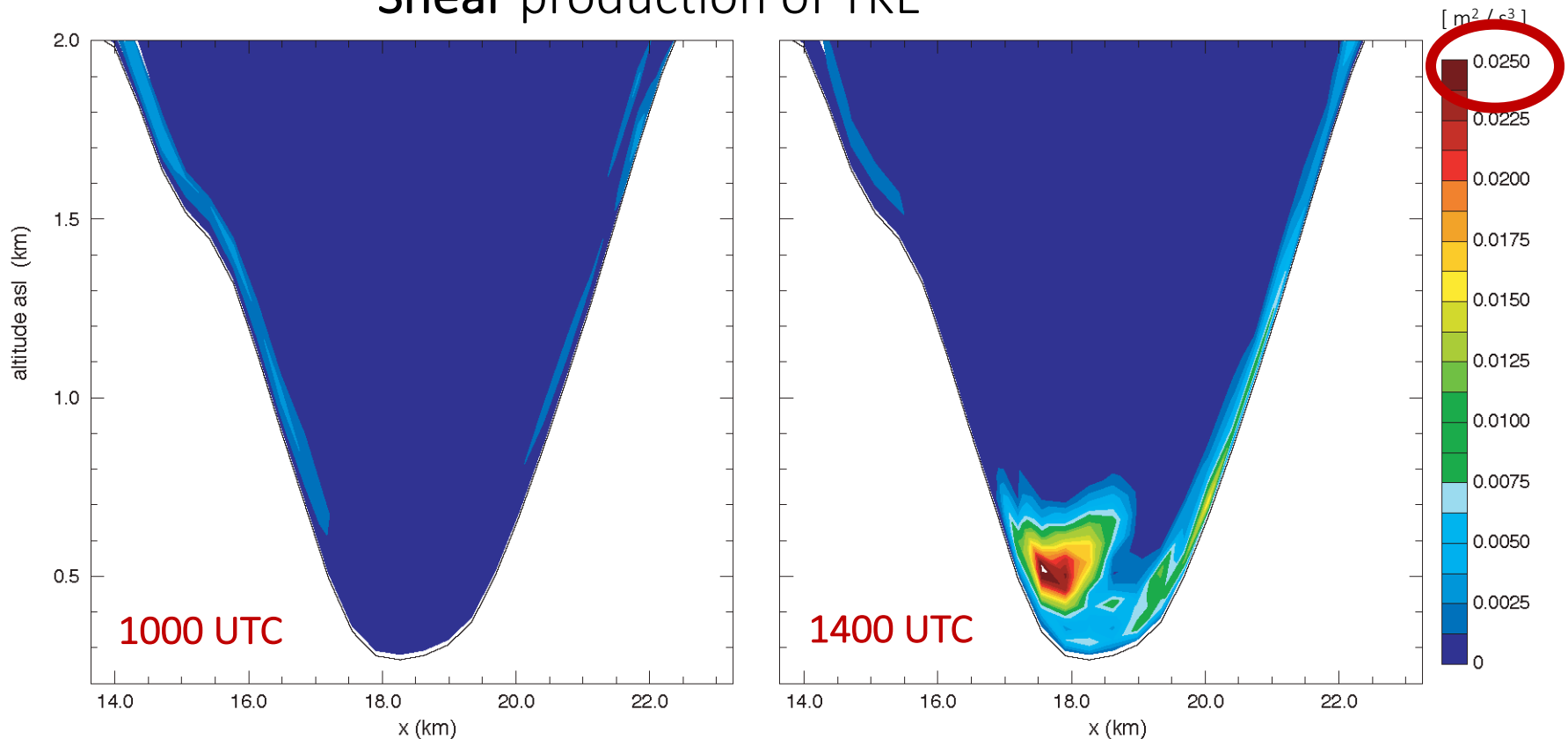
MAP: TKE production mechanisms

Buoyancy production of TKE



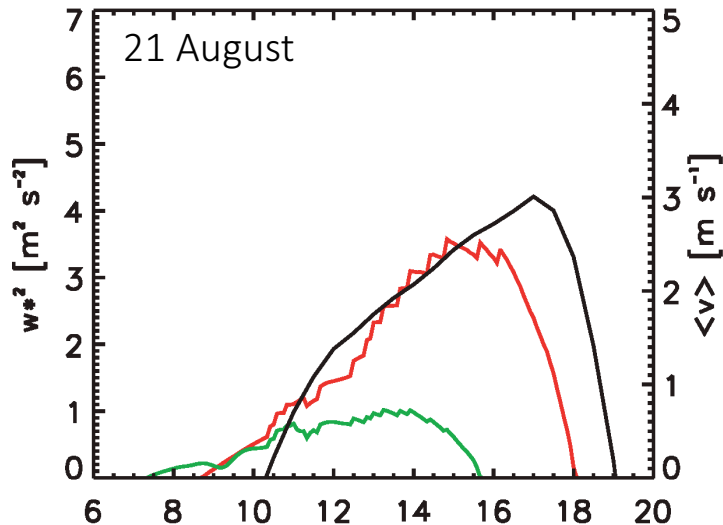
MAP: TKE production mechanisms

Shear production of TKE

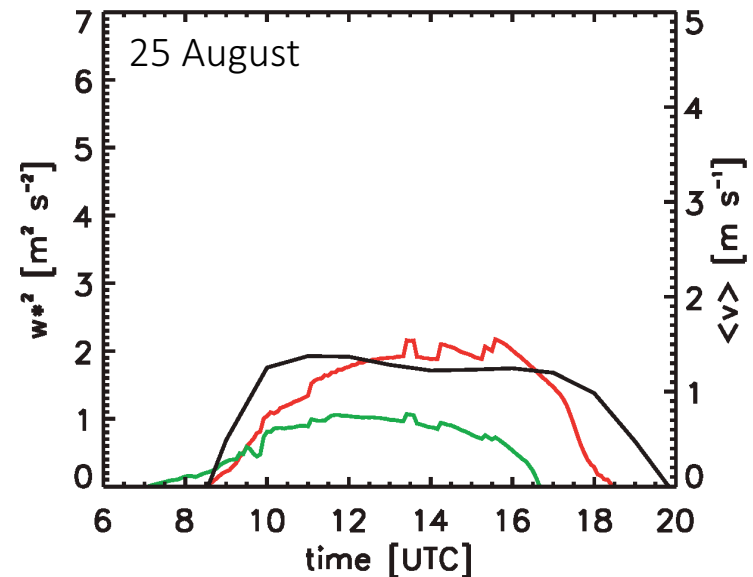
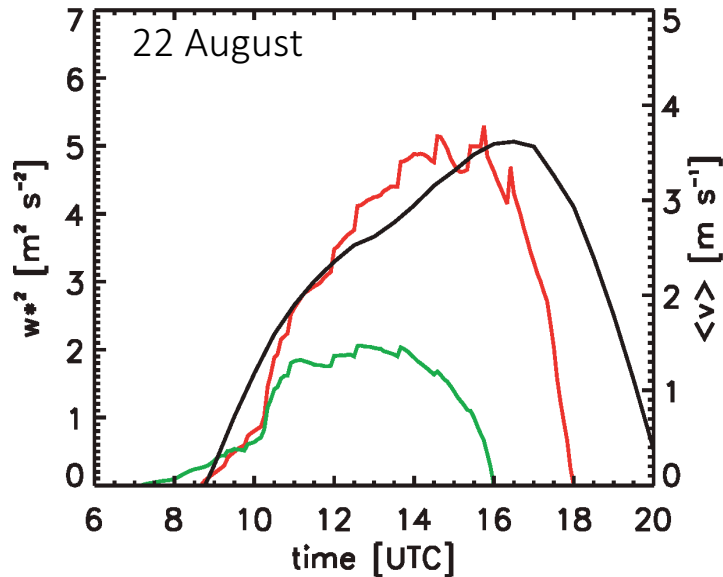


- TKE determined by strength of valley winds.
- But why does convective scaling approach still work ?

MAP: TKE production mechanisms



- up-valley wind speed $\langle v \rangle$
(averaged over valley volume)
- w_*^2 on eastern (sunlit) slope
- w_*^2 on valley floor

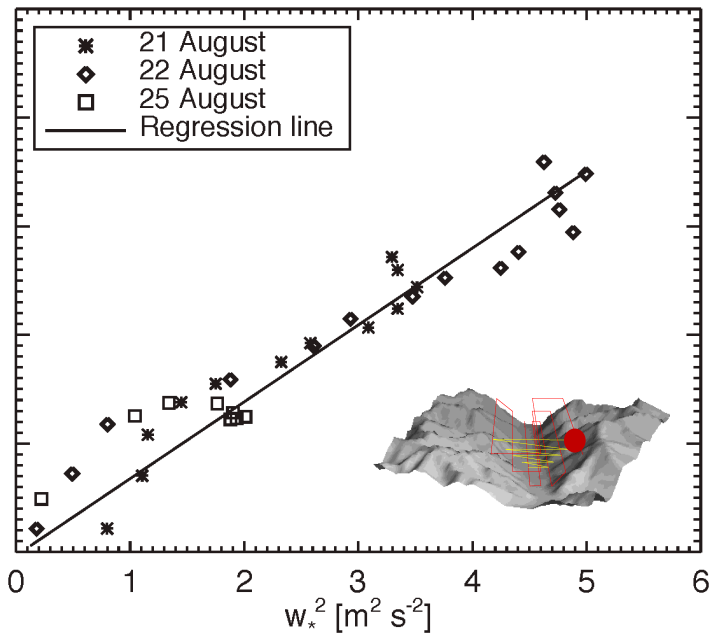
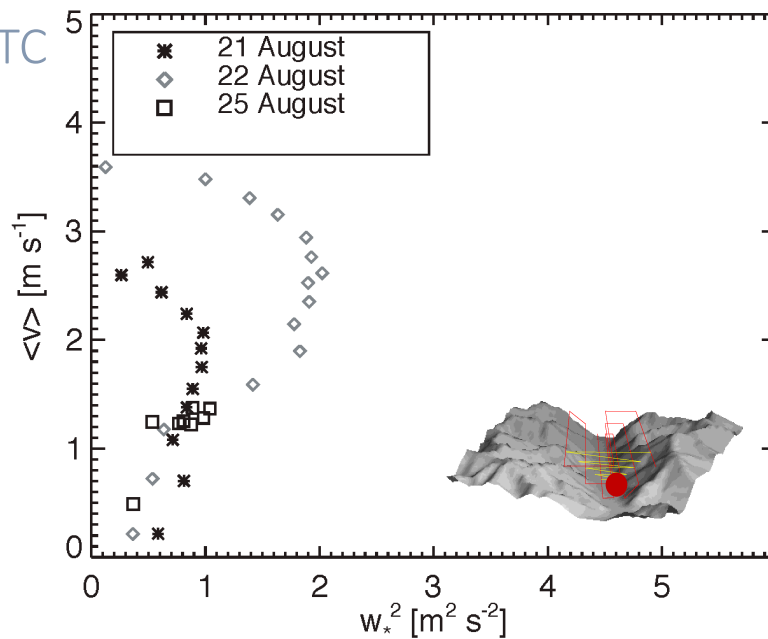


MAP: TKE production mechanisms

Correlation between $\langle v \rangle$ and w_*^2 $\langle v \rangle \sim w_*^2$

w_*^2 from valley floor

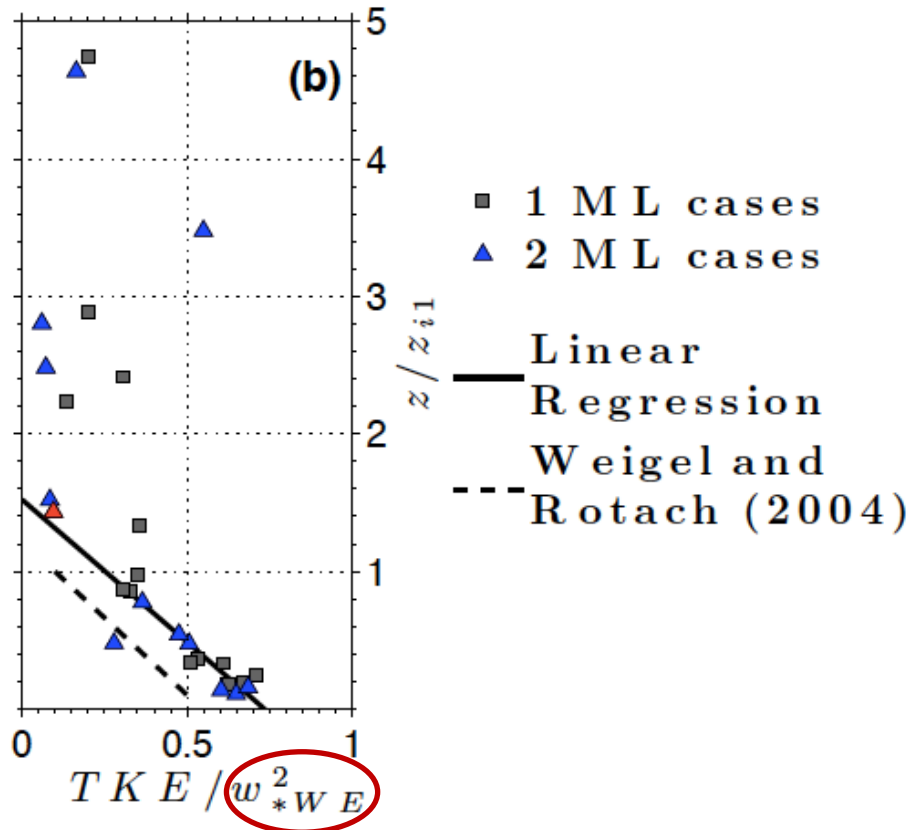
w_*^2 from sunlit slope



- Coincidence?
- Problem with scaling: $\langle v \rangle$ has wrong dimensions...

How general is this TKE scaling?

Inn Valley, 2013



- east-west valley
- also 'optimal' w_*
- similar (but not equal) behaviour

Baur, MSc thesis (2015)

Scaling in complex terrain

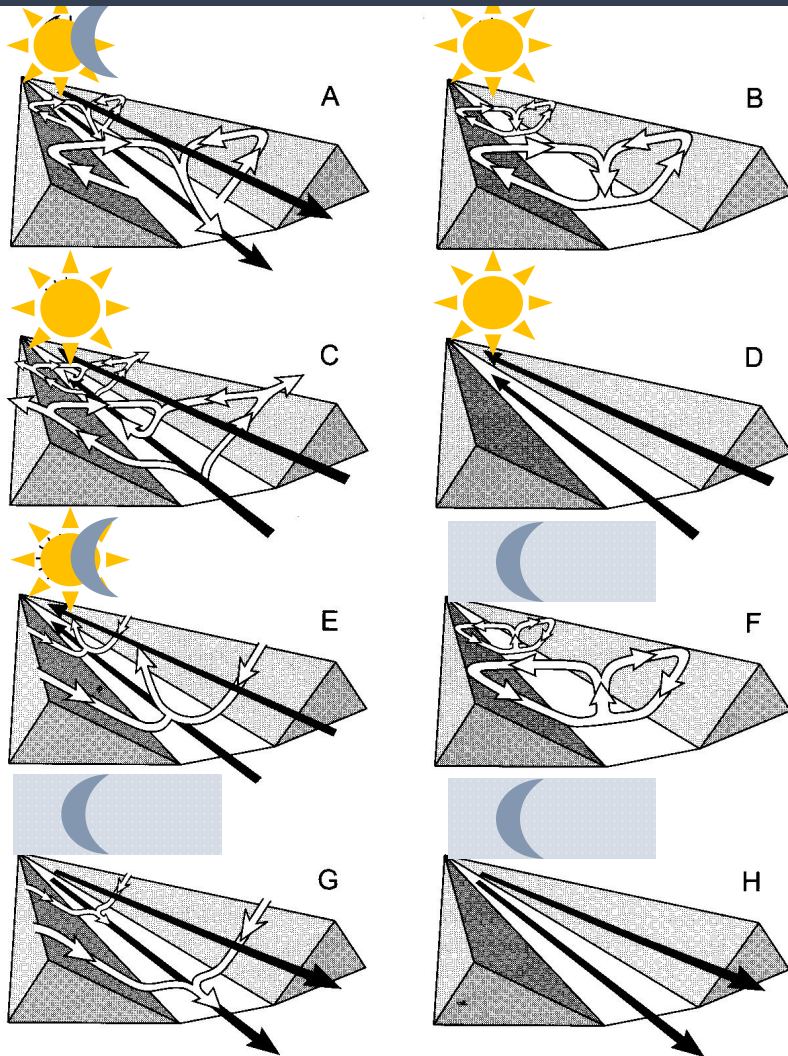
Application Examples:

1. Mixed-layer TKE scaling (MAP Riviera)
2. Wind directional change with height
3. Local scaling flux-variance relationships (12 Datasets)

→ friction velocity $u_* =: \left[\overline{(u' w'_o)^2} + \overline{(v' w'_o)^2} \right]^{1/4}$

→ if coordinate system aligned with mean wind and no directional shear with height: $\overline{v' w'_o} = 0$

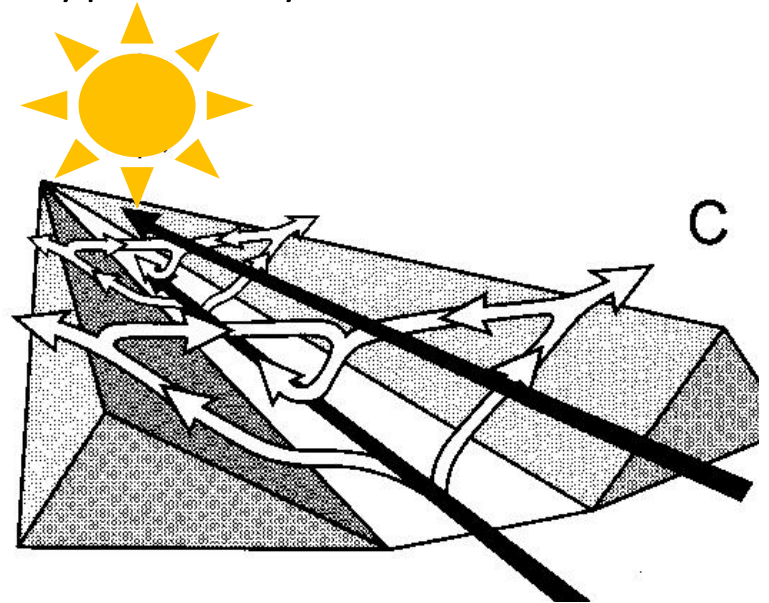
Valley wind – Slope wind



⇨ Hangwind

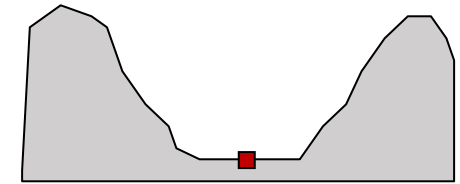
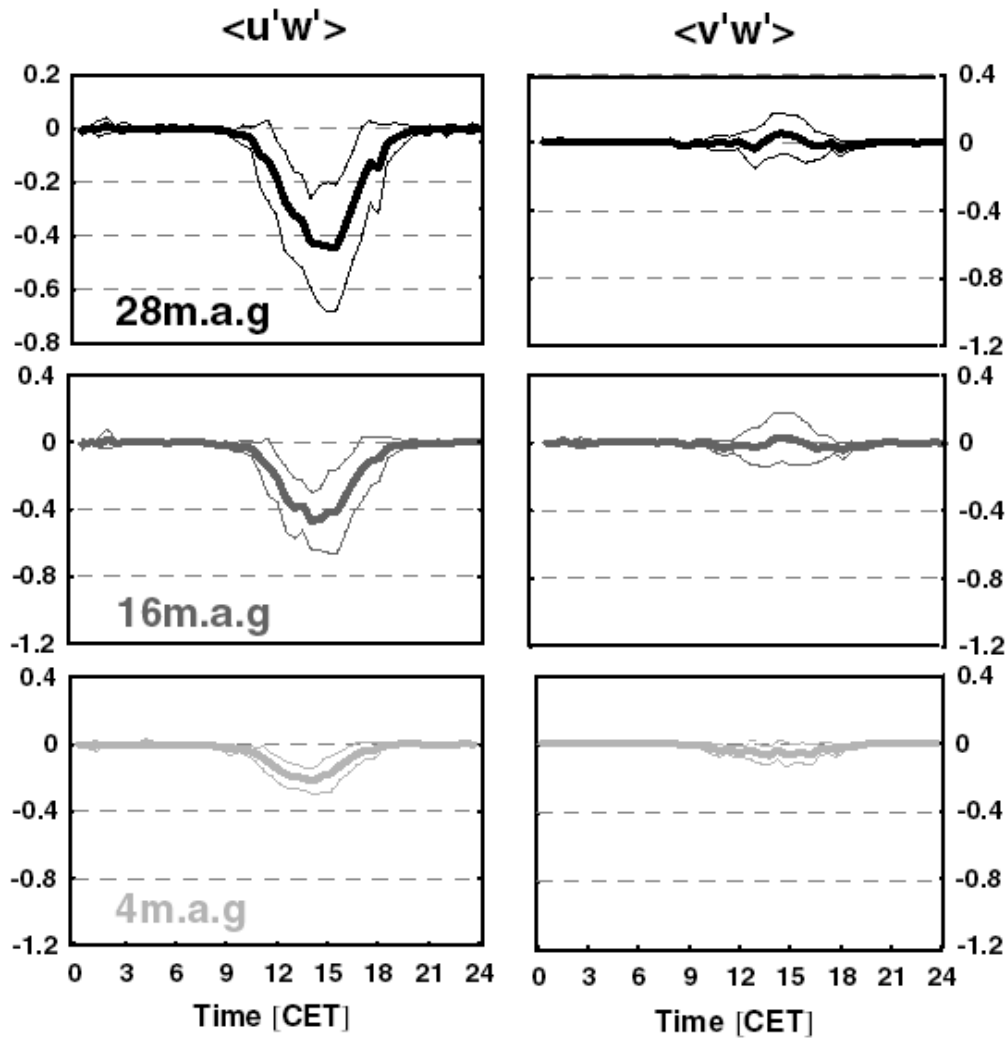
⇨ Berg-Talwind

typical daytime



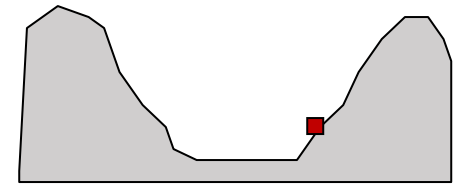
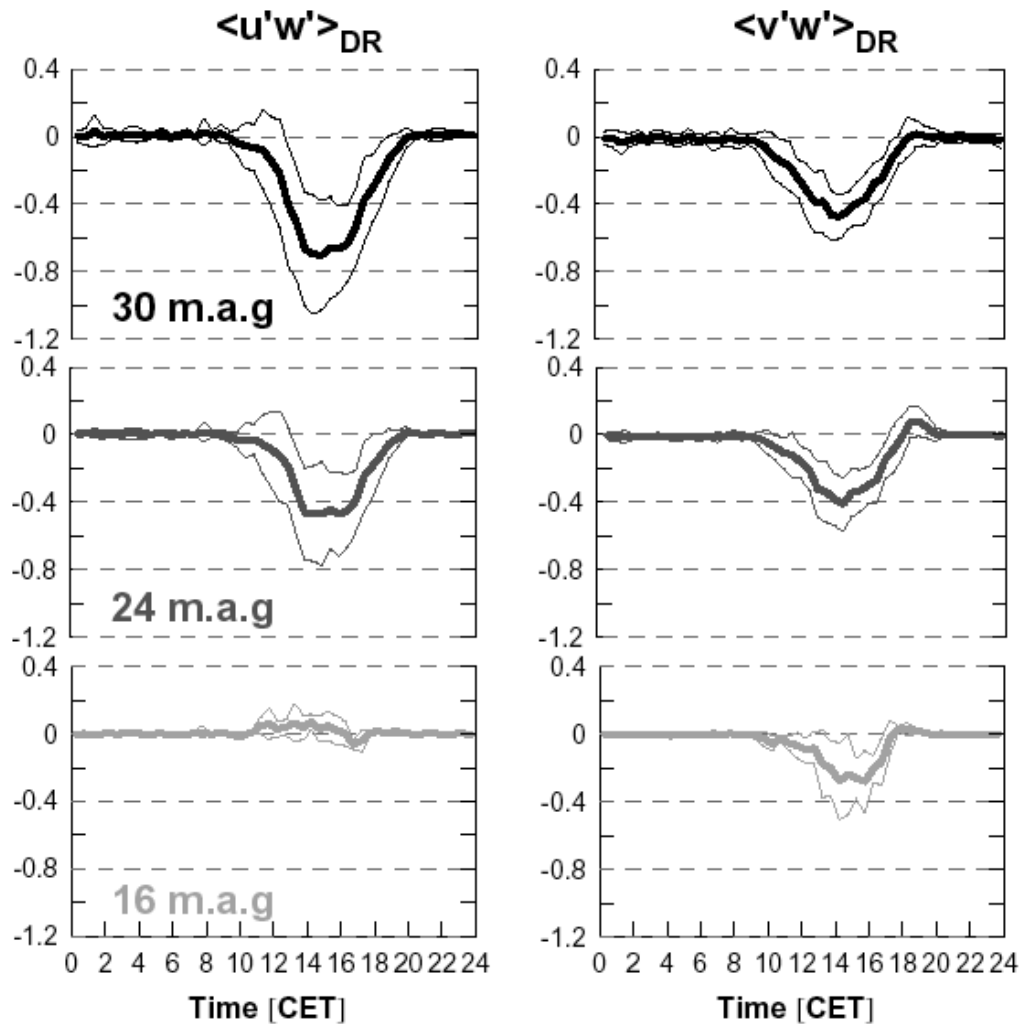
after Defant, 1949

Momentum transport: Valley floor



frictional stress: $\overline{u'w'}$

Momentum transport: Slope station



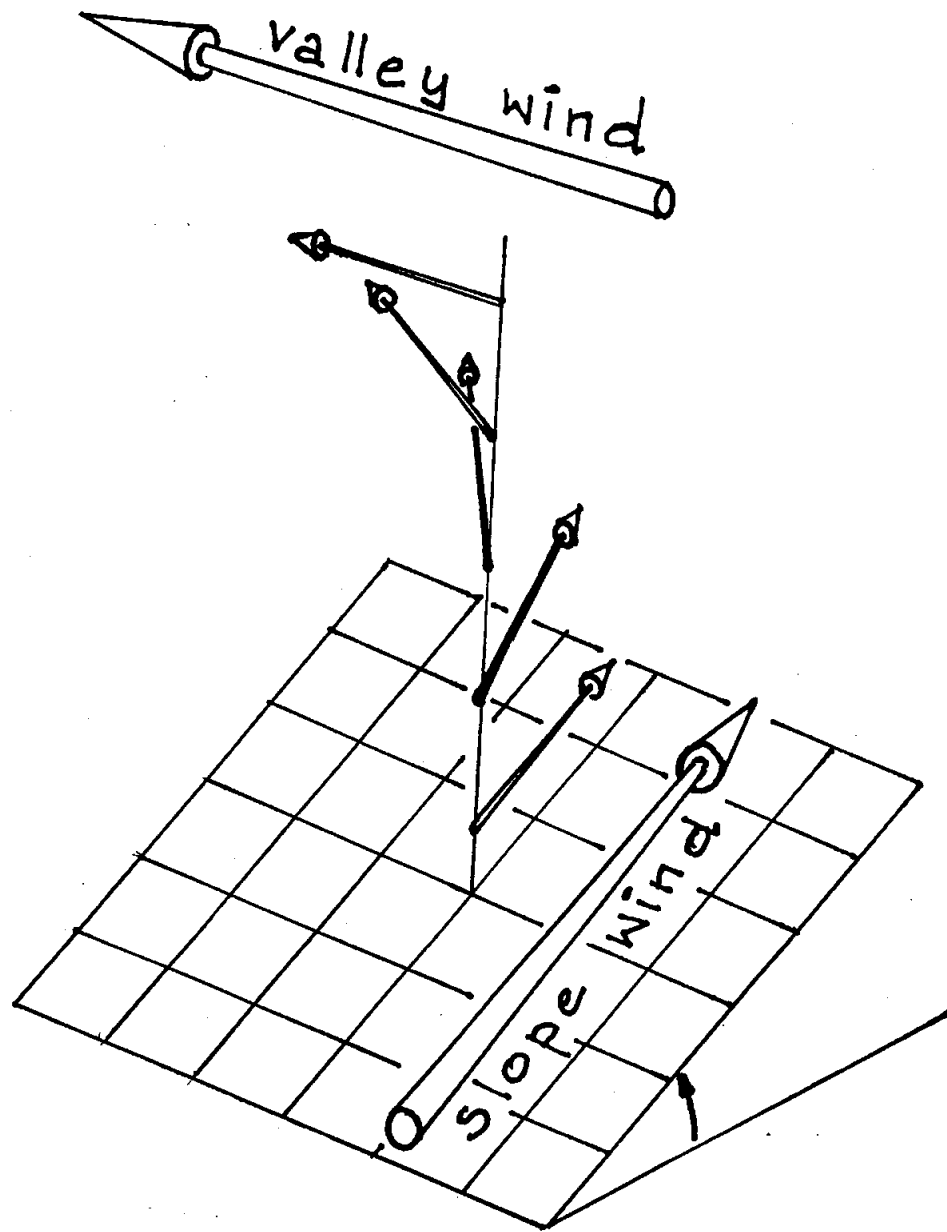
Local circulations cause :

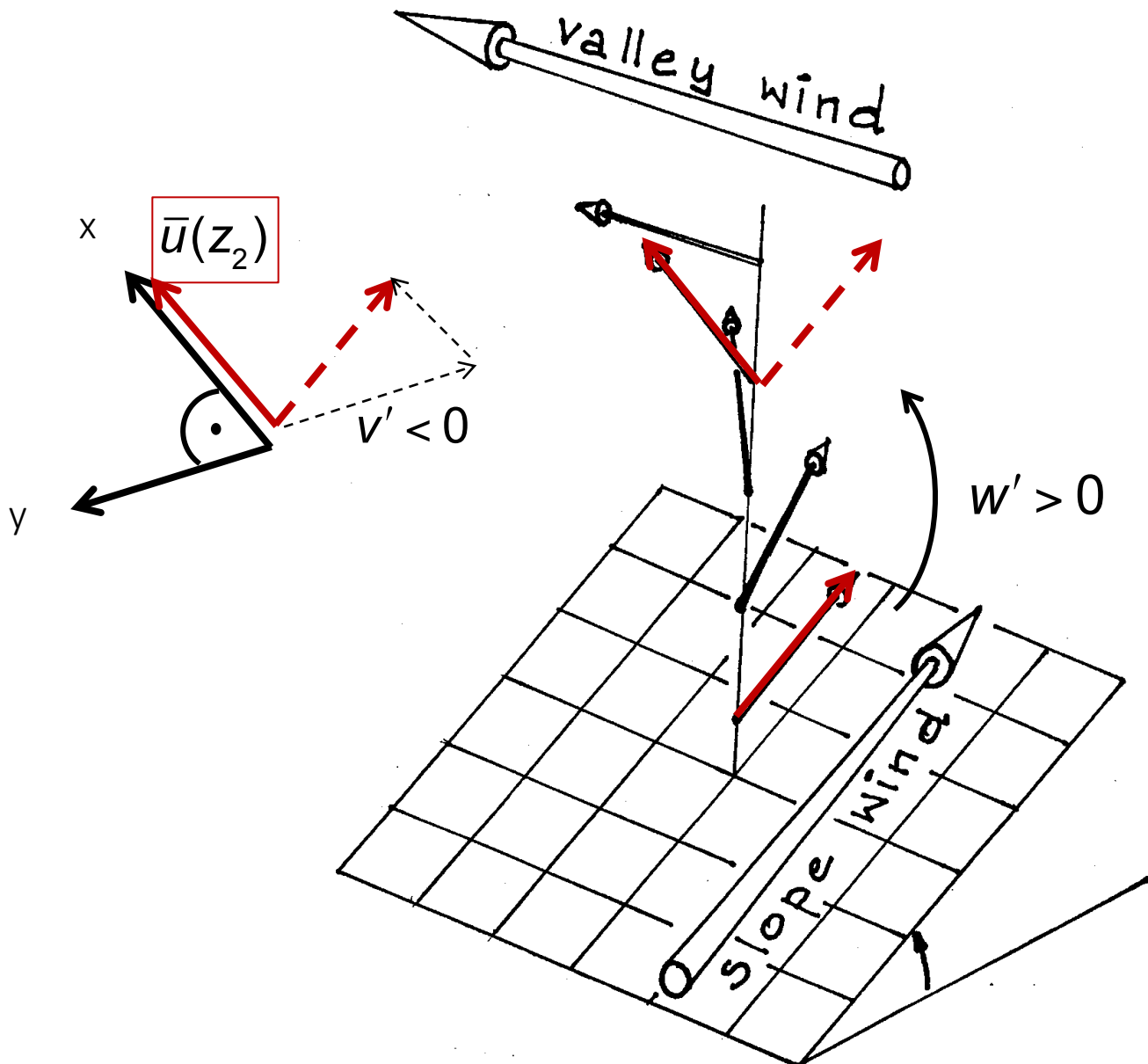
frictional stress: $\overline{u'w'}$

+

directional stress: $\overline{v'w'}$

Andetta et al. (2001)

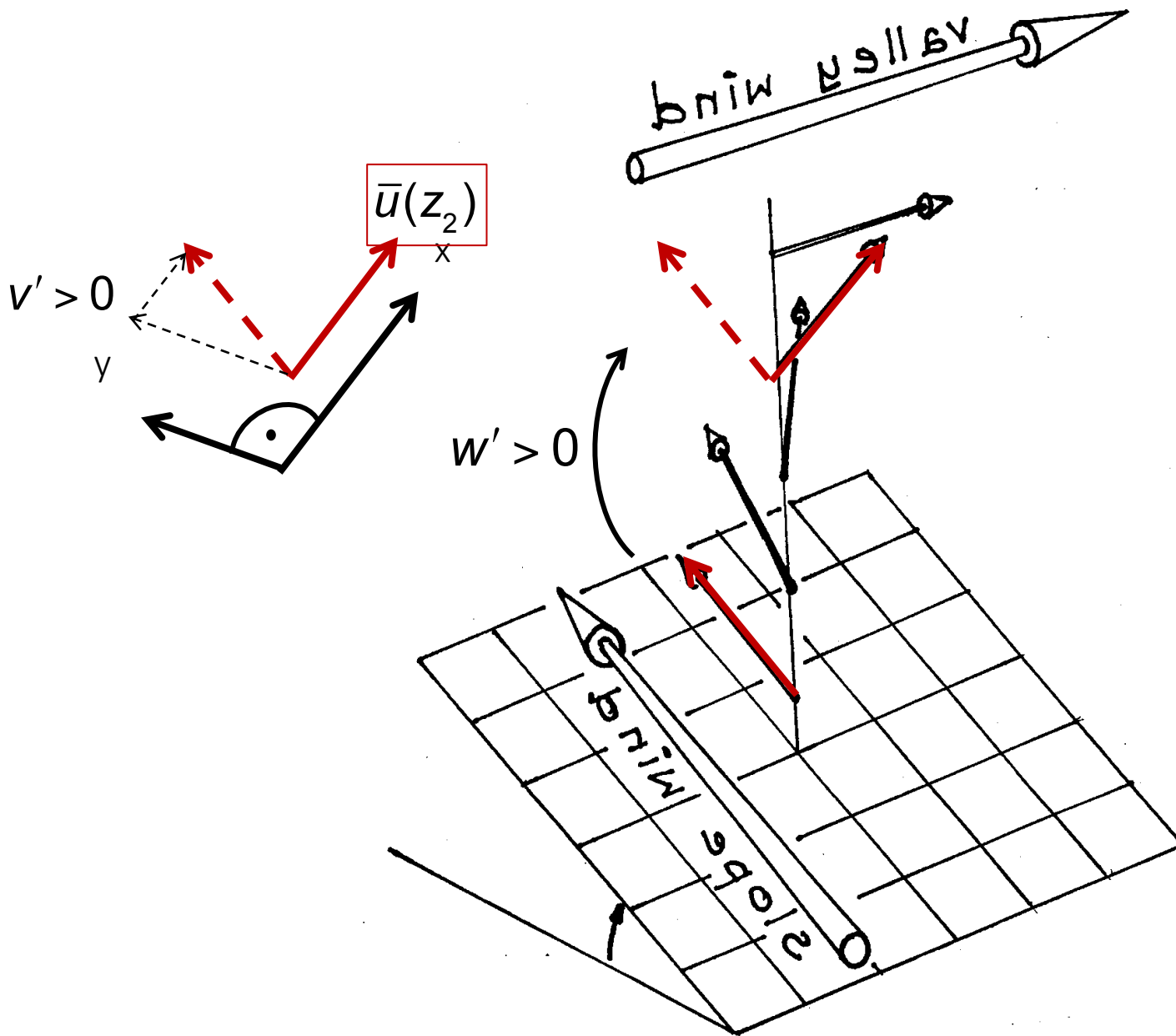




this eddy:
 $v'w' < 0$

many eddies:

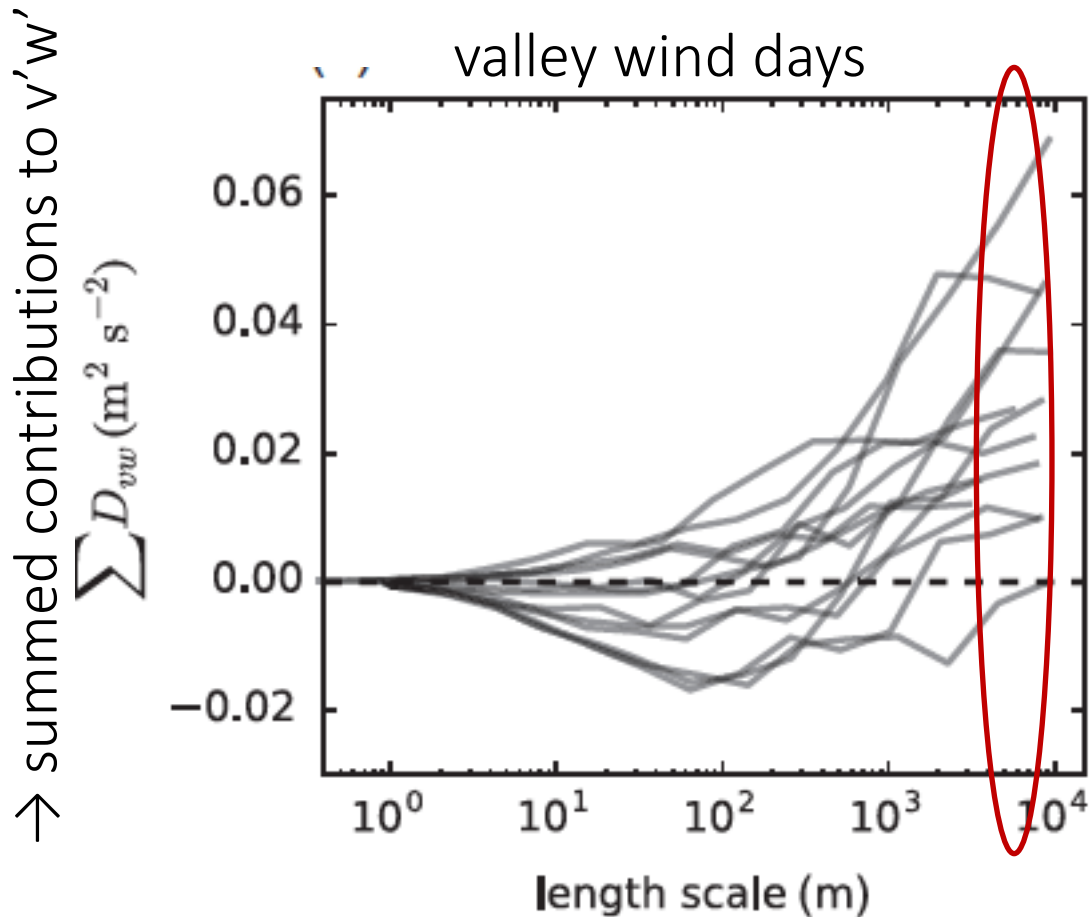
$$\overline{v'w'} < 0$$



other slope:

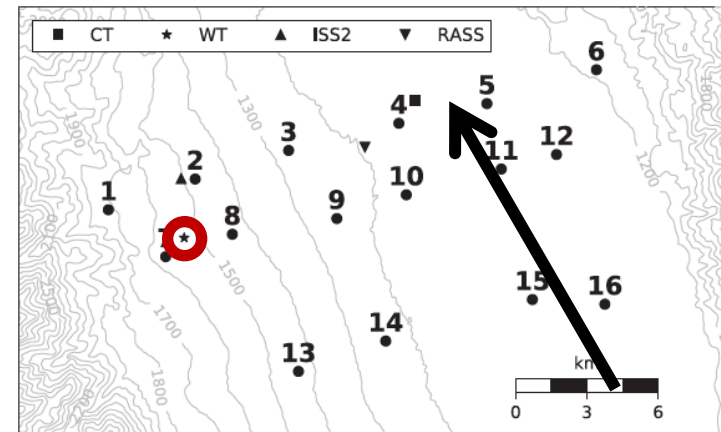
$$\overline{v'w'} > 0$$

Momentum transport: Opposite slope



→ $\langle v'w' \rangle$ all scales

Owens Valley, CA



Babic N et al. (2017)

Momentum transport: Scaling velocity

flat terrain:

→ in streamline coordinates

$$u_* = \left(-\overline{u'w'} \right)^{1/2} \quad \text{is enough, because} \quad \overline{v'w'} = 0$$

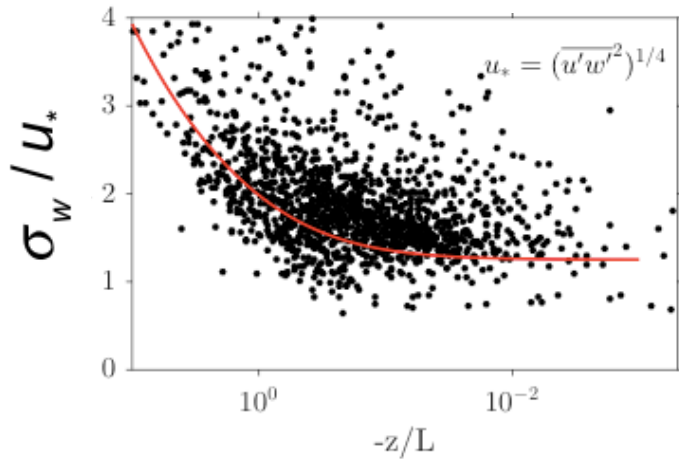
systematic turning with height (e.g., slope/valley wind)

→ even if streamline coordinates

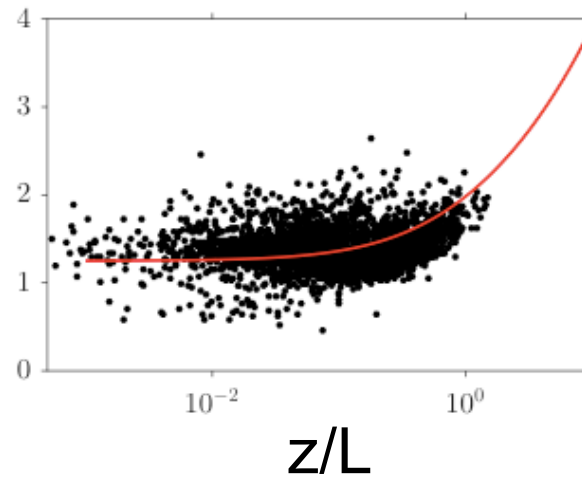
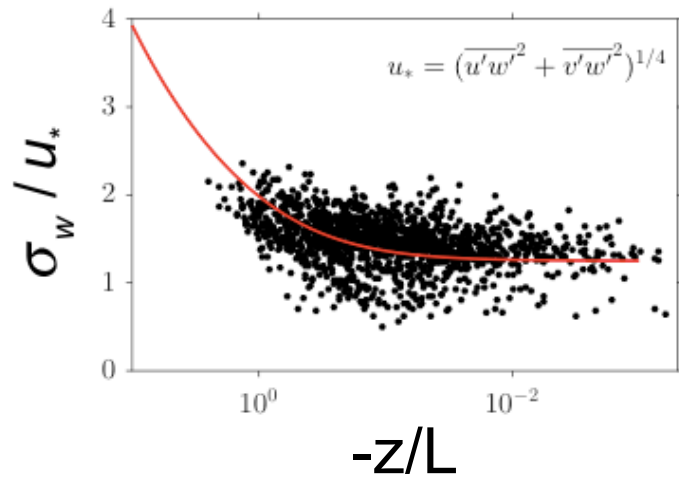
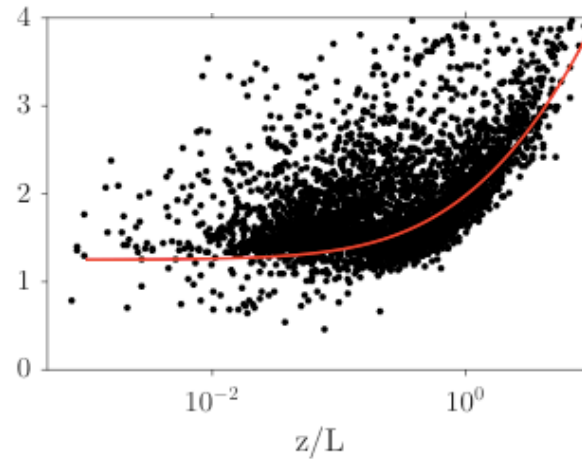
$$u_* = \left(\overline{u'w'}^2 + \overline{v'w'}^2 \right)^{1/4} \quad \text{necessary}$$

Momentum transport: Scaling velocity

unstable



stable



data from i-Box site ‚Hochhäuser‘, 2014-2016, HQ

Scaling in complex terrain

Application Examples:

1. Mixed-layer TKE scaling (MAP Riviera)
2. Wind directional change with height
3. Local scaling flux-variance relationships (12 Datasets)

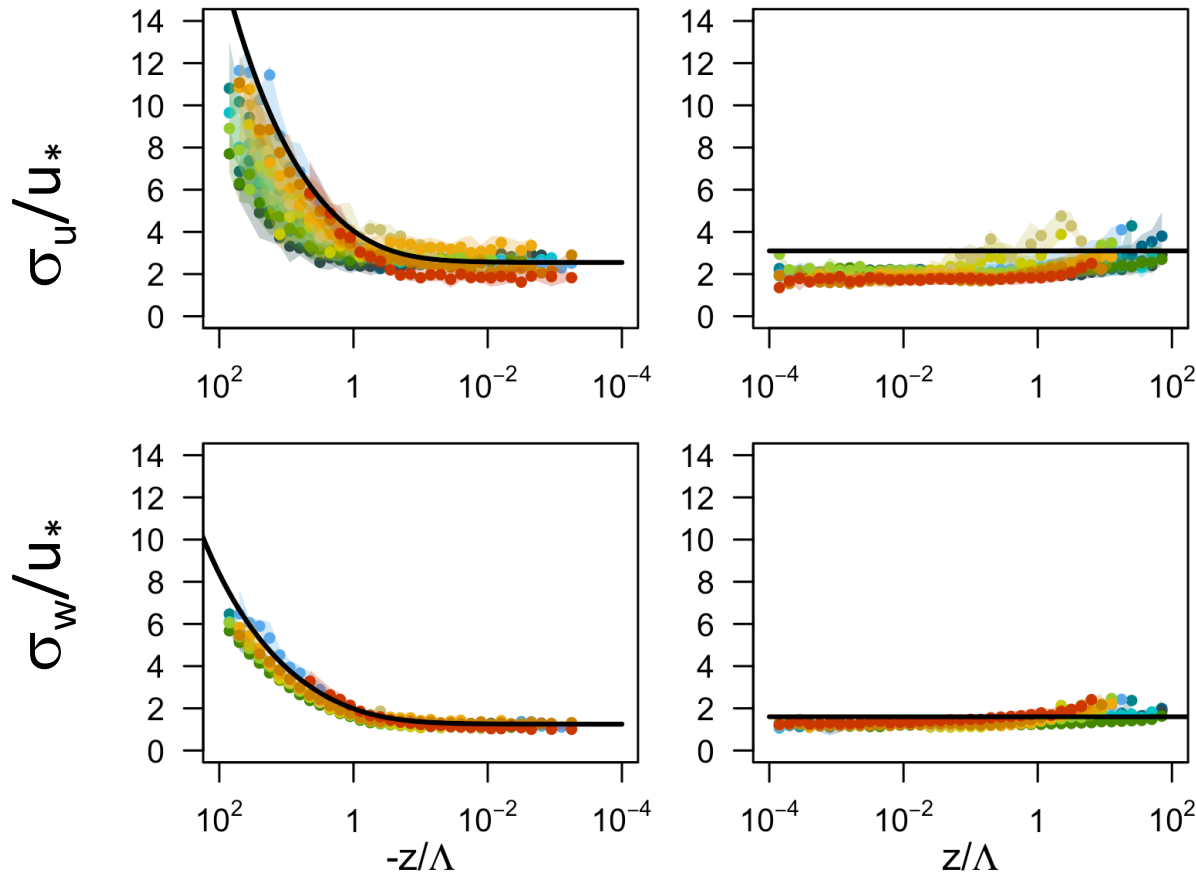
→ look at influence of anisotropy

Scaling in complex terrain

Standard deviation horizontal and vertical velocity

Unstable

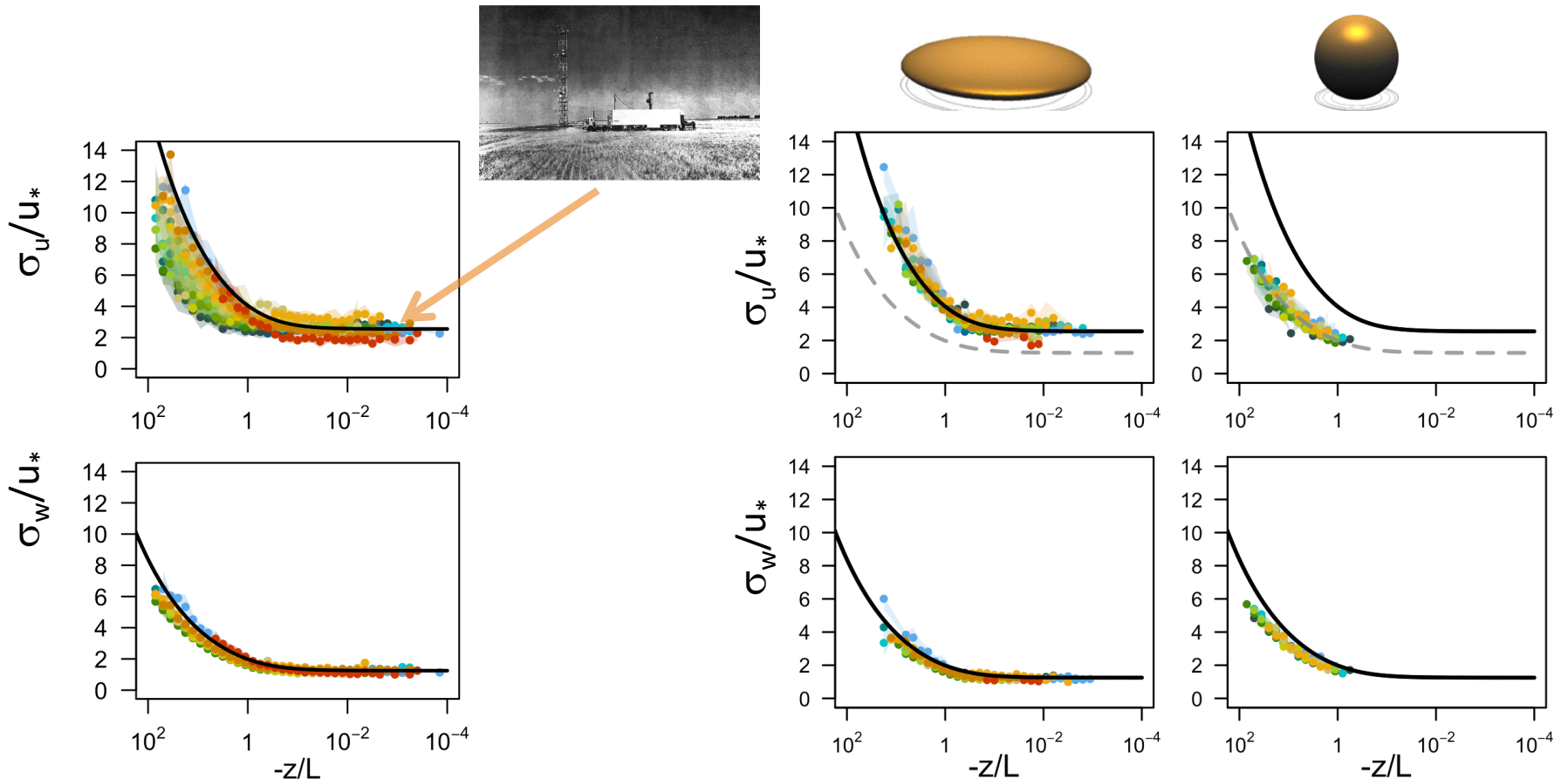
Stable



Stiperski et al. (2019)

Scaling in complex terrain

Standard deviation horizontal and vertical velocity

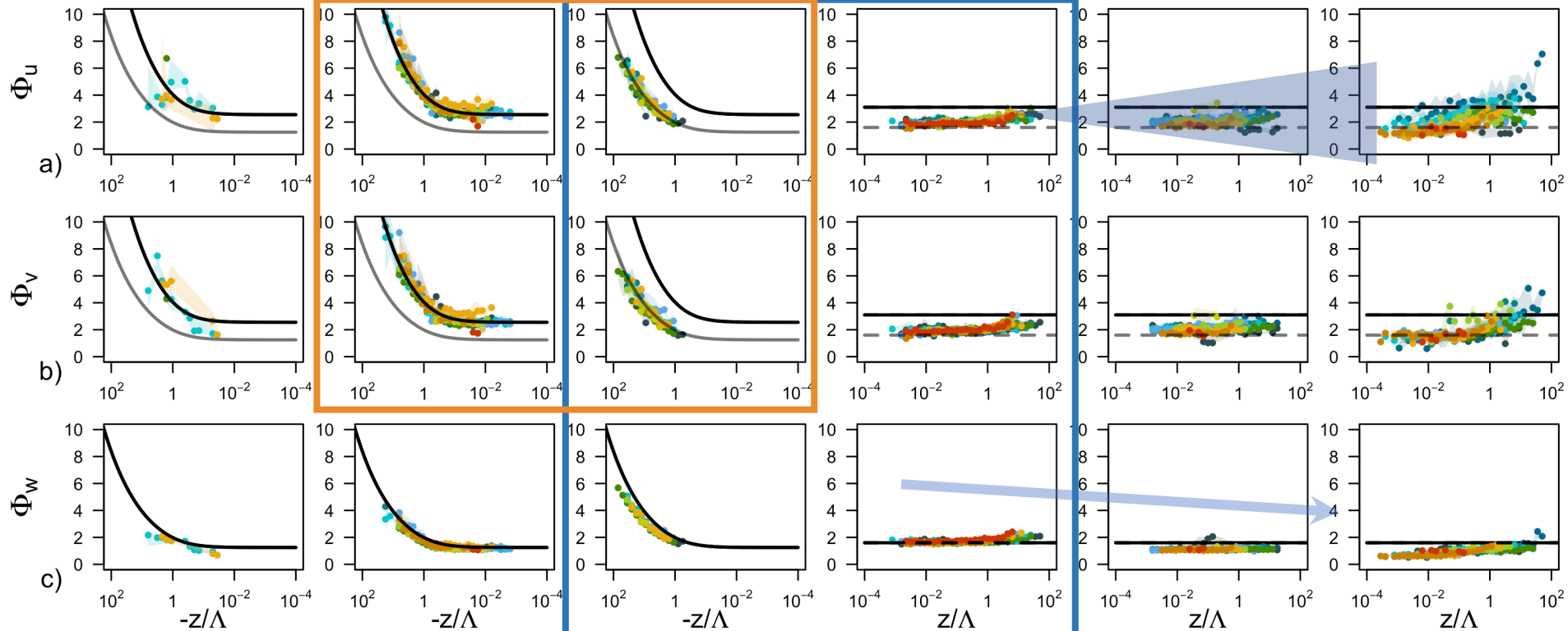


Stiperski et al. (2019)

Scaling in complex terrain

Unstable

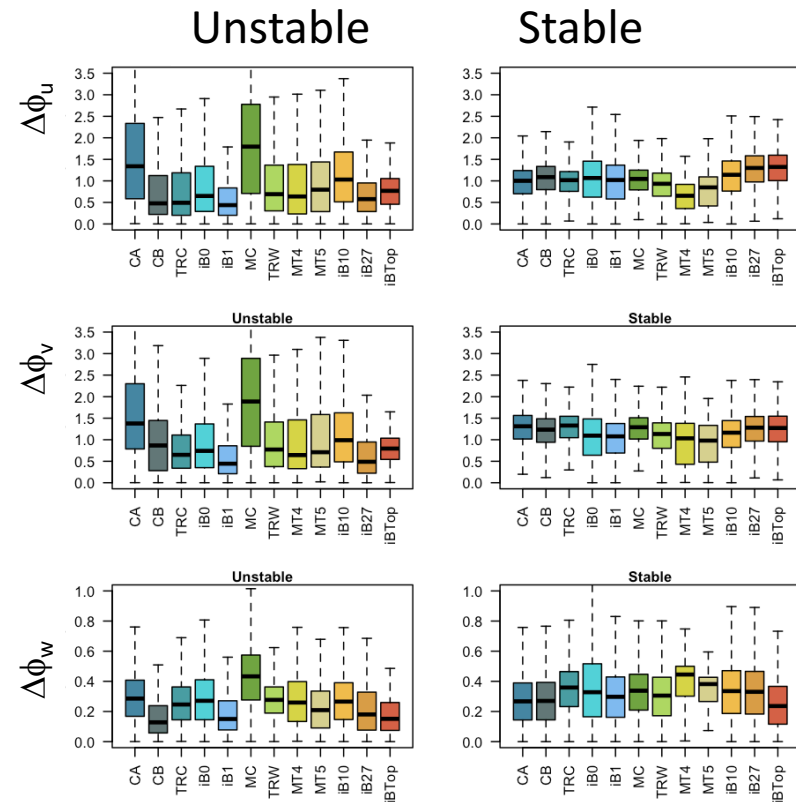
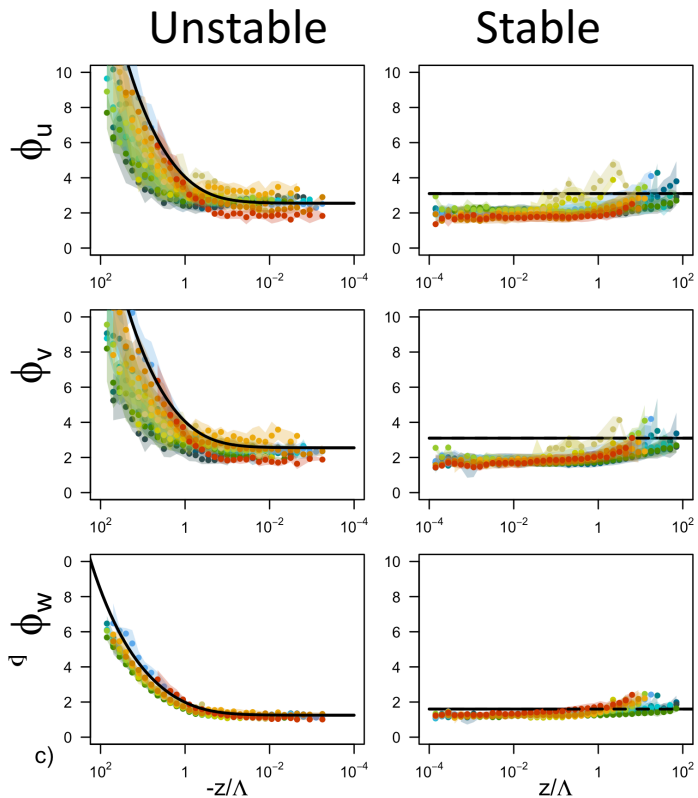
Stable



CASES-99 Cabauw T-RexC i-Box0 i-Box1 METCRAX T-RexW Materhorn4 Materhorn5 i-Box10 i-Box27 i-BoxTop

Stiperski et al. (2019)

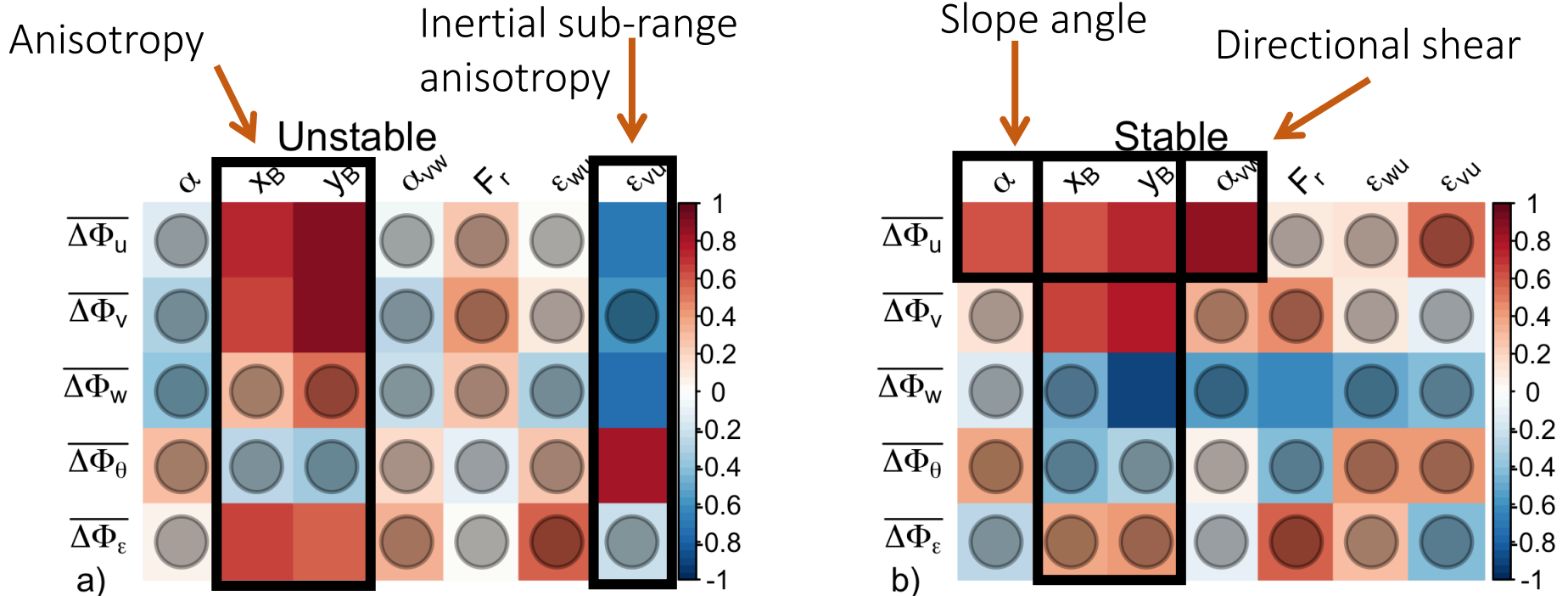
Scaling in complex terrain



We can define deviations from scaling curve as measure of process we are missing

Stiperski et al. (2019)

Scaling in complex terrain



Multilinear regression between deviation from scaling curve and

- Anisotropy
- Slope angle
- Directional shear

shows these are missing processes

Stiperski et al. (2019)

Summary: Similarity Theory

- 'direct approach' to determine turbulence state
- few (characteristic) scaling variables needed
- 'recipe' available
- Scaling regimes
 - Surface Layer (MOST) (→ free convection limit)
 - Mixed Layer
 - Local Scaling Layer (z-less Scaling)
- horizontally homogeneous: established
- complex terrain / heterogeneous surface: only evolving